Missing beauty of proton-proton interactions

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Faculty of Physics

הפקולטה לפיסיקה

In this talk

QGP signature is small systems

Overview of LHC results

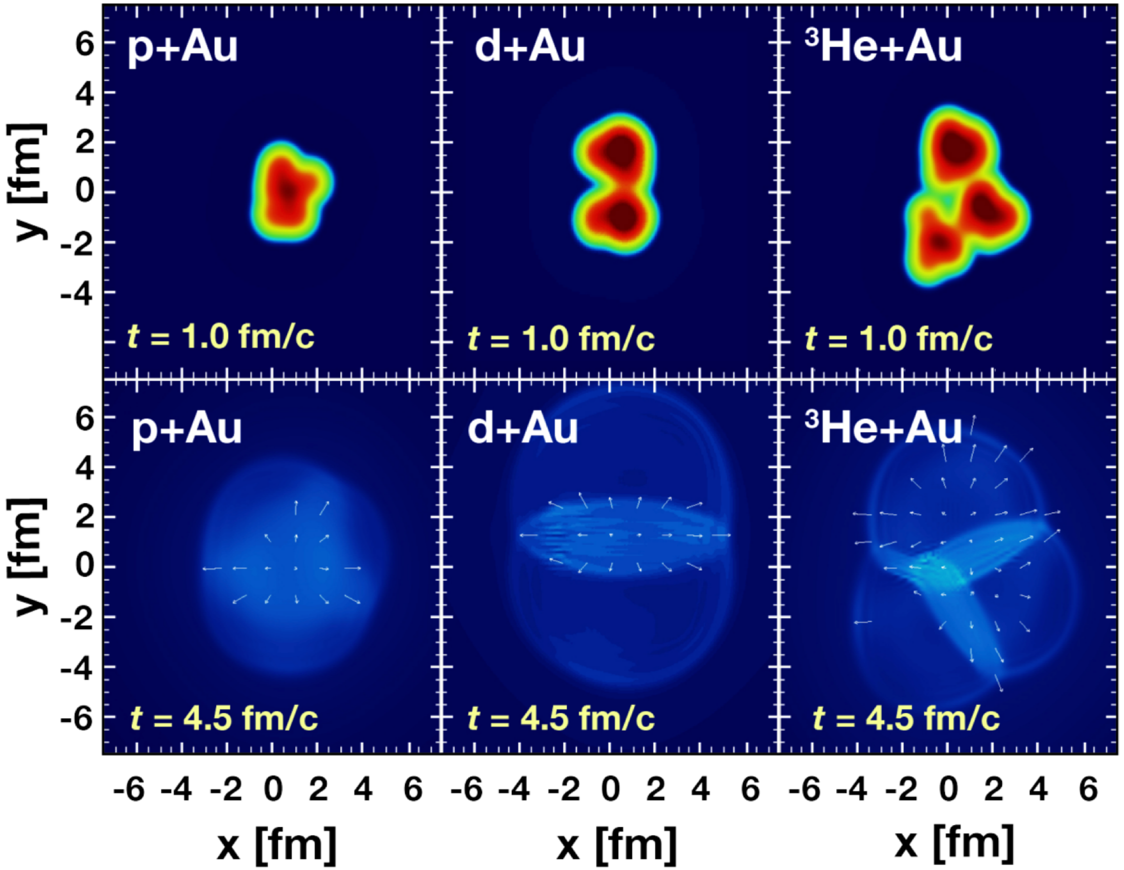
Global analysis of meson production in pp

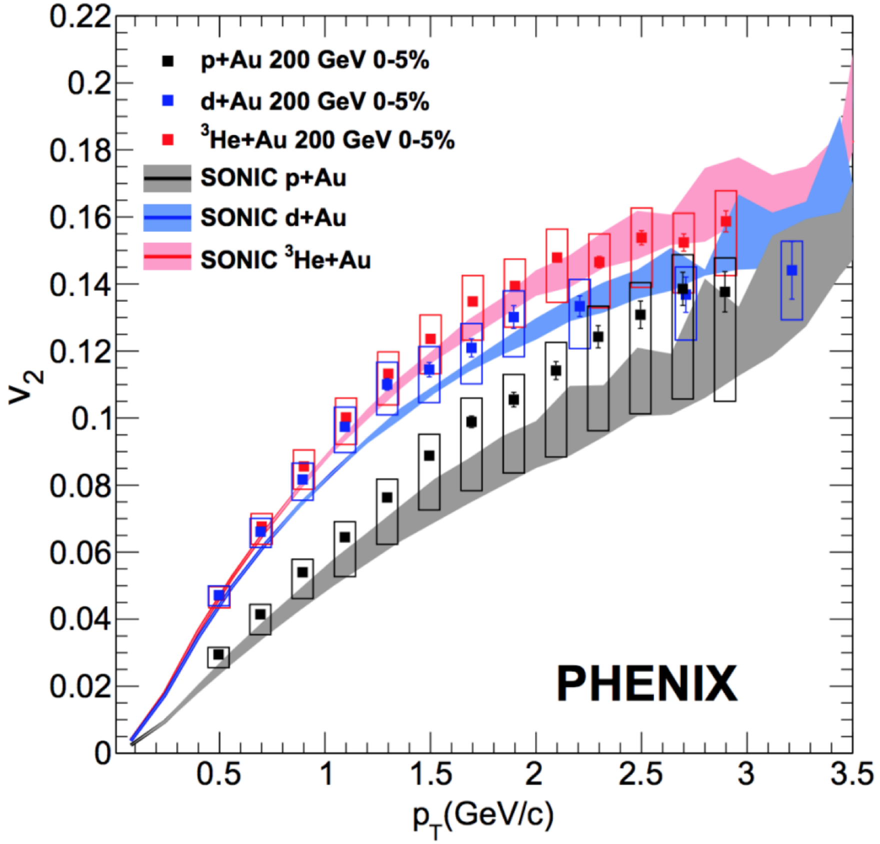
Bringing things together

Conclusions and questions

QGP signatures in small systems

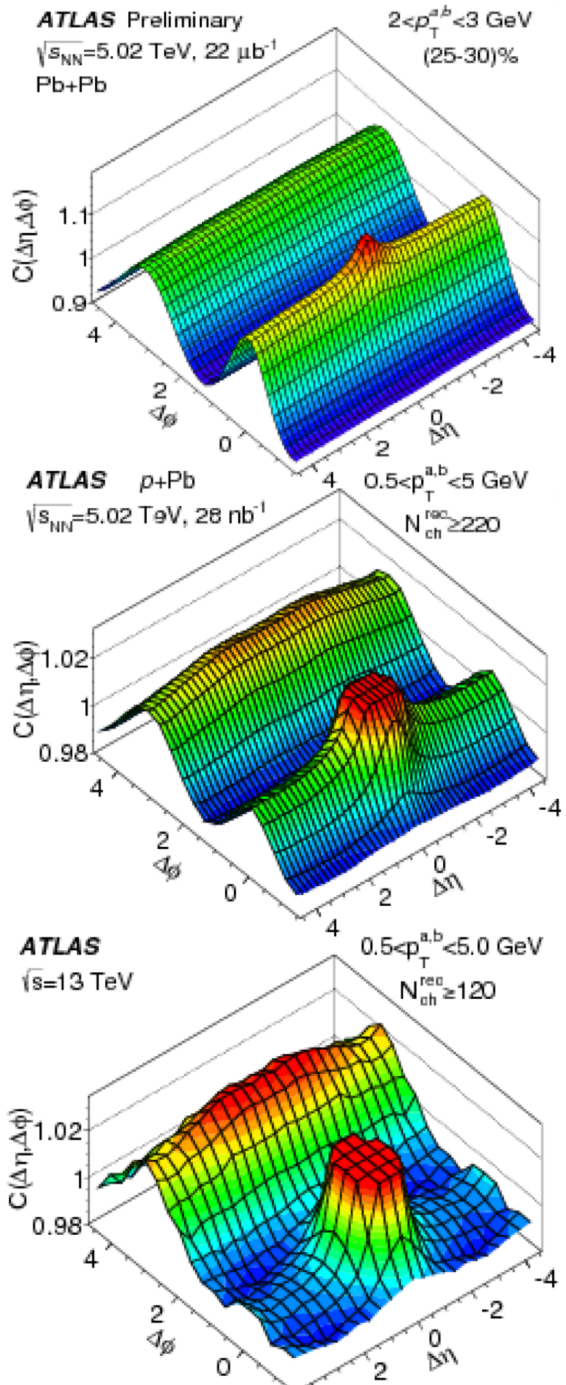
ARNPS 68 (2018) 211-235

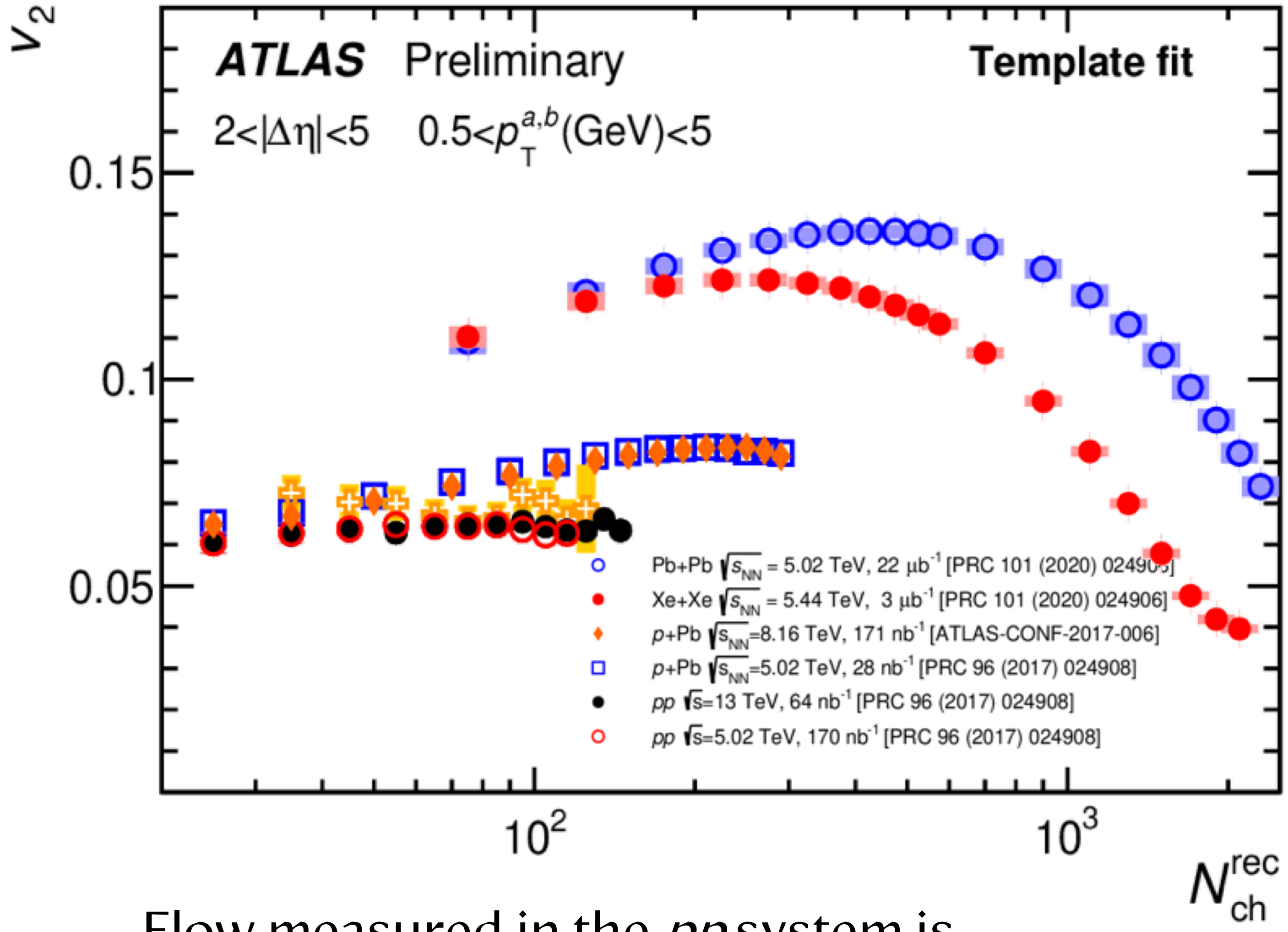




Accounting for geometry and using hydro model with the same η/s allows for simultaneous description of the flow in three different systems

QGP signatures in small systems

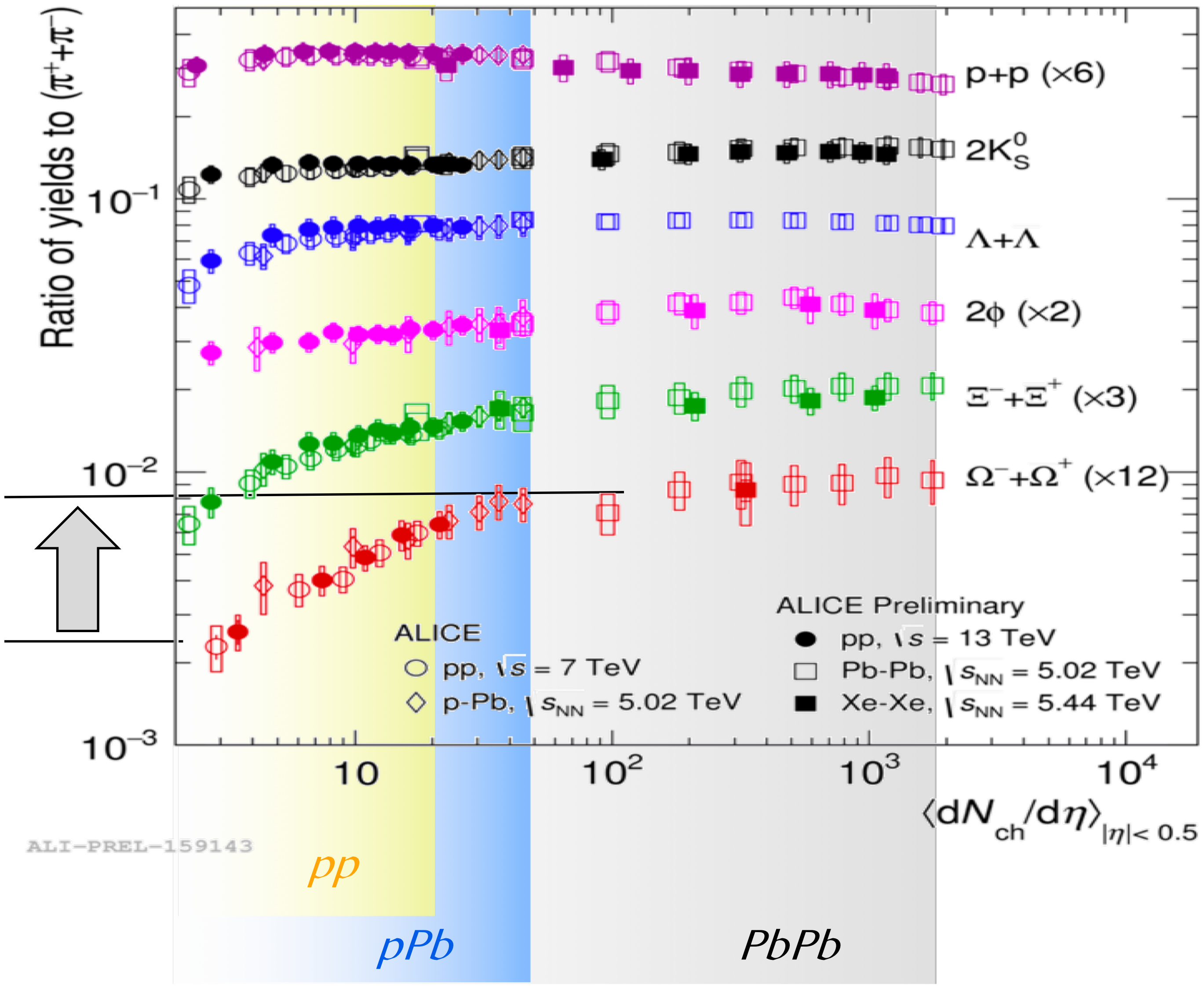




Flow measured in the *pp* system is comparable to much larger systems

QGP signatures in small systems

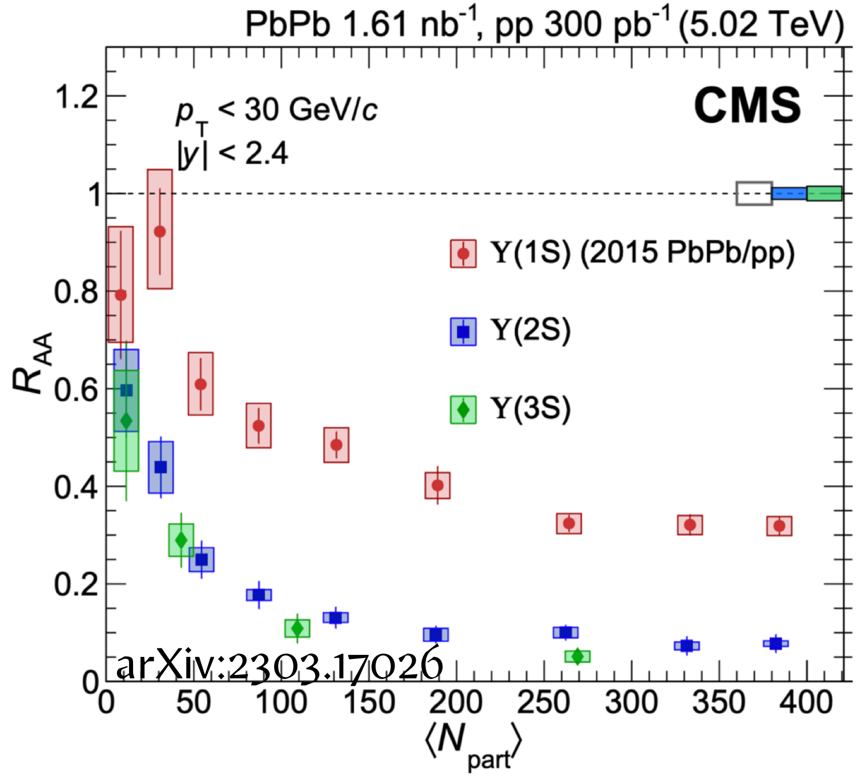
see NP **13**, 535–539 (2017)



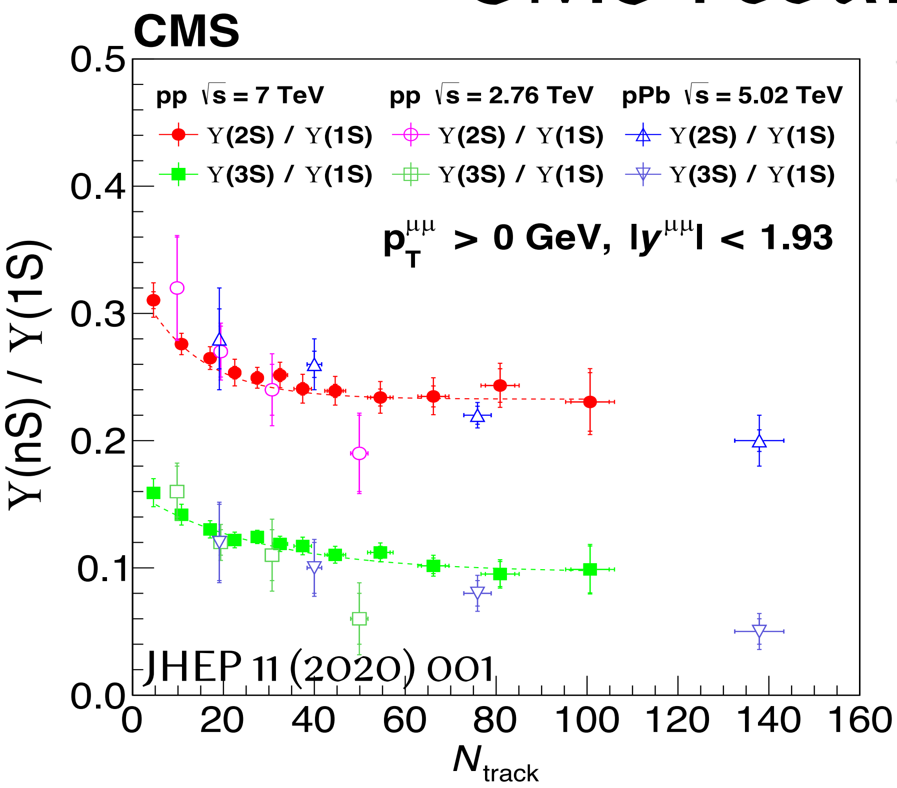
Strangeness enhancement happens in the range pf multiplicities of small systems

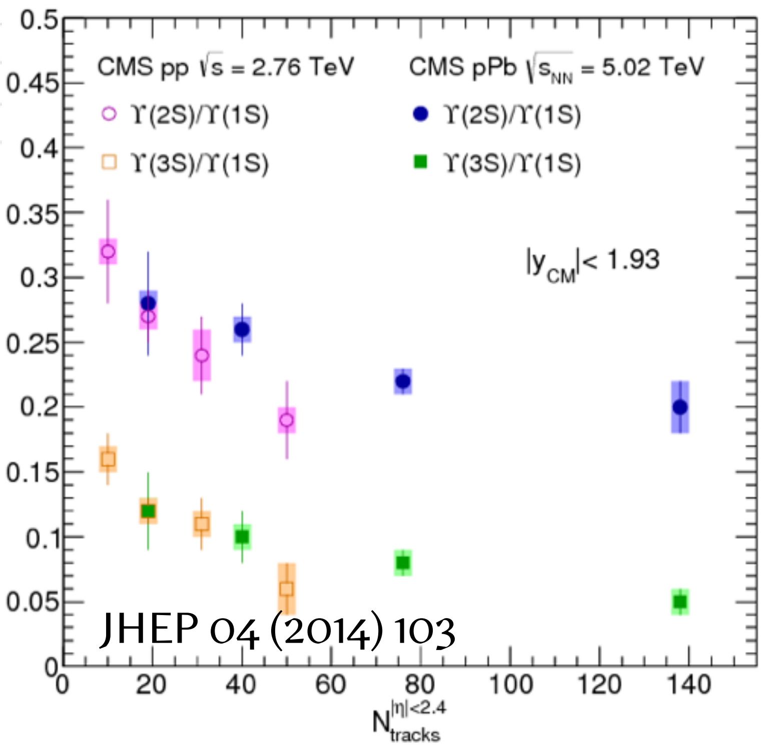
What about hard probes?

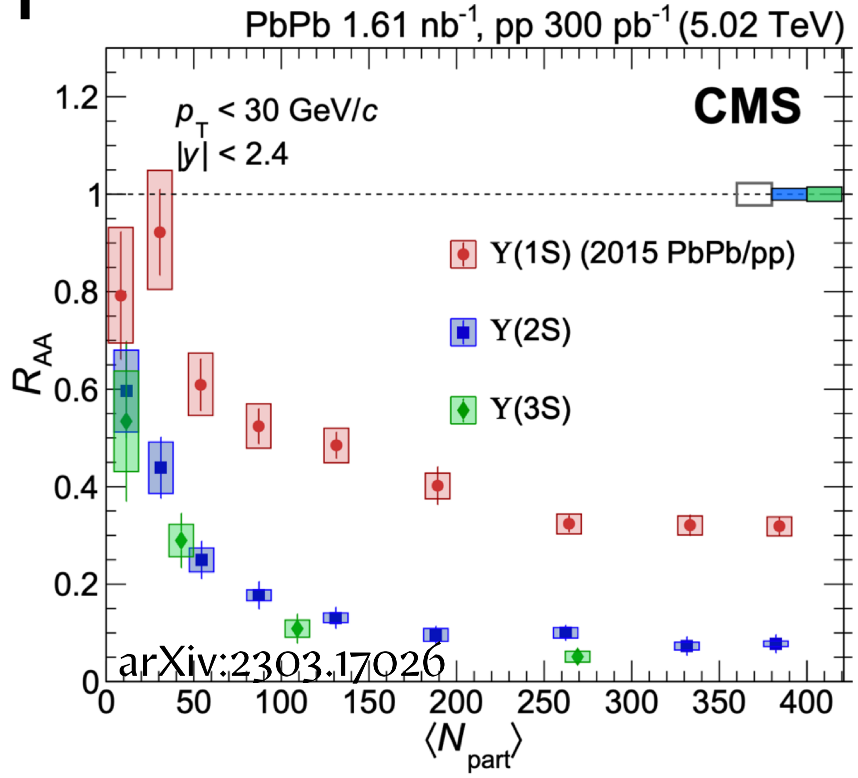
If we are to look for the most sensitive QGP hard probe the obvious suspect would be the $\Upsilon(nS)$ family…



CMS results with $\Upsilon(nS)$ suppression

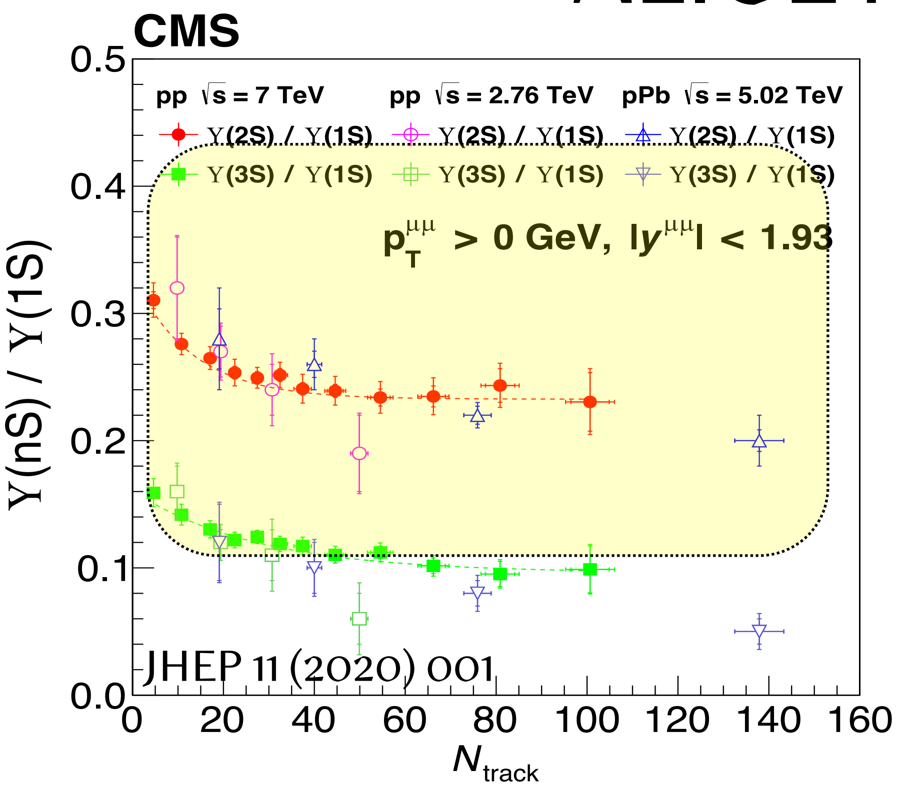


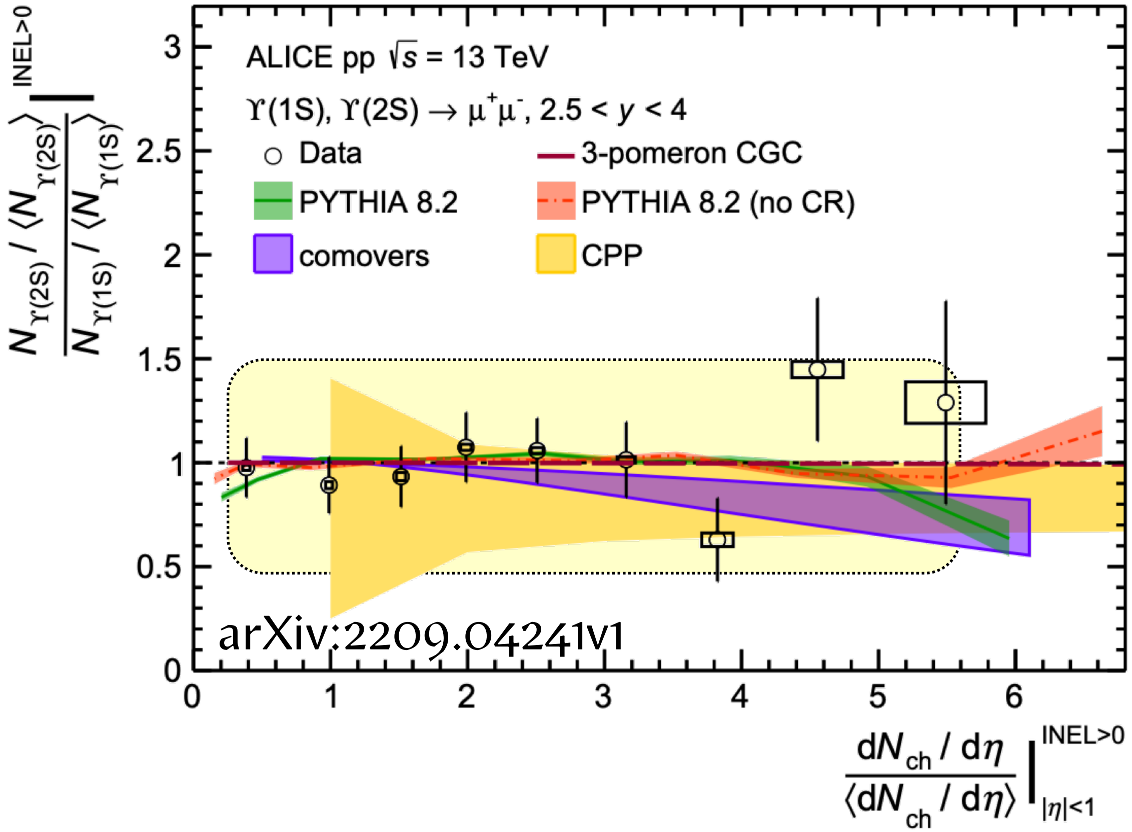




CMS: "It was concluded that the feed-down contributions cannot solely account for this feature. This is also seen in the present analysis, where the $\Upsilon(1S)$ meson is accompanied by about one more track on average ($\langle N_{track} \rangle = 33.9 \pm 0.1$) than the $\Upsilon(2S)$ ($\langle N_{track} \rangle = 33.0 \pm 0.1$), and about two more than the $\Upsilon(3S)$ ($\langle N_{track} \rangle = 32.0 \pm 0.1$). [...] On the other hand, it is also true that, if we expect a suppression of the excited states at high multiplicity, it would also appear as a shift in the mean number of particles for that state (because events at higher multiplicities would be missing)."

ALICE result with a rapidity gap



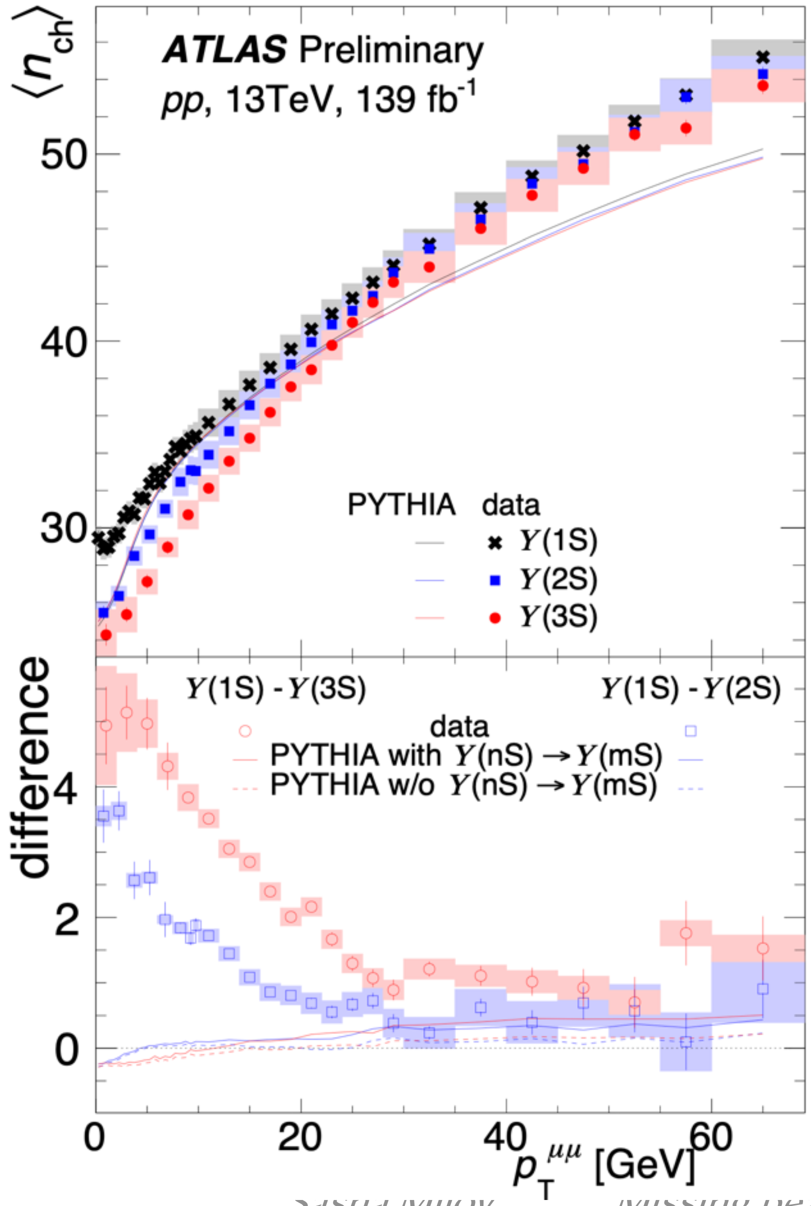


ALICE result on forward $\Upsilon(2S)/\Upsilon(1S)$ vs. tracks at midrapidity shows rather different behavior when quarkonia and multiplicity measured at different rapidities

Statistics is too low to warrant any gap dependence

Multiplicity dependence on Υ–momentum

ATLAS-CONF-2022-023



Multiplicity is different for different $\Upsilon(nS)$ states

Can't be explained by feed downs or p_T, conservation

Pythia has no effect like this

At the lowest p_T, where the effect is the strongest:

$\Upsilon(1S) - \Upsilon(2S)\ \Delta\langle n_{ch}\rangle = 3.6 \pm 0.4$ 12% of $\left\langle n_{ch}^{\Upsilon(1S)}\right\rangle$

$\Upsilon(1S) - \Upsilon(3S)\ \Delta\langle n_{ch}\rangle = 4.9 \pm 1.1$ 17% of $\left\langle n_{ch}^{\Upsilon(1S)}\right\rangle$

It diminishes with p_T, but remains visible at 20–30 GeV
And actually above that as well

Where the differences are coming from?

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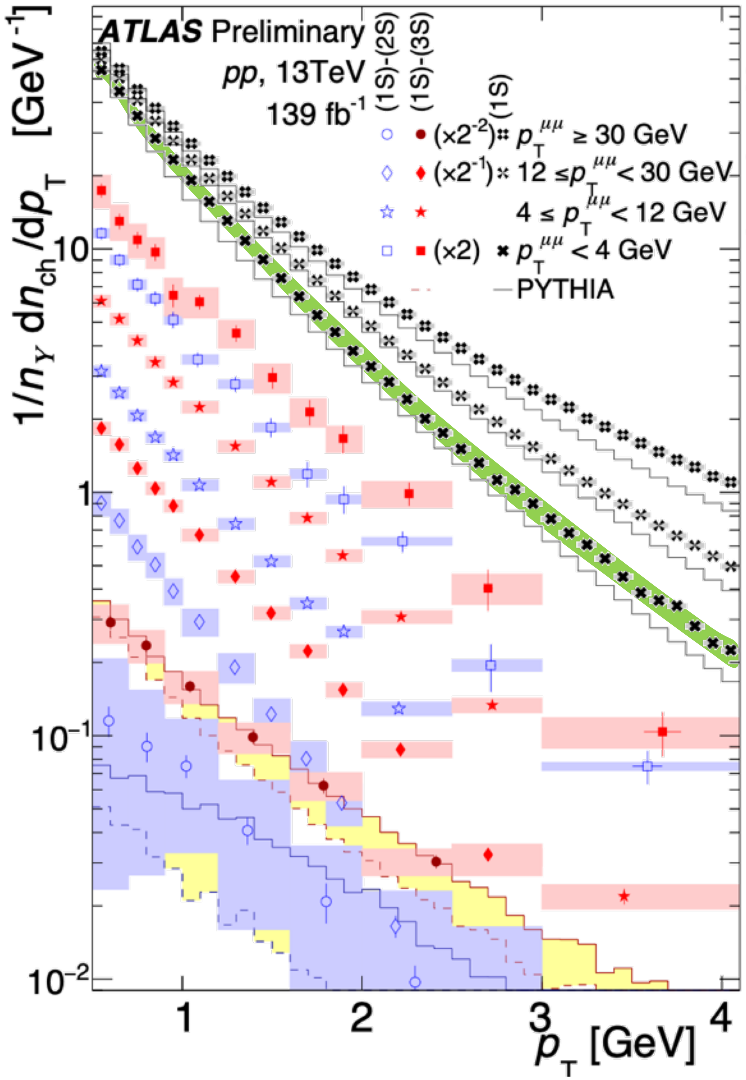
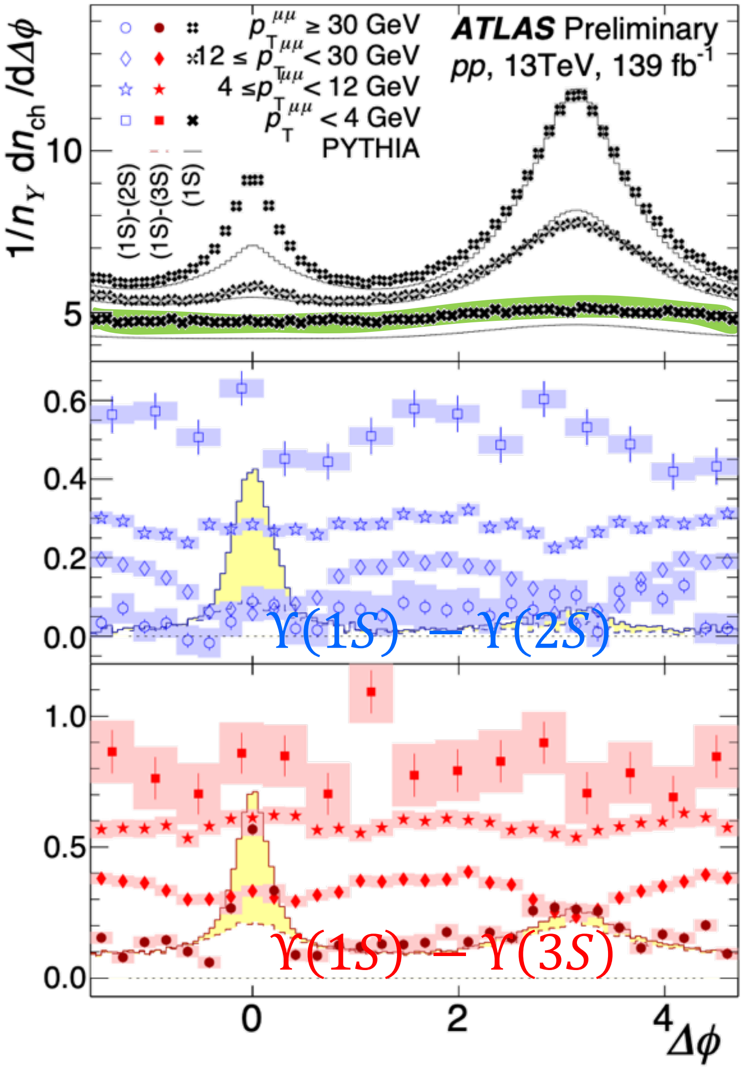
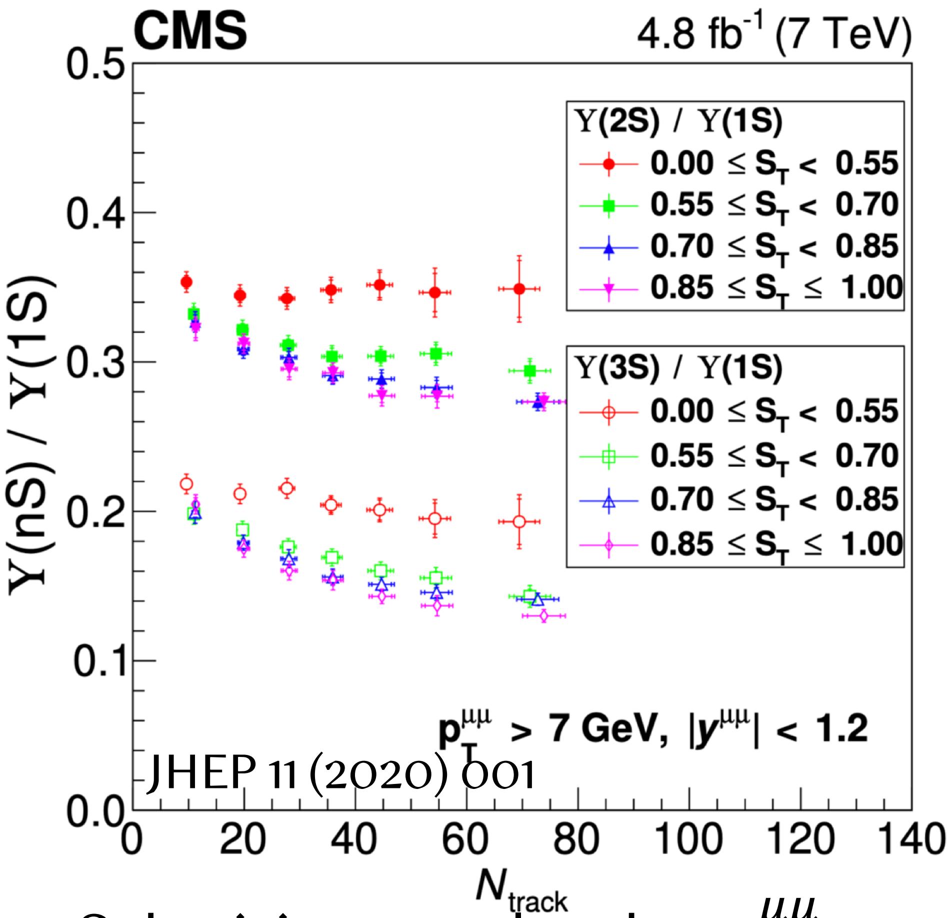




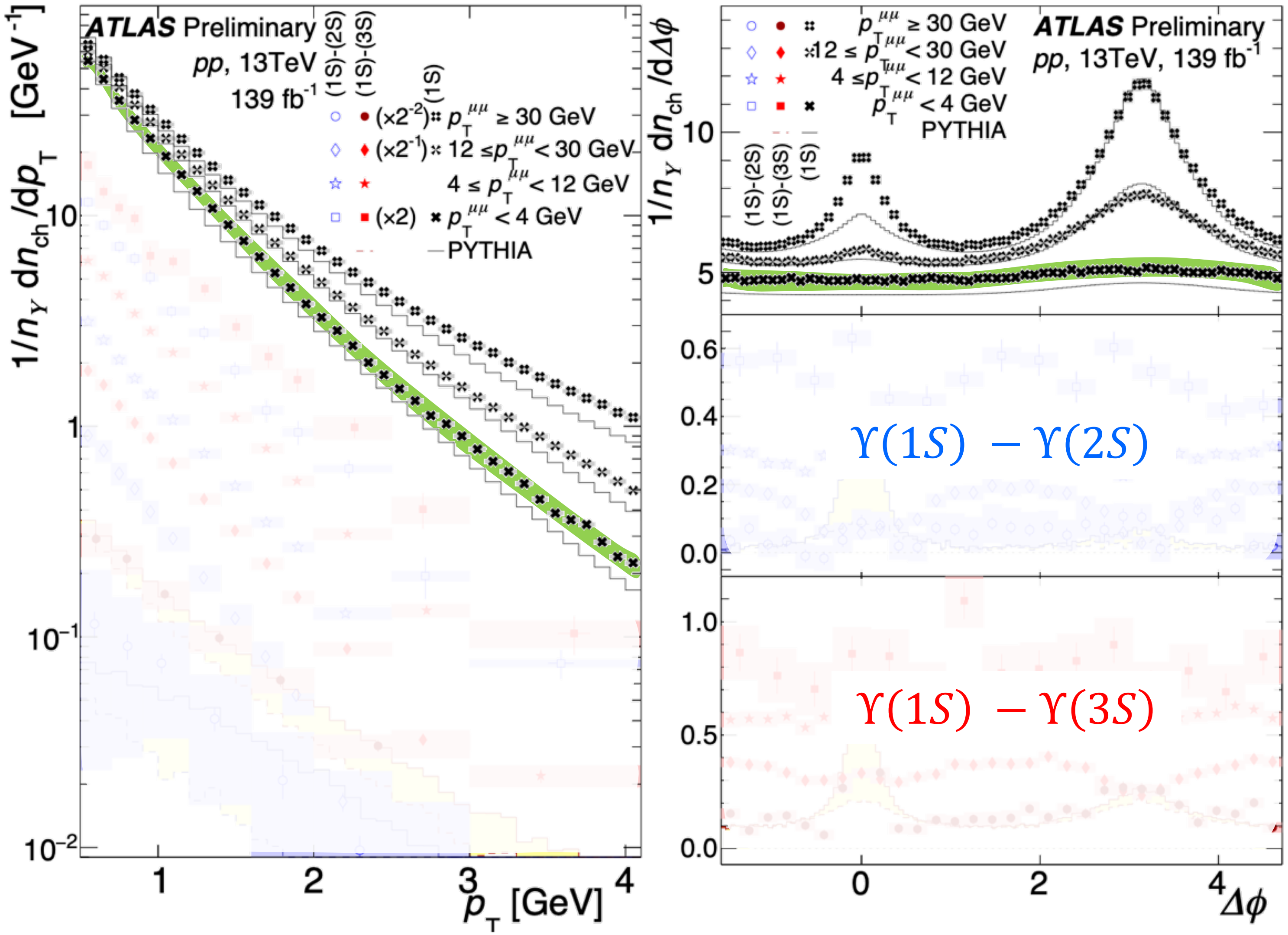


Spherisity, correlated to $p_T^{\mu\mu}$

$S_T < 1$ jet-like

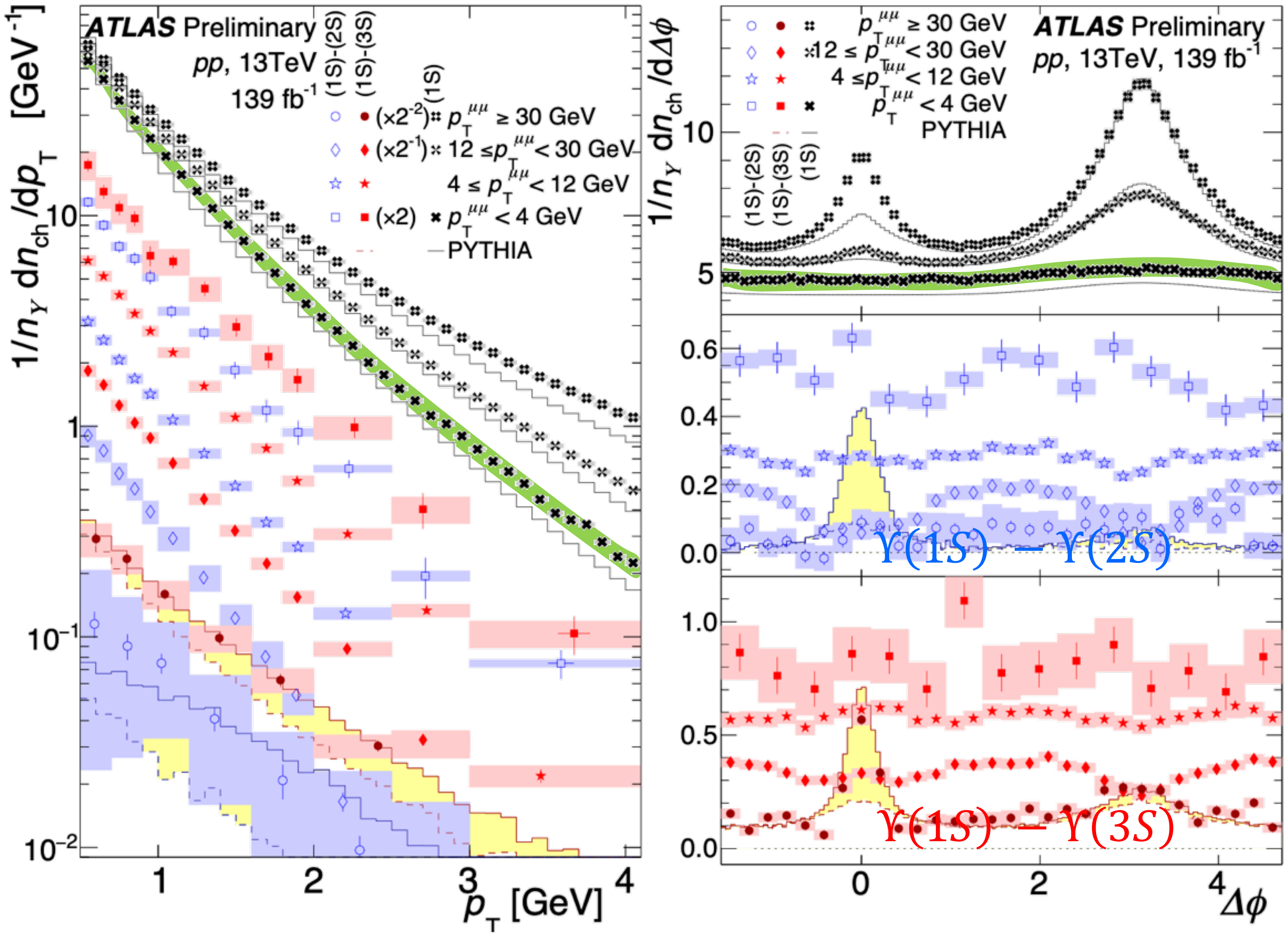
$S_T \sim 1$ no jet ← stronger effect

Kinematic distributions of $\Upsilon(1S)$



One cannot measure the UE, but $p_T < 4$ GeV is the closest to it, jet part that is correlated to $\Upsilon(nS)$

Kinematic distributions of the differences



One cannot measure the UE, but p_T < 4 GeV is the closest to it, jet part that is correlated to $\Upsilon(nS)$

Subtracted distributions look like UE at rather high $\Upsilon(nS)$ p_T. At the highest p_T there are feed-downs

Away from jets there are regions with charged particles

The effect is related to the UE

Is it a deficit for $\Upsilon(nS)$ or an excess for $\Upsilon(1S)$?

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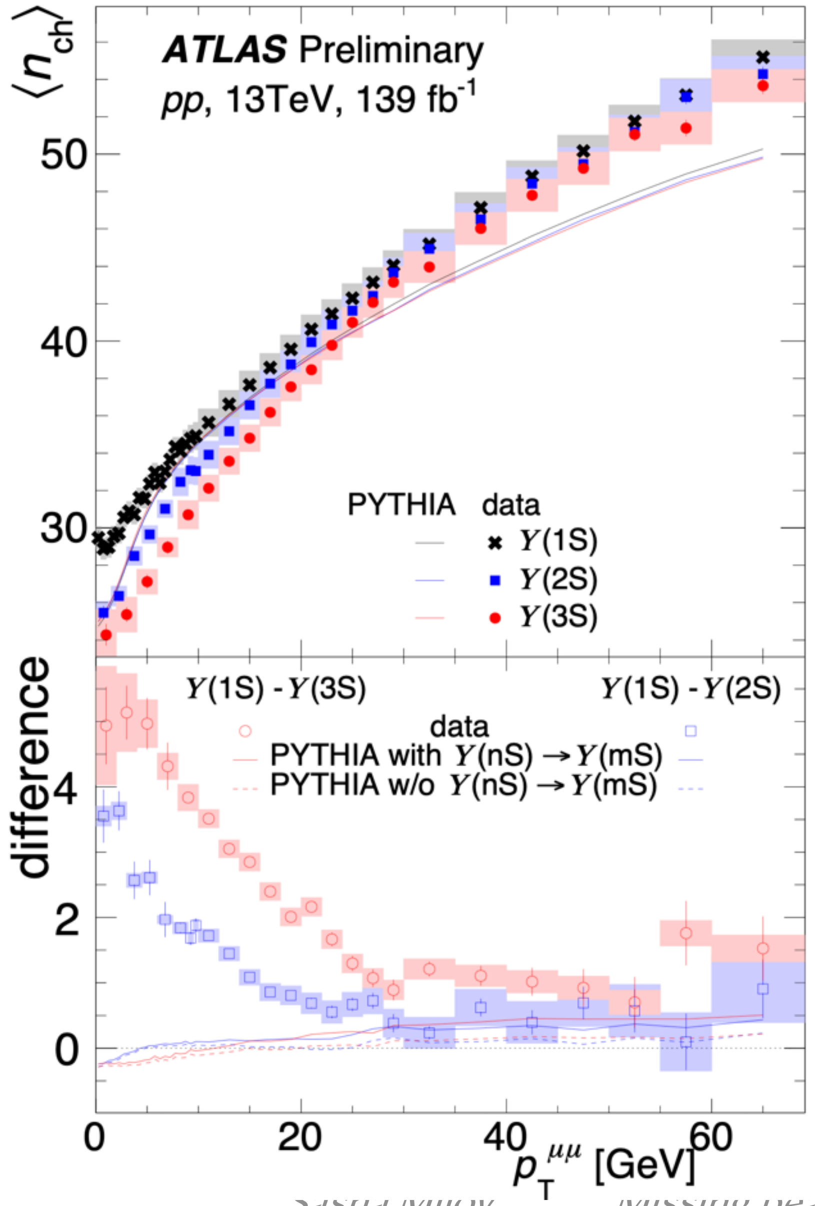


How large is the UE in the presence of $\Upsilon(nS)$?

Inclusive *pp* collisions: $\langle n_{ch} \rangle \approx 14$
Drell-Yan with $40\ \mathrm{GeV} < m < m_Z$ $\langle n_{ch} \rangle = 24 - 28$
Jets with leading particles $m < \frac{1}{2} m_\Upsilon$ $\langle n_{ch} \rangle \approx 27$

On the other hand, a p_T – dependence of the $\Delta\langle n_{ch} \rangle$ points to the modification of p_T spectrum. What shall be the p_T spectrum of $\Upsilon(nS)$?

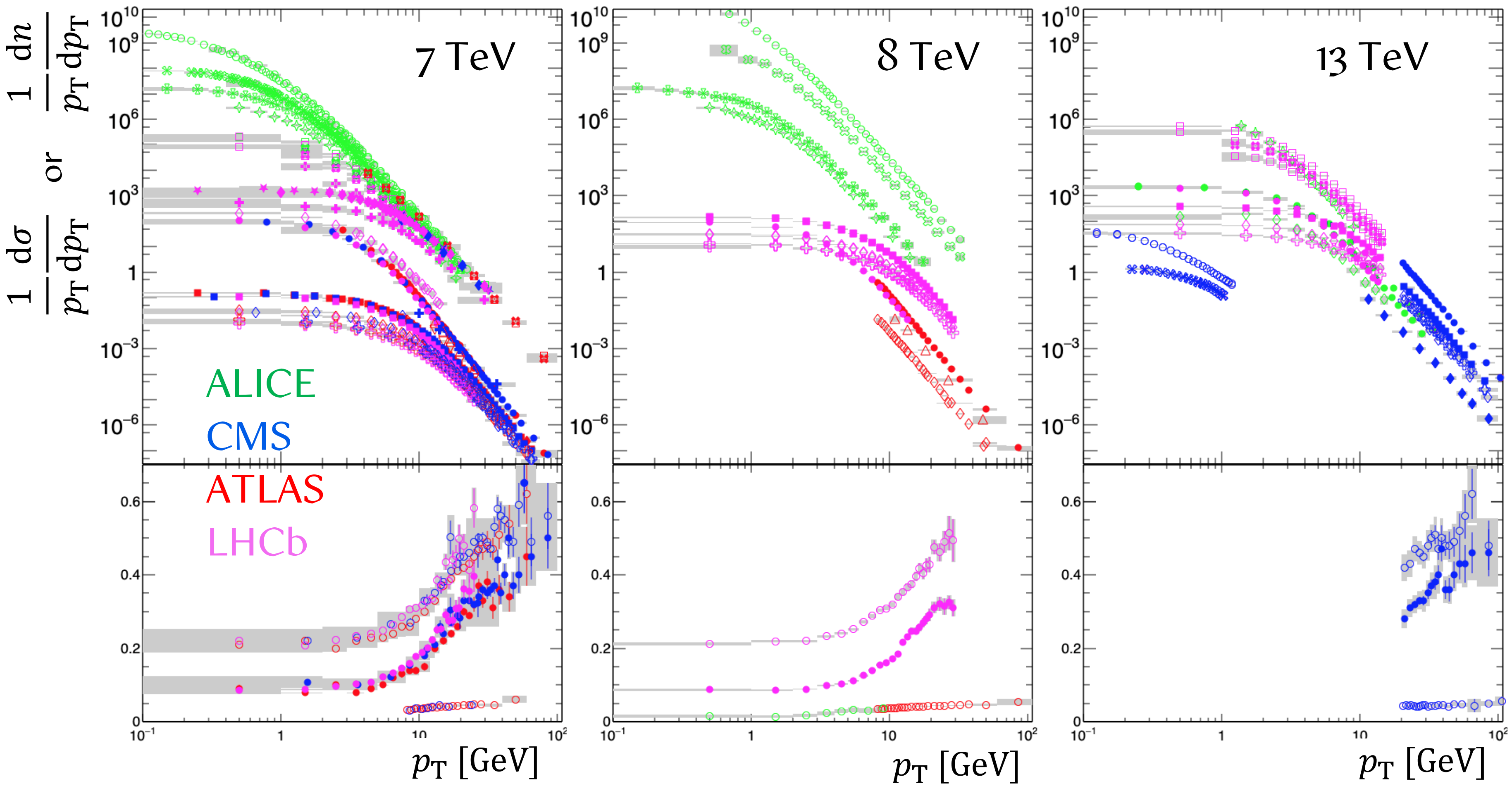
Basic assumption:

If particles have the same quark content and the same mass, they must have the same kinematics.

For small Δm between particles one can use m_T – scaling

LHC data for m_T - scaling

$$\frac{d\sigma}{dm_T} \propto \left[1+\frac{m_T}{nT}\right]^{-n}$$



Only mesons

4 LHC experiments

$\sqrt{s} = 7, 8, 13$ TeV

18 species + iso-partners

72 data samples with 1509 data points

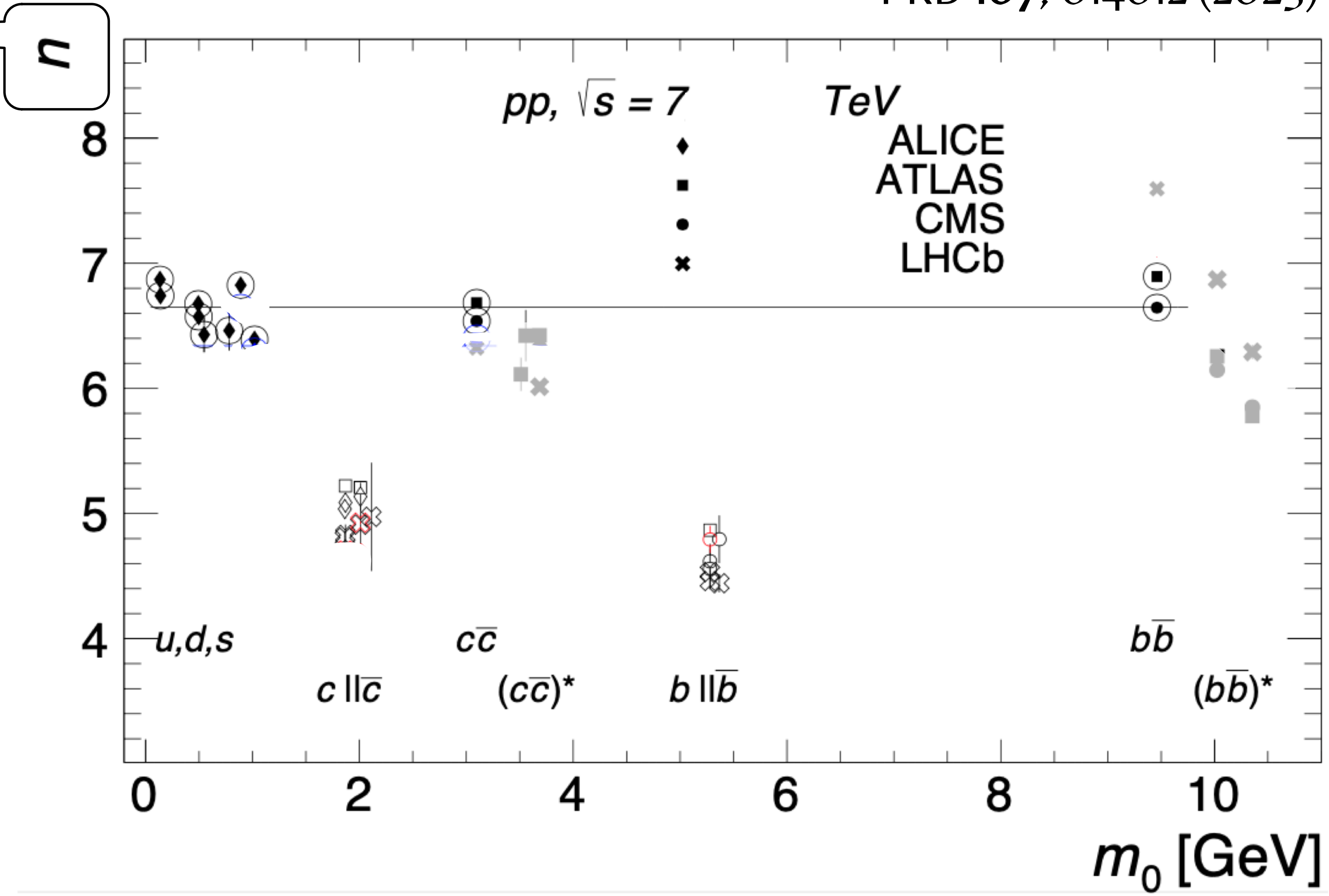
15 quarkonia ratios with 327 data points

T is fixed to 254 MeV

Spectra at high p_T

PRD **107**, 014012 (2023)

$$\frac{d\sigma}{dm_T} \propto \left[1 + \frac{m_T}{nT}\right]^{-n}$$



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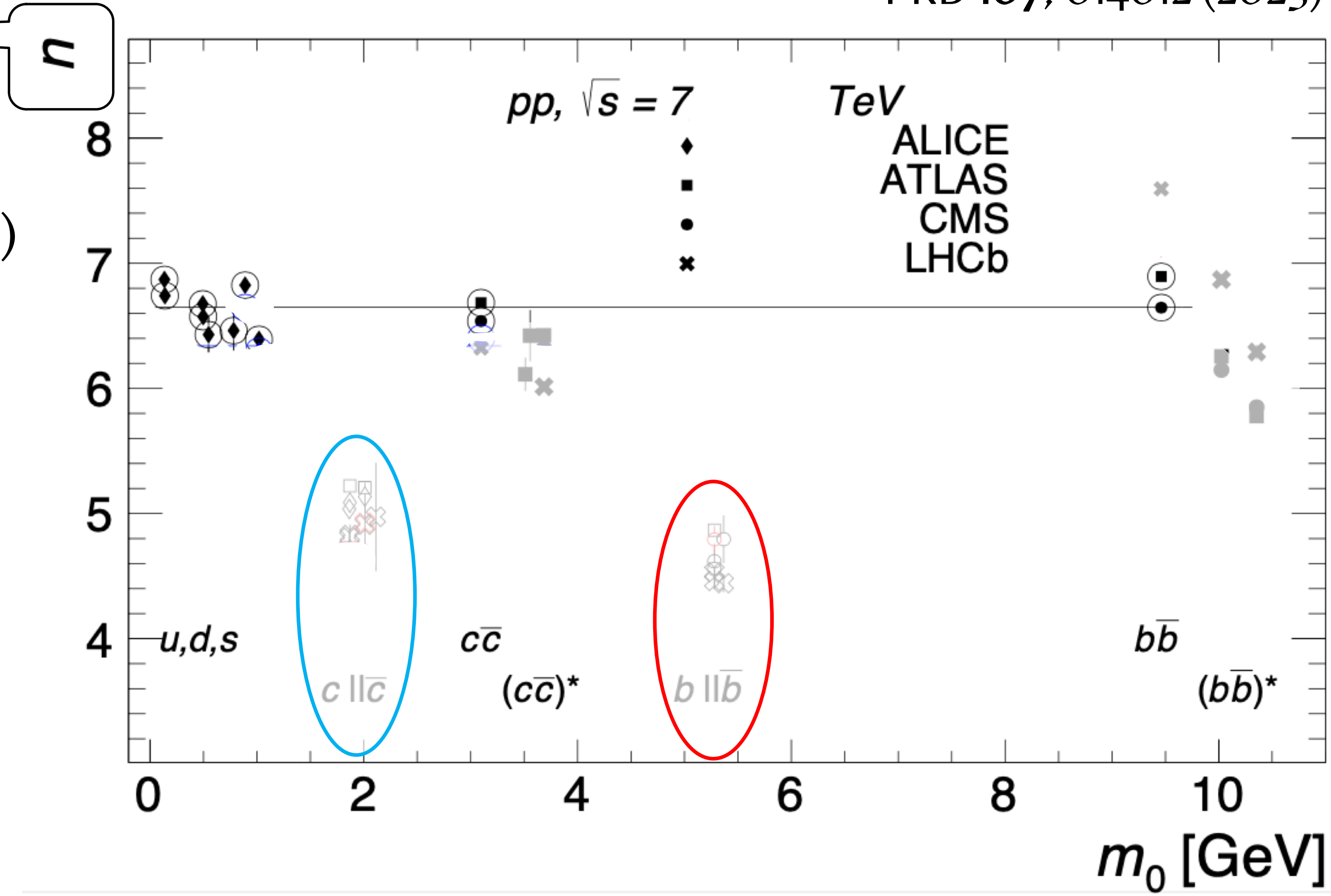

Open flavor mesons ($c||\bar{c}$ and $b||\bar{b}$) has harder spectra (lower n)

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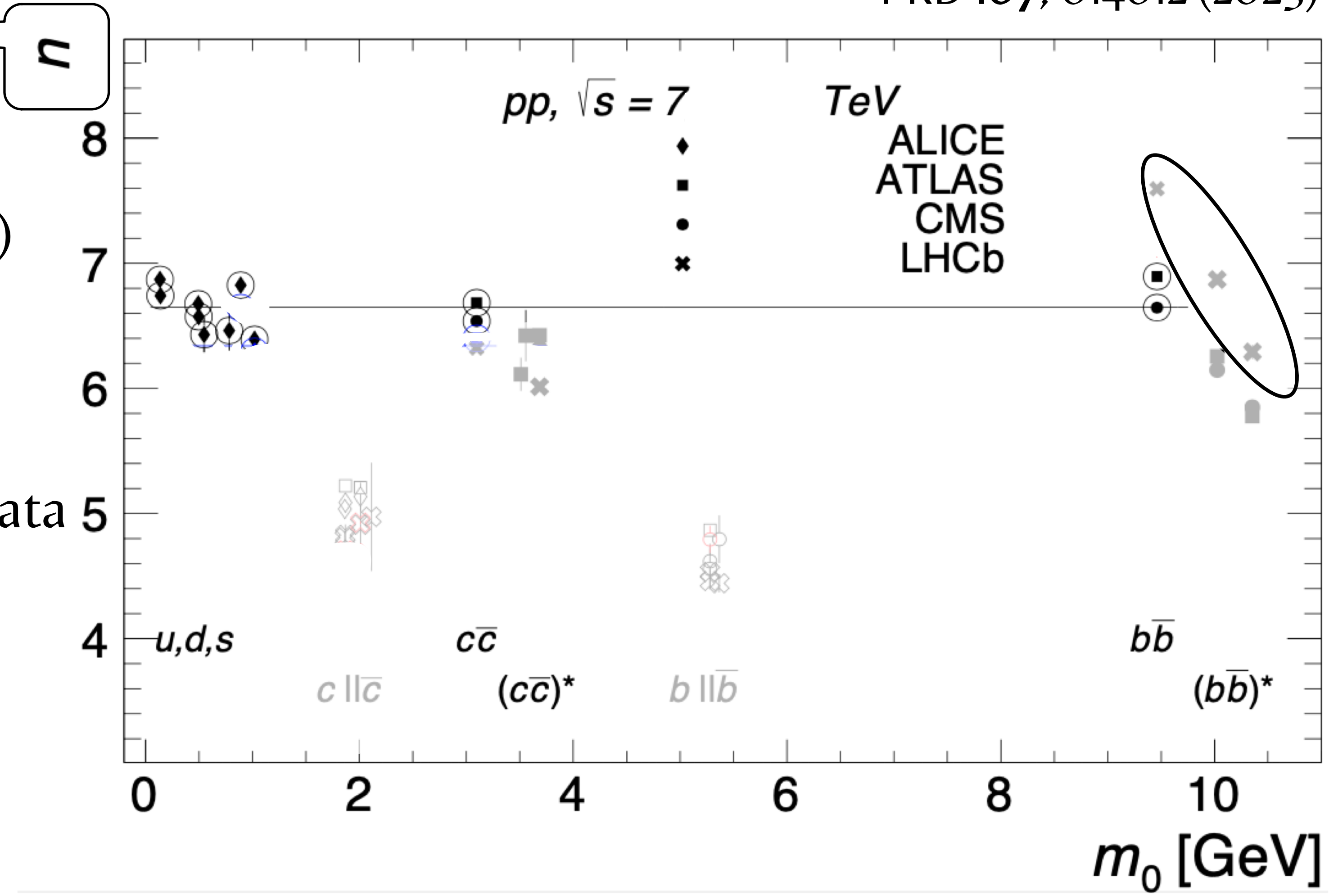
LHCb data (high-rapidity) are typically higher than midrapidity data



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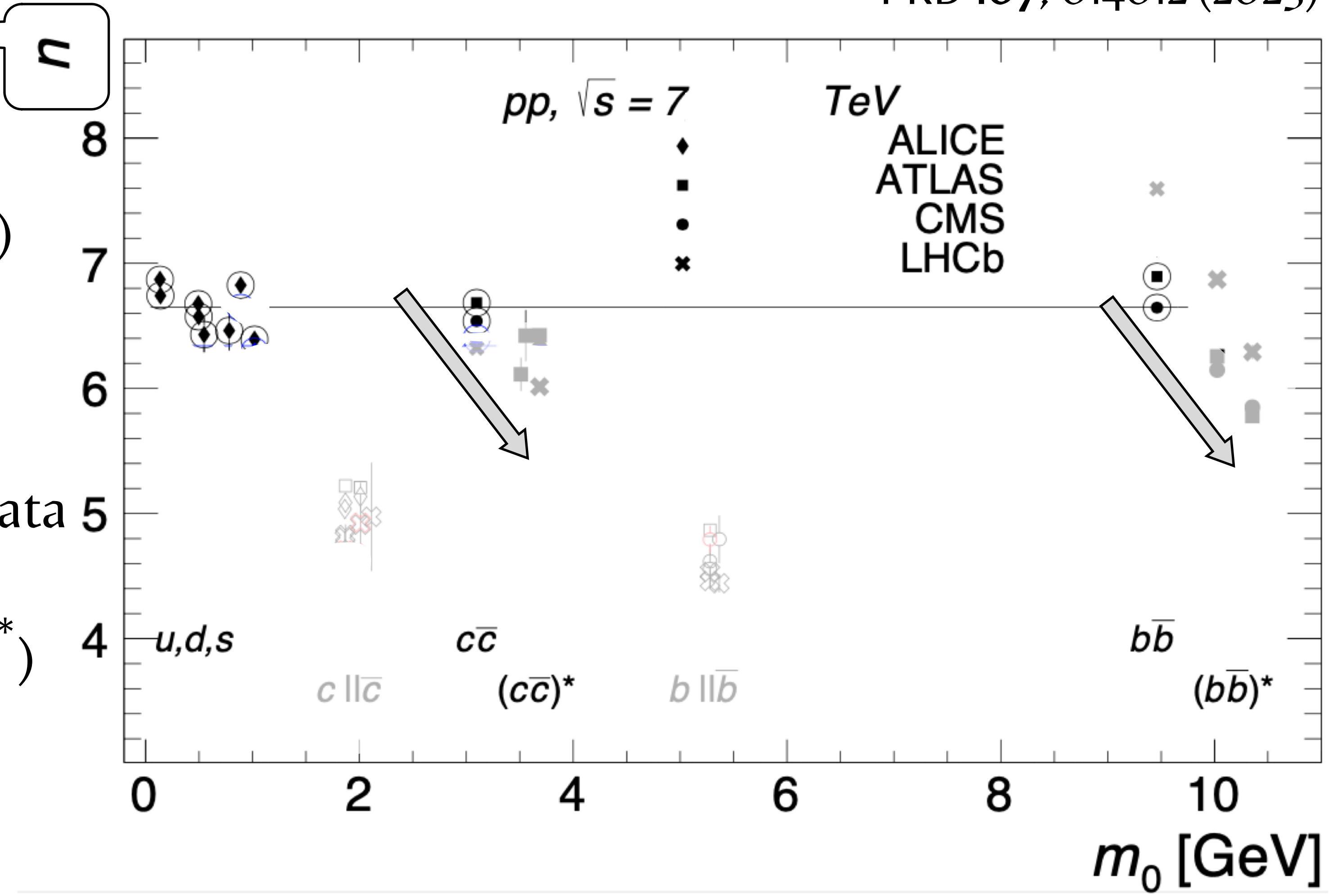
LHCb data (high-rapidity) are typically higher than midrapidity data

Excited quarkonia ($(c\bar{c})^*$ and $\left(b\bar{b}\right)^*$) have lower n

u,d,s & $q\bar{q}$ are fit simultaneously
$n = 6.65$ $\sqrt{s} = 7$ TeV



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Spectra at high p_T

$$\frac{d\sigma}{dm_T} \propto \left[1 + \frac{m_T}{nT}\right]^{-n}$$

n

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Excited quarkonia ($(c\bar{c})^*$ and $(b\bar{b})^*$) have lower n

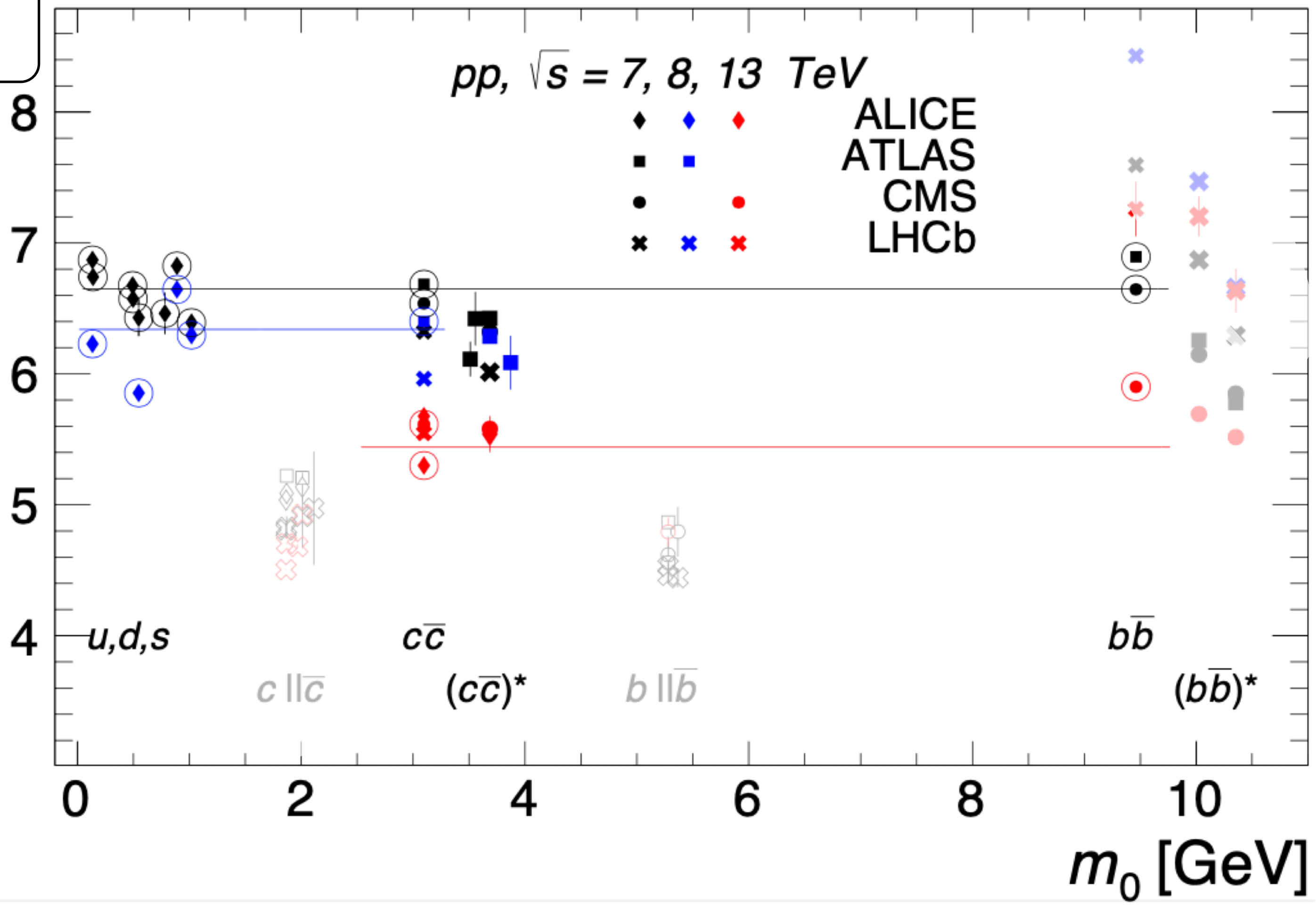
u,d,s & $q\bar{q}$ are fit simultaneously

$n = 6.65$ $\sqrt{s} = 7$ TeV

$n = 6.34$ $\sqrt{s} = 8$ TeV

$n = 5.44$ $\sqrt{s} = 13$ TeV

PRD **107**, 014012 (2023)



T is fixed to 254 MeV

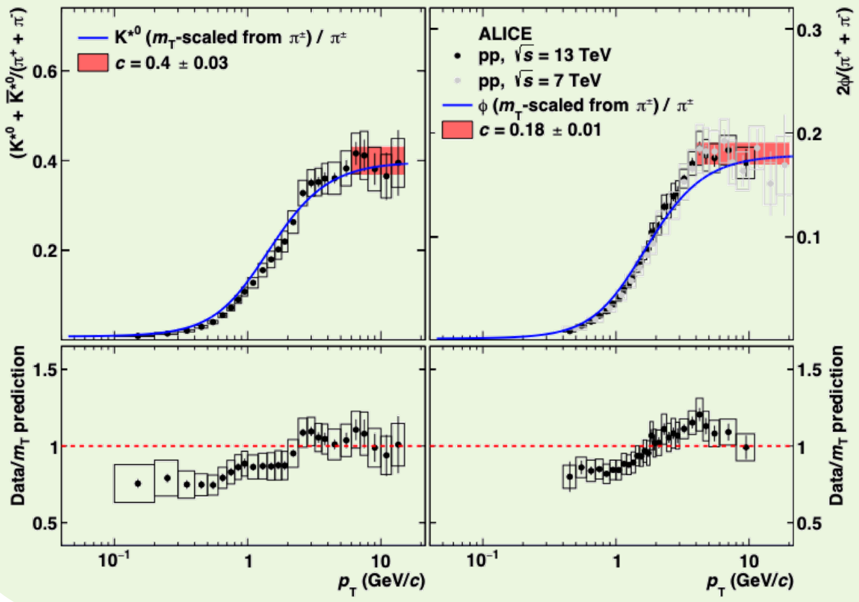
Quarkonia ratios: expected & measured

Normalized at $p_{\mathrm{T}} > 50$ GeV

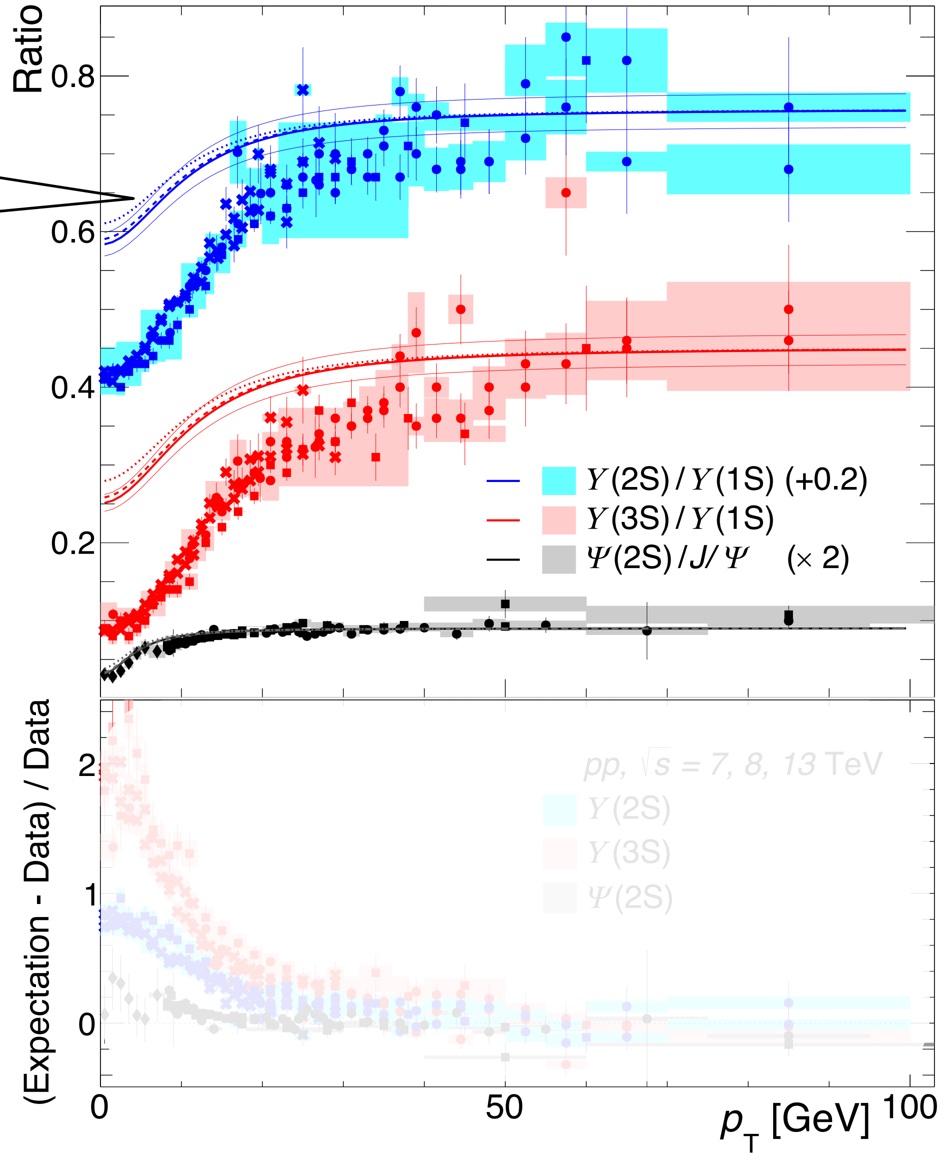
$$\lim_{\Delta m, p_T \ll m_{q\bar{q}}} \left[\frac{nT + \sqrt{p_{\mathrm{T}}^2 + (m_{q\bar{q}} + \Delta m)^2}}{nT + \sqrt{p_{\mathrm{T}}^2 + m_{q\bar{q}}^2}} \right]^{-n} = 1 - \frac{n\Delta m}{nT + m_{q\bar{q}}}$$

Measured $\Upsilon(nS)/\Upsilon(1S)$ are not as derived

No known effects can bridge differences for $(b\bar{b})^*$



An example
EPJC81 (2021) 256



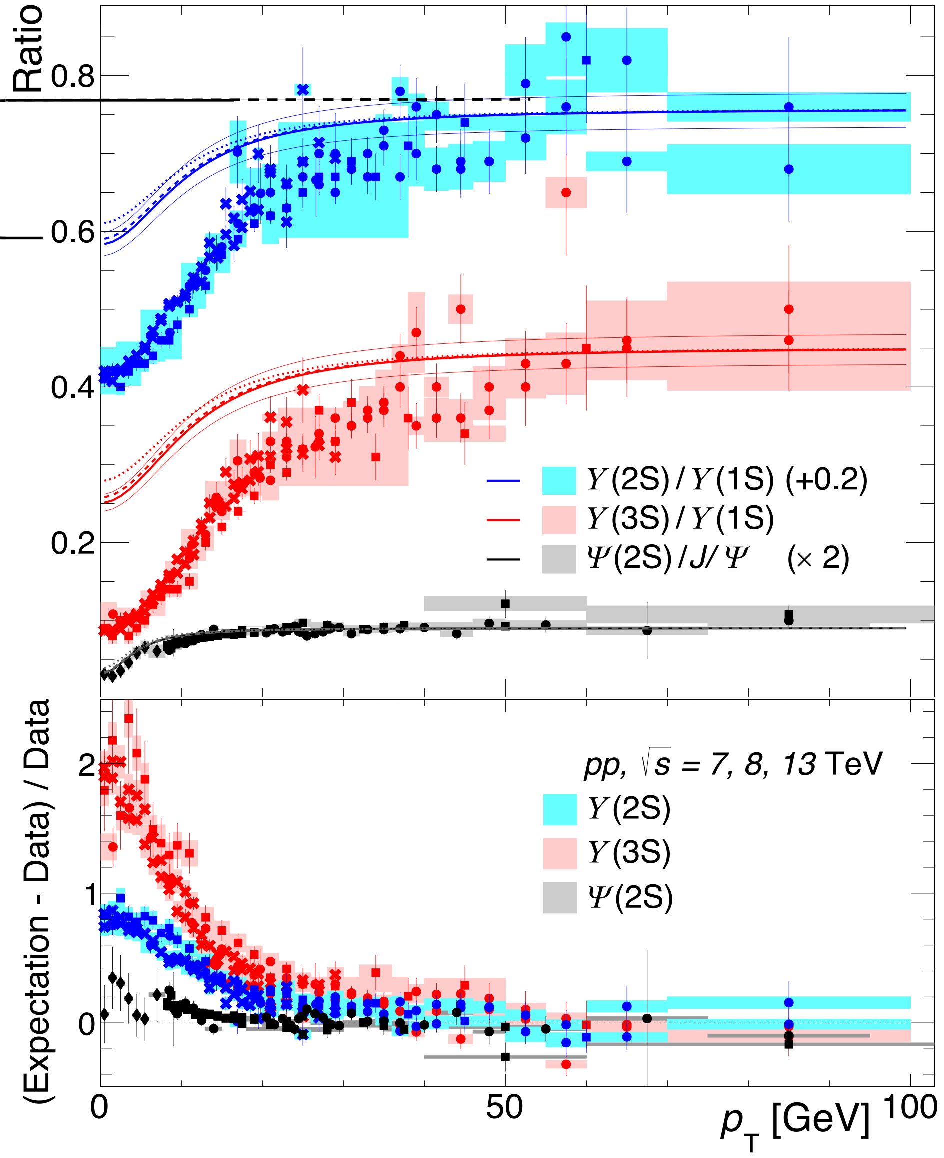
Quarkonia ratios: expected & measured

Normalized at $p_T > 50$ GeV

Measured ratios are not as derived $1 - \frac{n\Delta m}{nT + m_{q\bar{q}}}$

No known effects can bridge differences for $(b\bar{b})^*$

$$\text{Missing beauty} = \frac{\text{Expected}}{\text{Measured}} - 1$$



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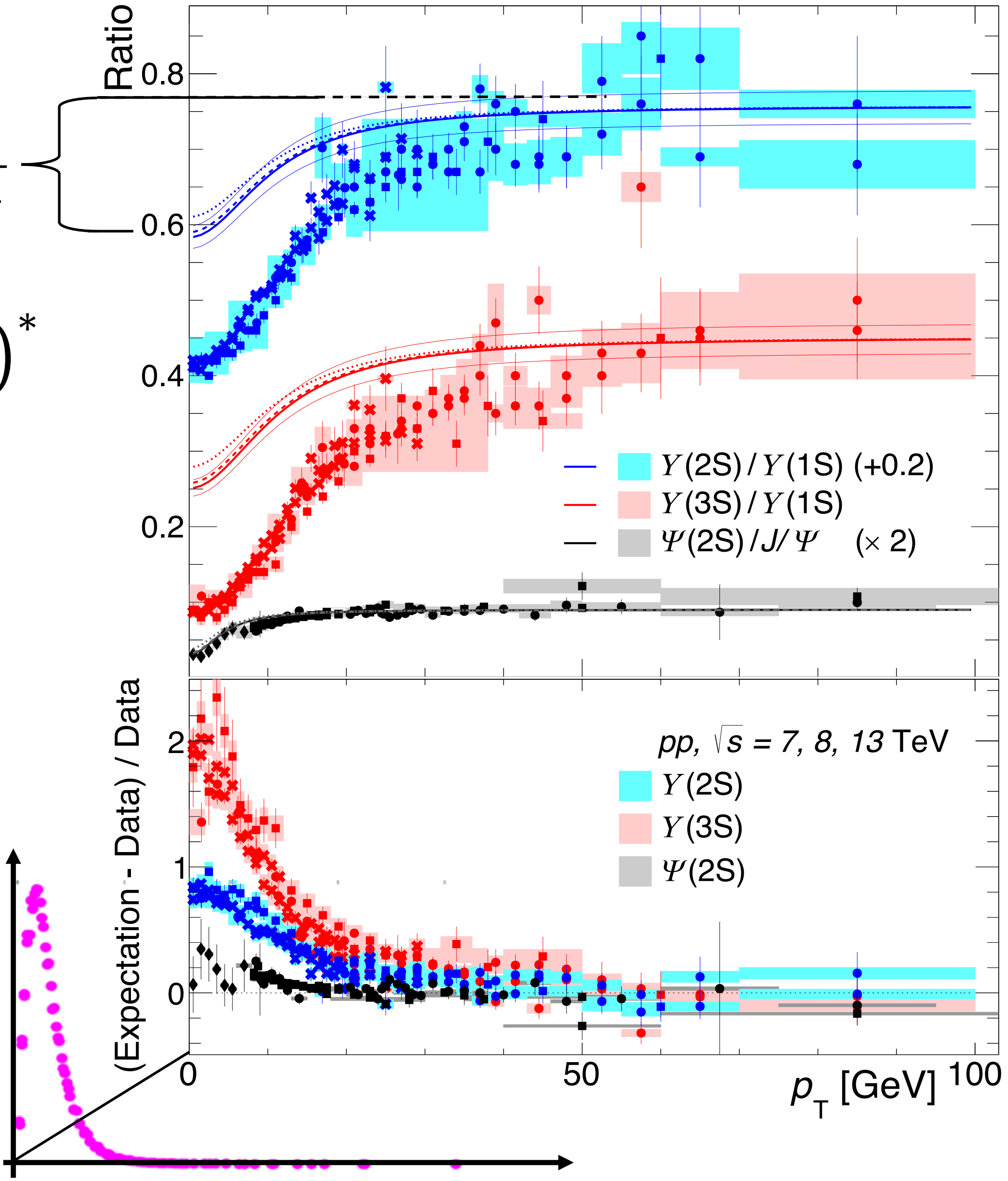
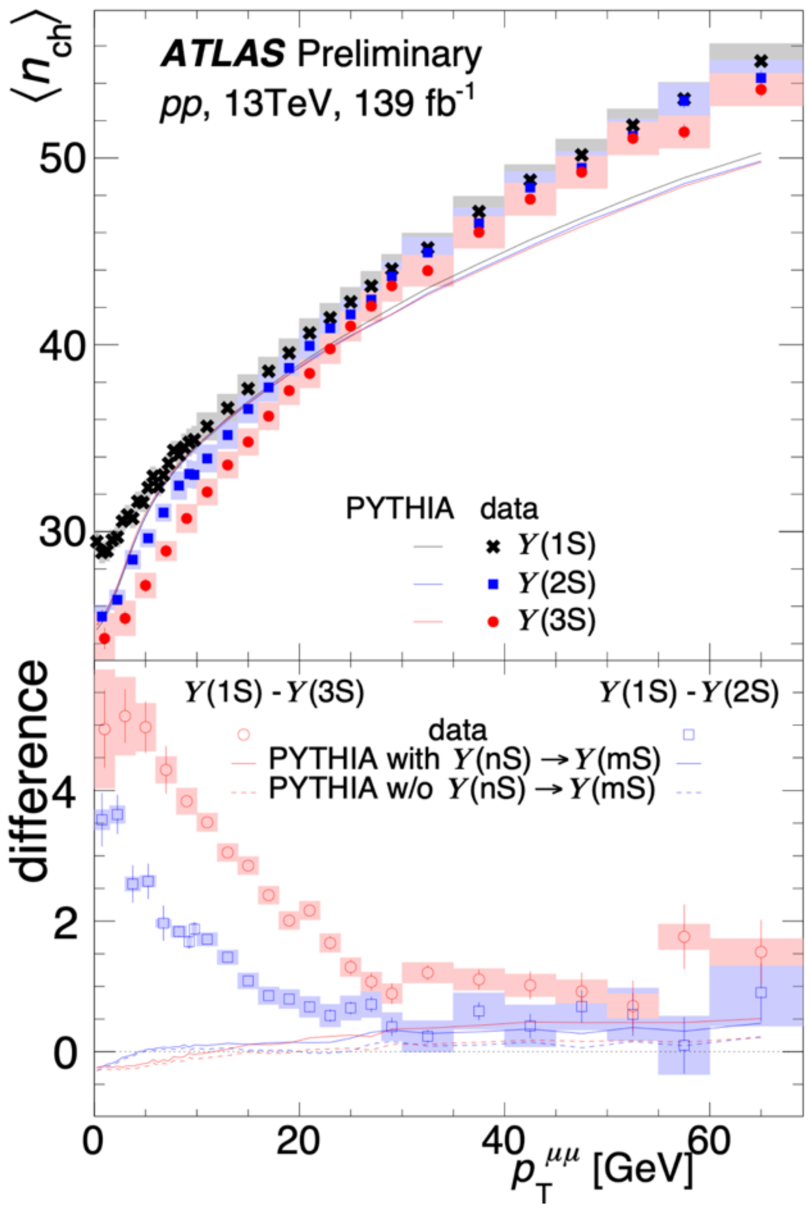
Multiplying by experimental spectra

- $\Upsilon(2S)$ should be factor 1.6 larger!
- $\Upsilon(3S)$ factor 2.4!



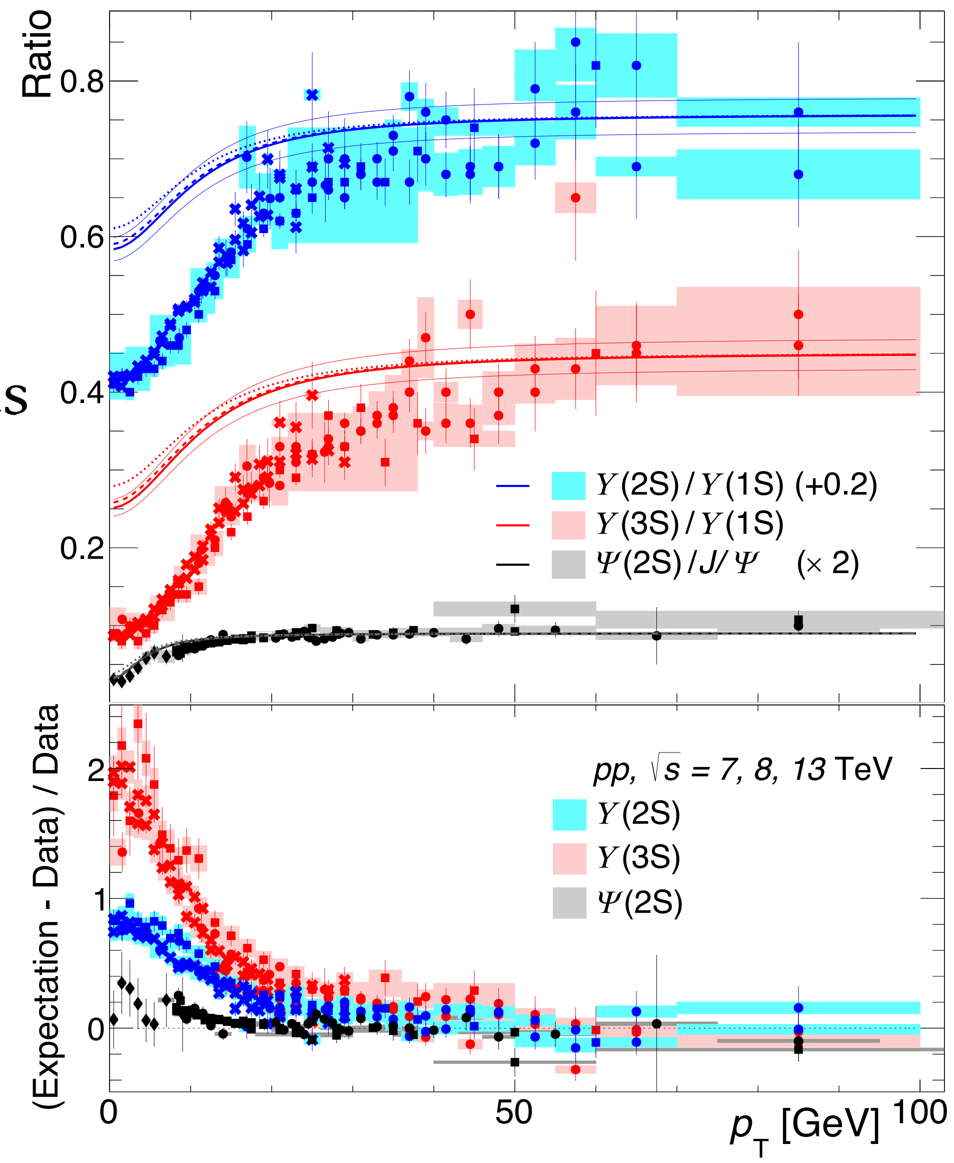
Bringing pieces together



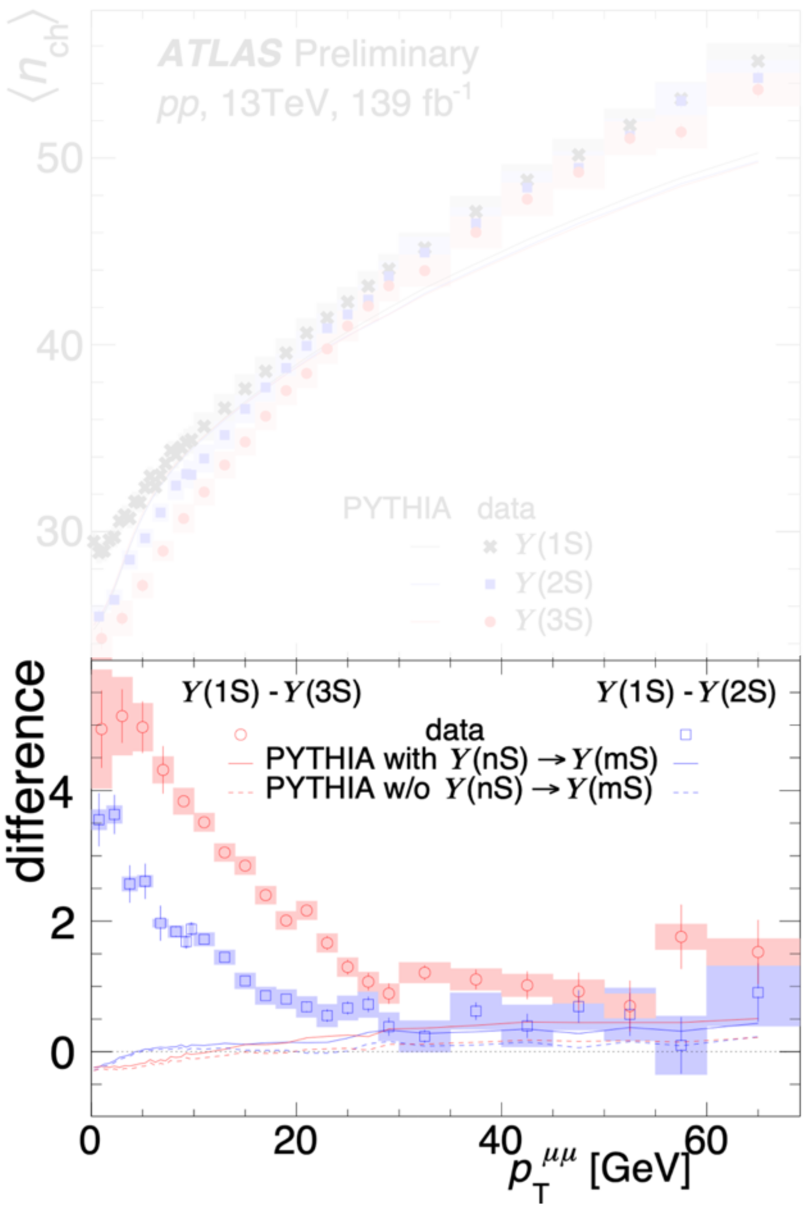
Independent analyses by CMS and ATLAS

Link the $\Upsilon(nS)$ production to the UE
-- ATLAS by kinematics
-- CMS by sphericity

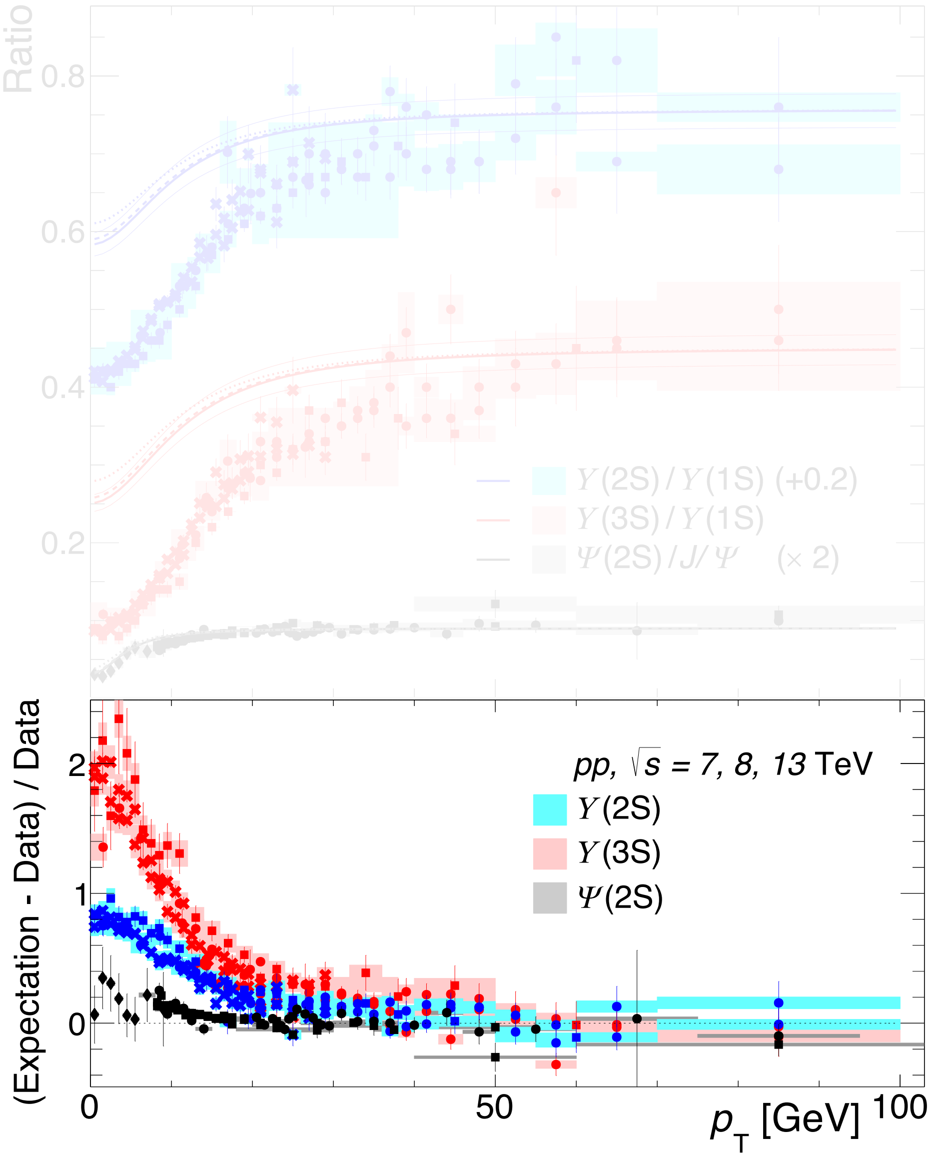
Deficit of the excited $\Upsilon(nS)$ with similar
-- p_T dependence
-- specie ordering



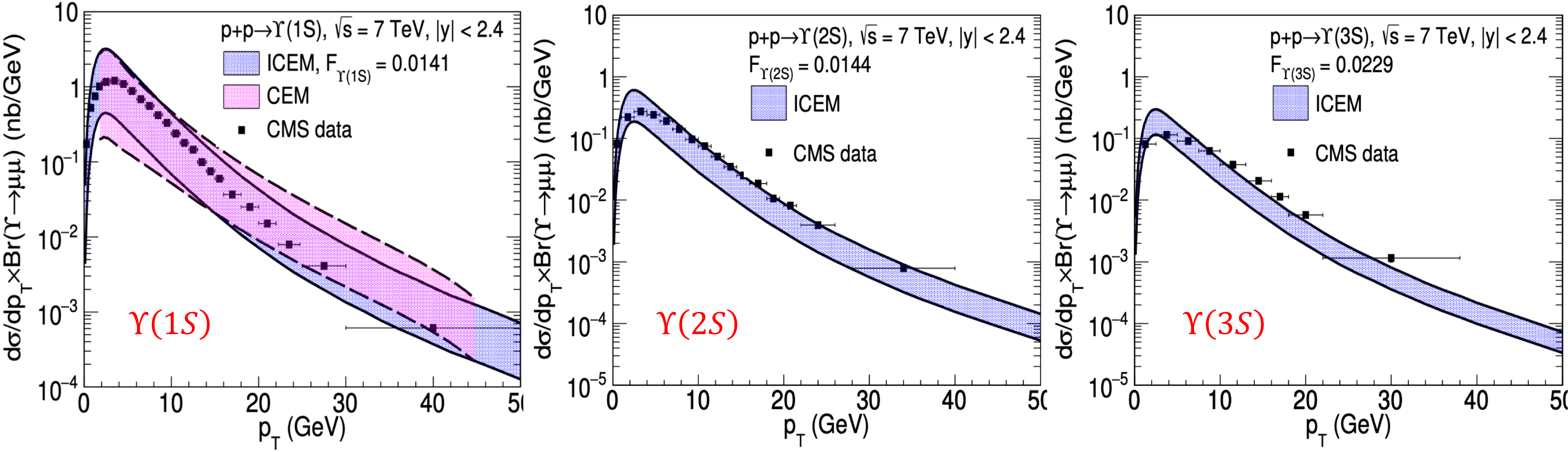
Final state interaction in *pp*

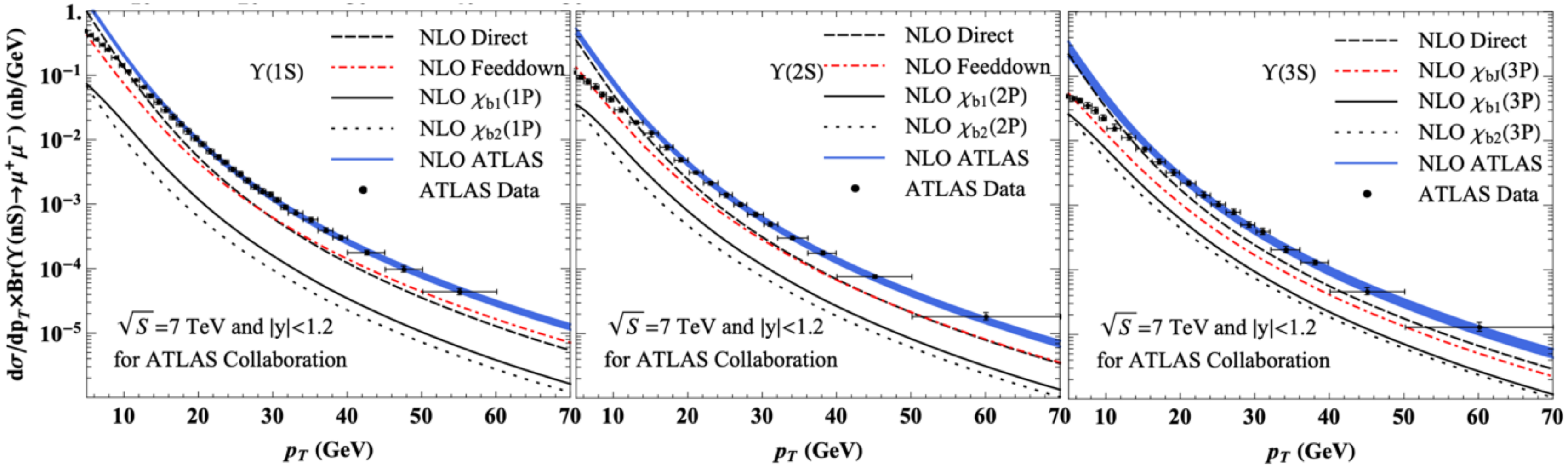


$\Upsilon(nS)$ production in *pp* is suppressed by the UE



Cross-section calculations





PRD94, 014028 (2016)
PPNP 131, (2023), 104044
arXiv:2305.13177

What about charmonia?

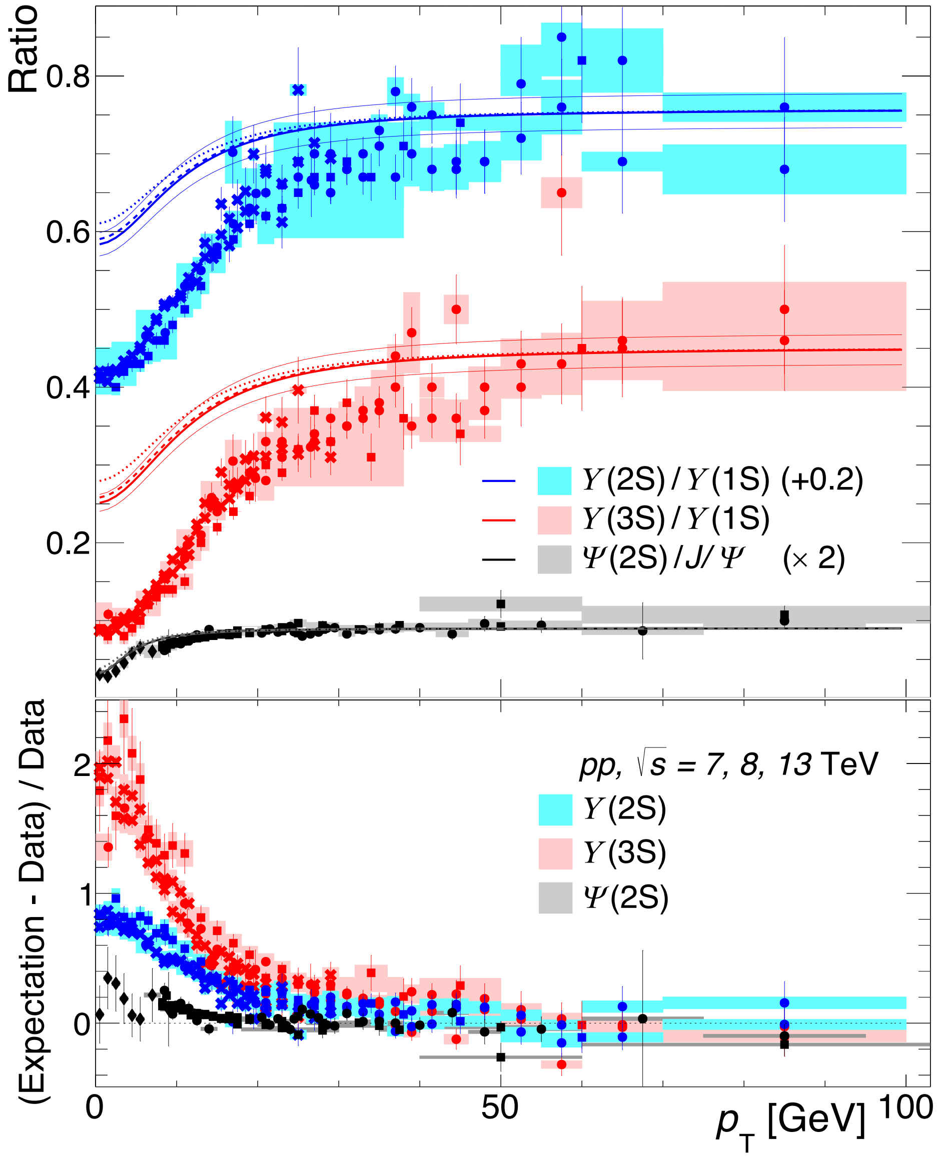
It would be logical to assume that the effect is related to the $q\bar{q}$ binging energy, but then $\psi(2S)$ must show a lot more suppression.

n_{ch} for $\psi(2S)$ shall be measured

EPJC (2018) 78:731

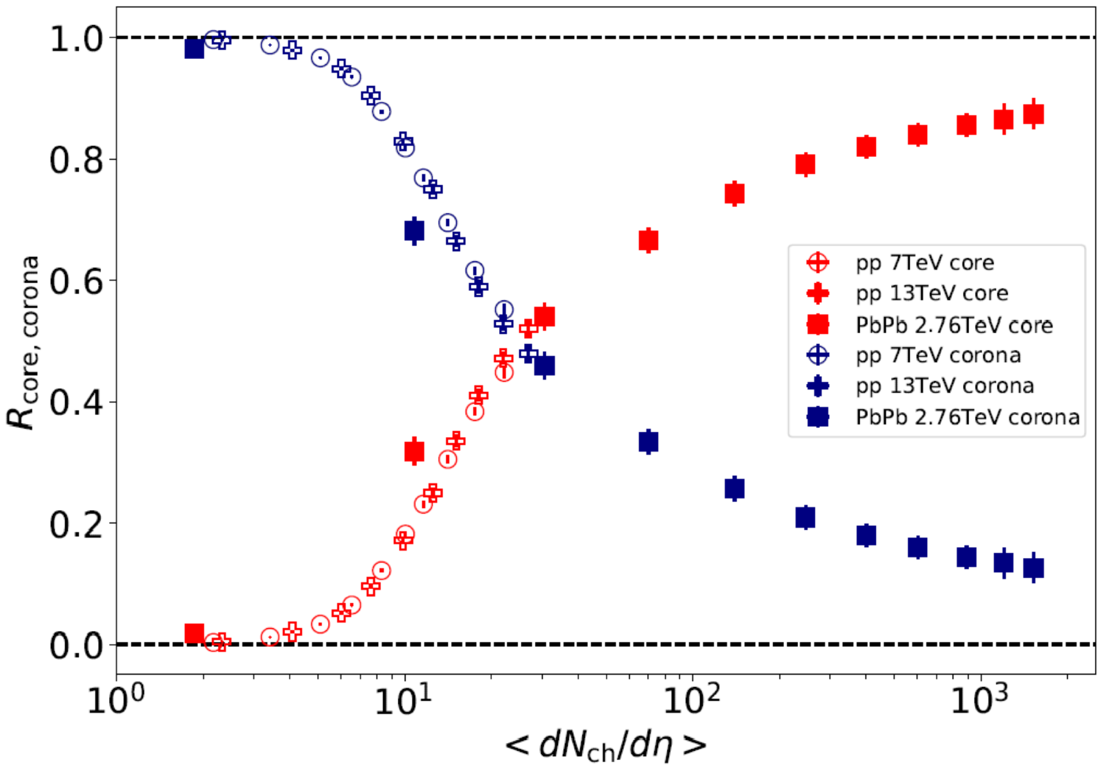
Table 1 Binding energies of the quarkonia shown in Fig. 3

Quarkonium	E_b (MeV)	Quarkonium	E_b (MeV)
$\chi_{b2}(3P)$	36	χ_{c0}	315
$\psi(2S)$	44	$\chi_{b0}(3P)$	326
$\chi_{b1}(3P)$	47	$\Upsilon(2S)$	536
$\chi_{b0}(3P)$	62	J/ψ	633
χ_{c2}	174	$\chi_{b2}(1P)$	647
$\Upsilon(3S)$	204	$\chi_{b1}(1P)$	666
χ_{c1}	219	$\chi_{b0}(1P)$	700
$\chi_{b2}(2P)$	290	$\Upsilon(1S)$	1099
$\chi_{b1}(3P)$	304		

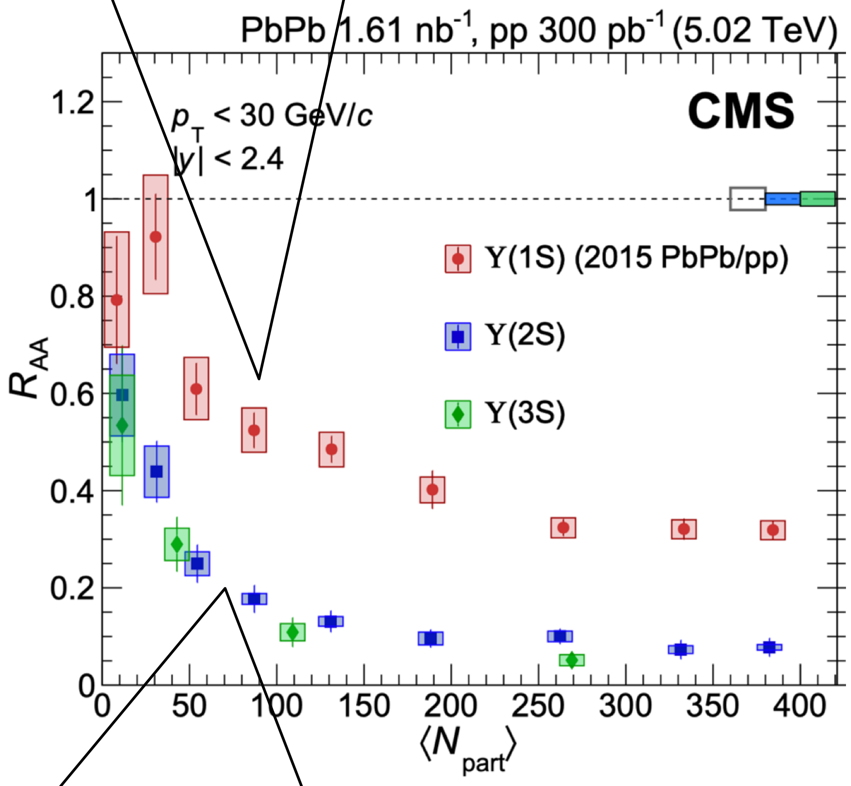


How it can look like in larger systems

from Yuuka Kanakubo's talk on Monday

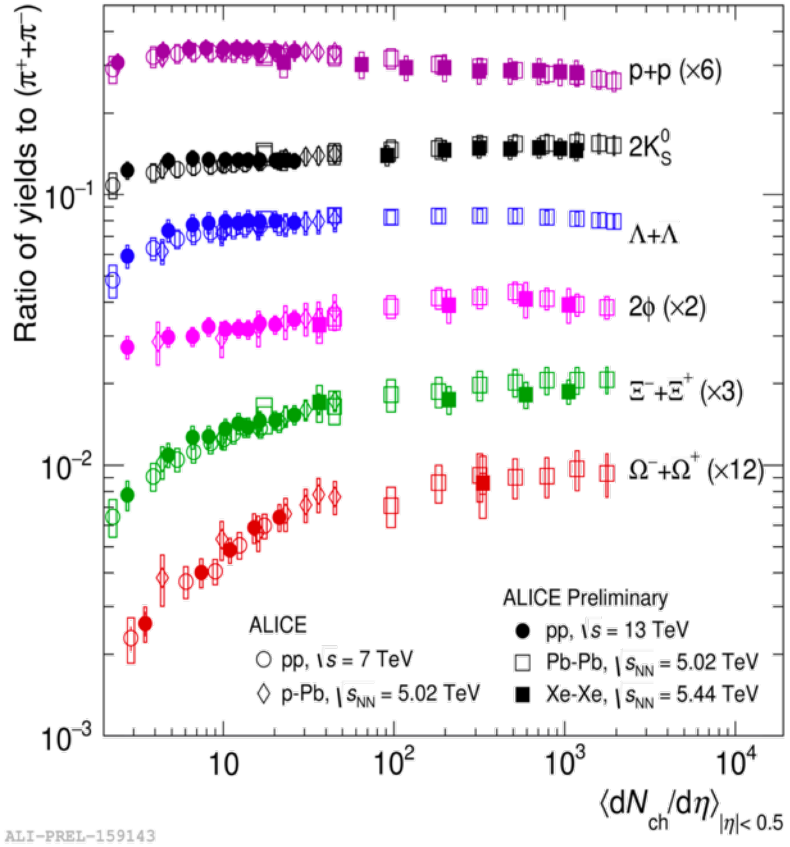


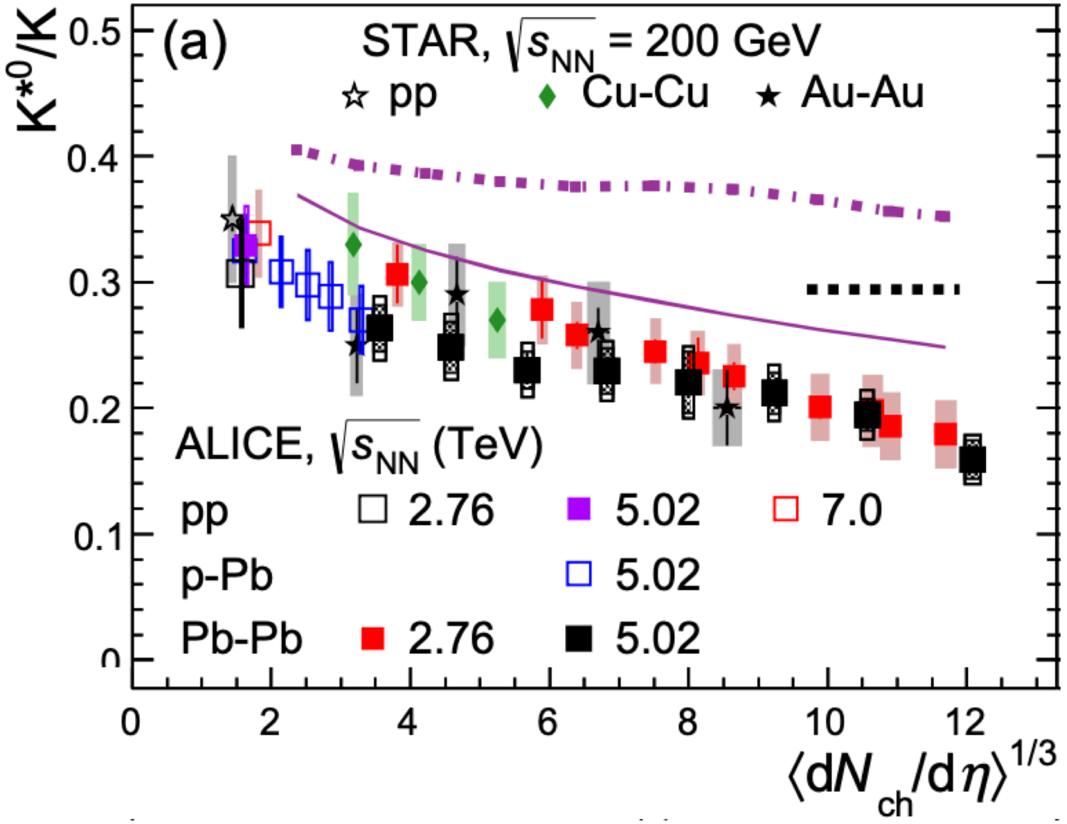
Core + corona: $\Upsilon(1S)$ resembles other particles, or we can't say better



Pure corona: medium is nearly opaque to $\Upsilon(2S)$ and $\Upsilon(3S)$) even in *pp*

QGP signatures in small systems





All strange-to-non-strange particle ratios go up

And K*/K ratio goes down…

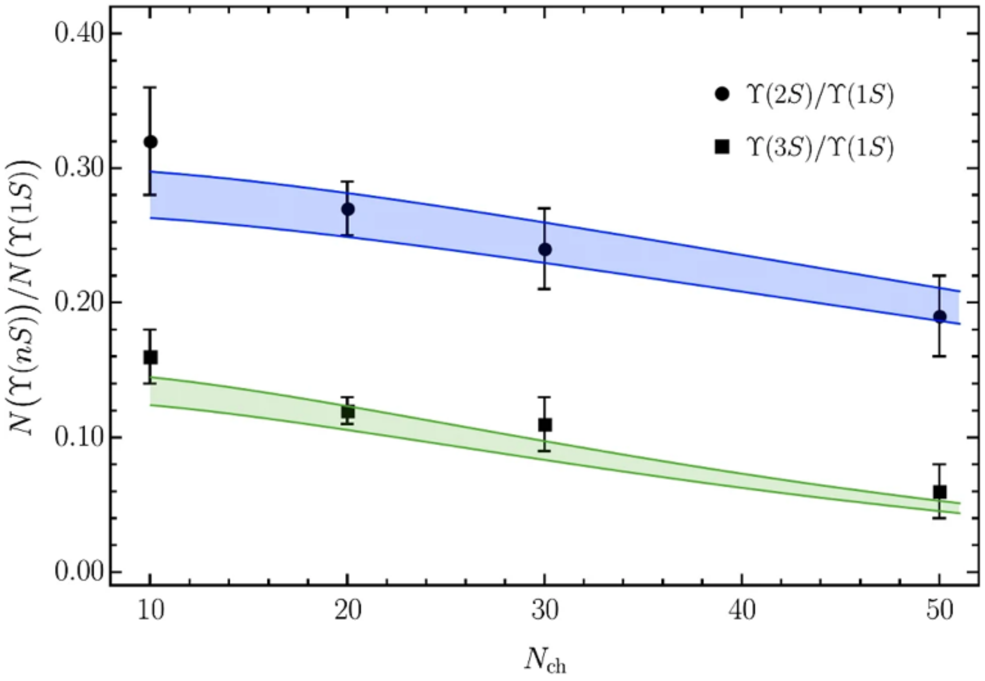
It might be that the effect is wider than just quarkonia

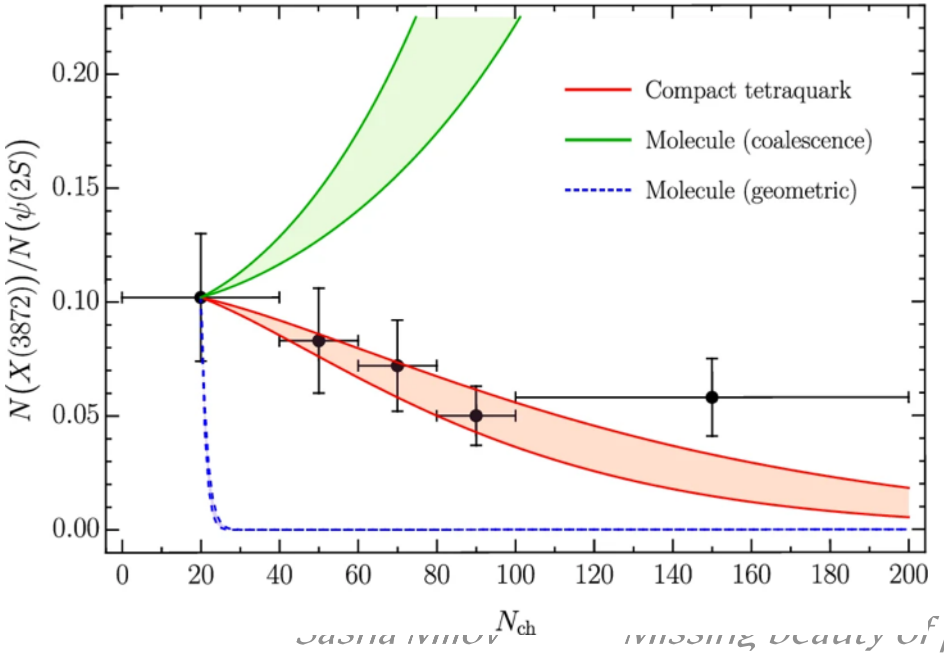
Comover interaction model

EPJC 81, 669 (2021)

Within CIM, quarkonia are broken by collisions with comovers – i.e. final state particles with similar rapidities.

CIM is typically used to explain *p*+A and A+A systems, although recently it was successfully applied to *pp*.





It looks like the effect isn't limited to only $\Upsilon(nS)$, at least χ_c can be affected as well, and possibly $\Psi(2S)$

In summary

Excited $\Upsilon(nS)$ states are destroyed in *pp* collisions by interactions with the UE

Only ~60% of $\Upsilon(2S)$ and only ~40% of $\Upsilon(3S)$ get out of the pp collisions at the LHC energies, based on what should be there from measured $\Upsilon(1S)$

Other particles can be affected as well. Some indirect hints exist for $\Upsilon(1S)$, $\Psi(2S)$, and even K*

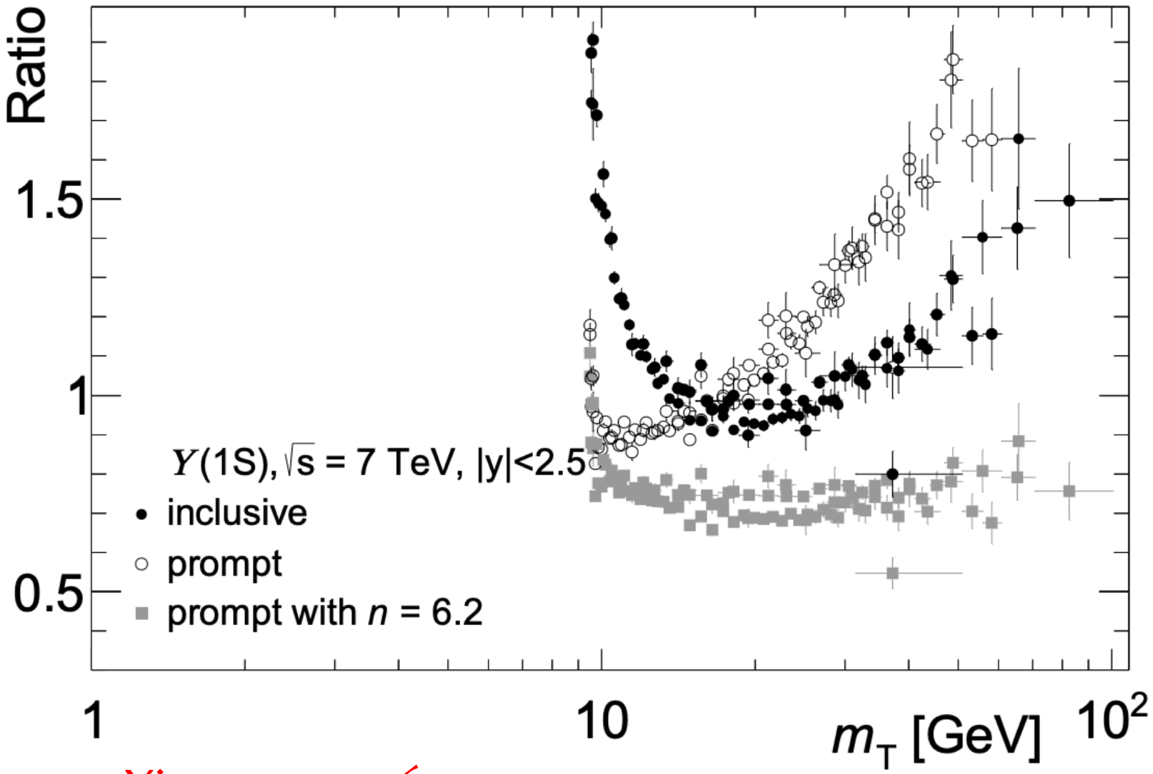
At the moment we do not know much about the observed phenomenon, but many signatures can be measured and not only at the LHC

Comover model explains one curve. More theoretical guidance is badly needed!

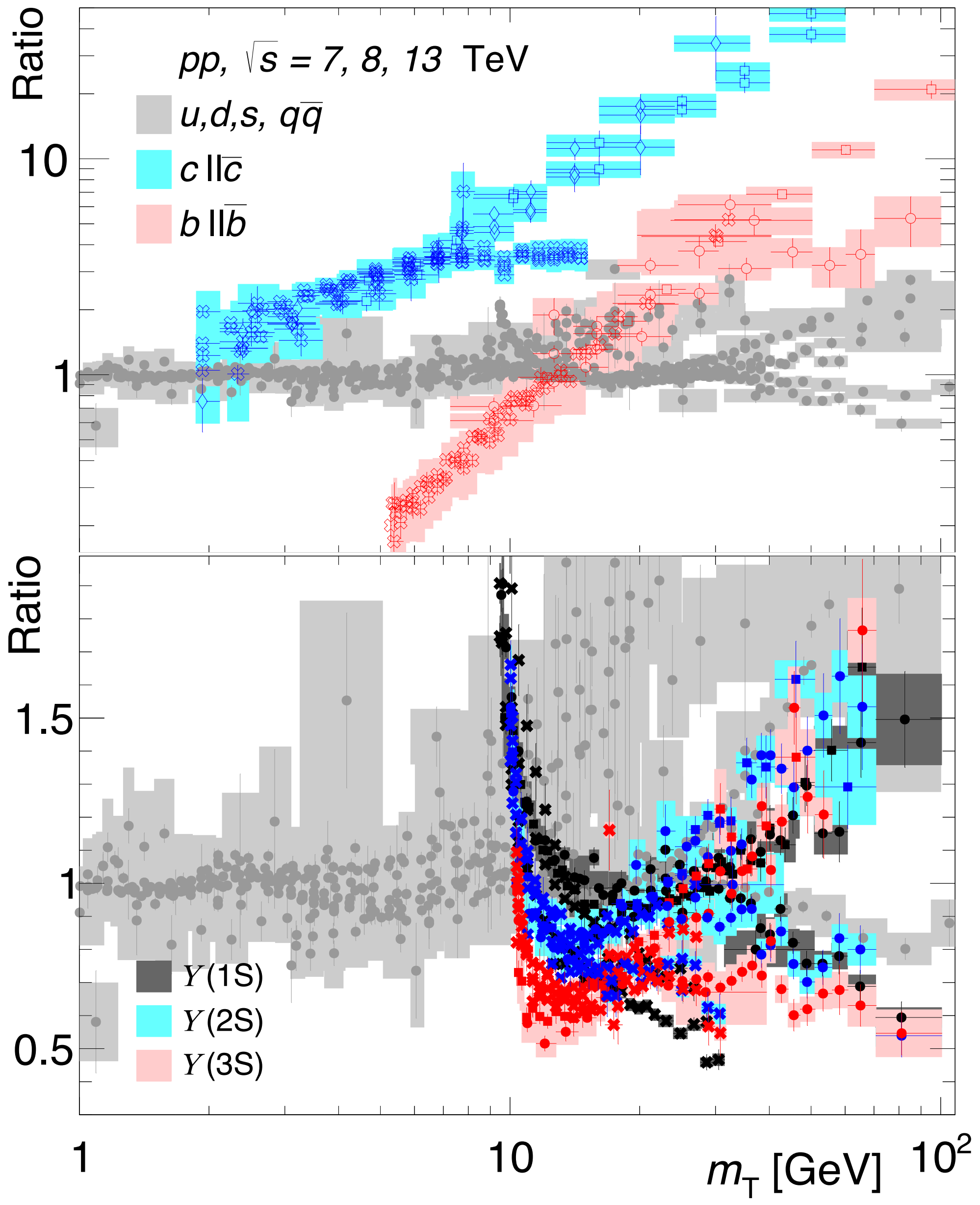
backups

Common fit

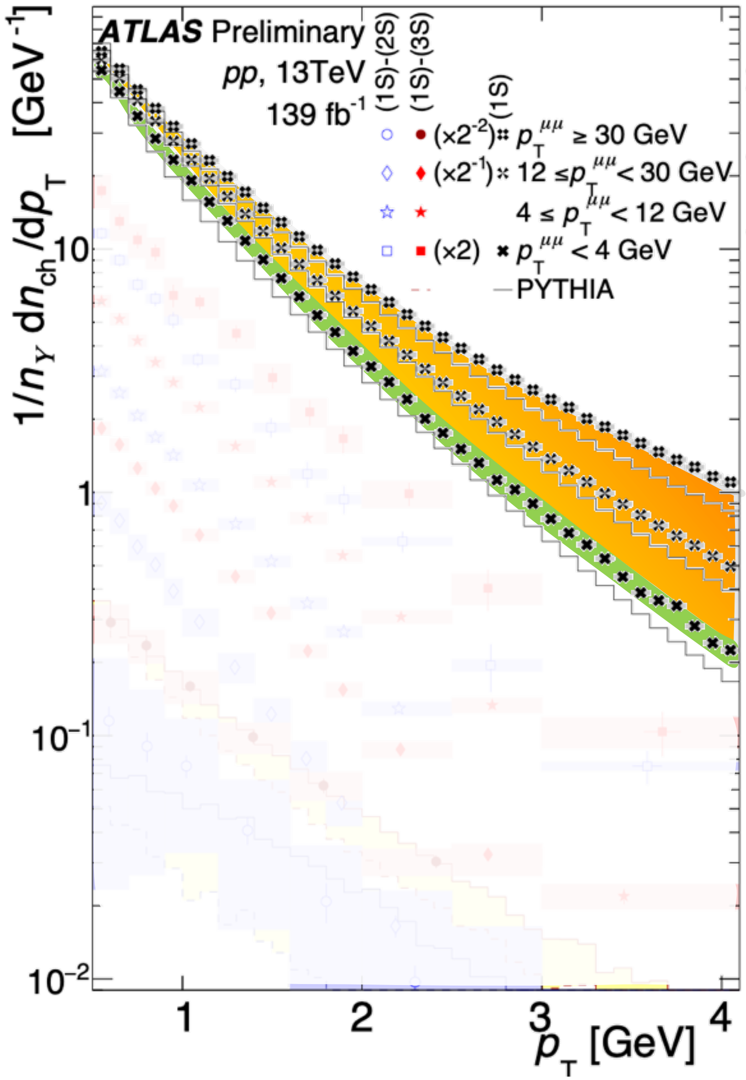
$$\lim_{\Delta m, p_T \ll m_{q\bar{q}}} \left[\frac{nT + \sqrt{p_T^2 + (m_{q\bar{q}} + \Delta m)^2}}{nT + \sqrt{p_T^2 + m_{q\bar{q}}^2}} \right]^{-n} = 1 - \frac{n\Delta m}{nT + m_{q\bar{q}}}$$

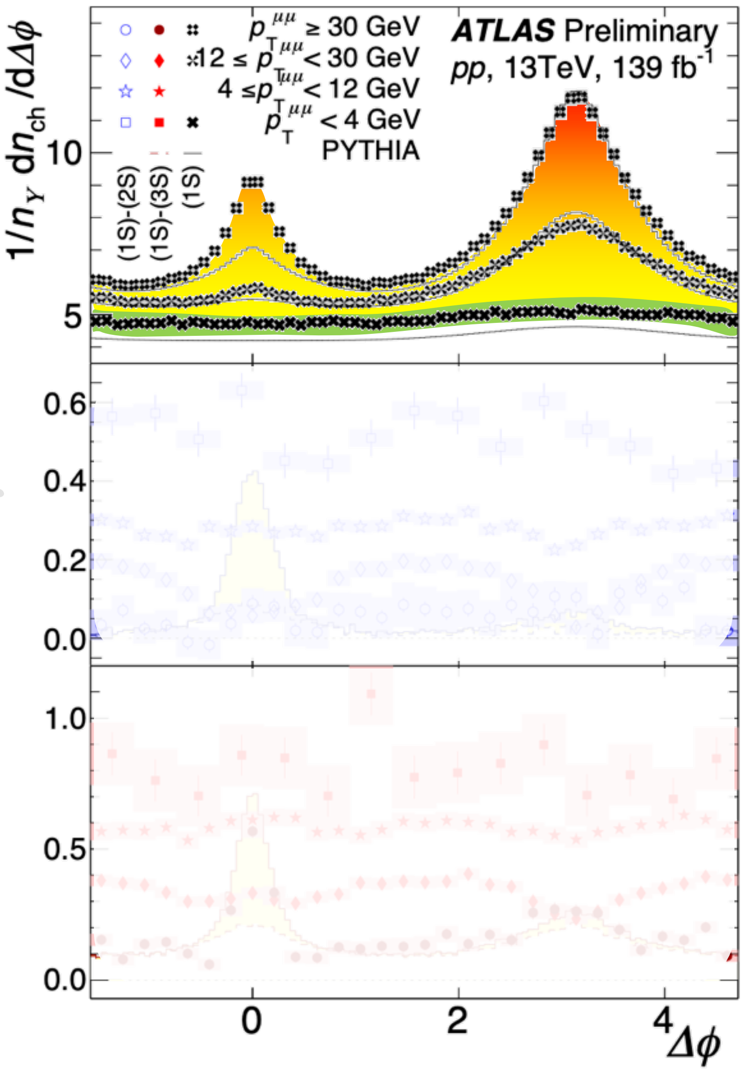


arXiv:2210.11026v1



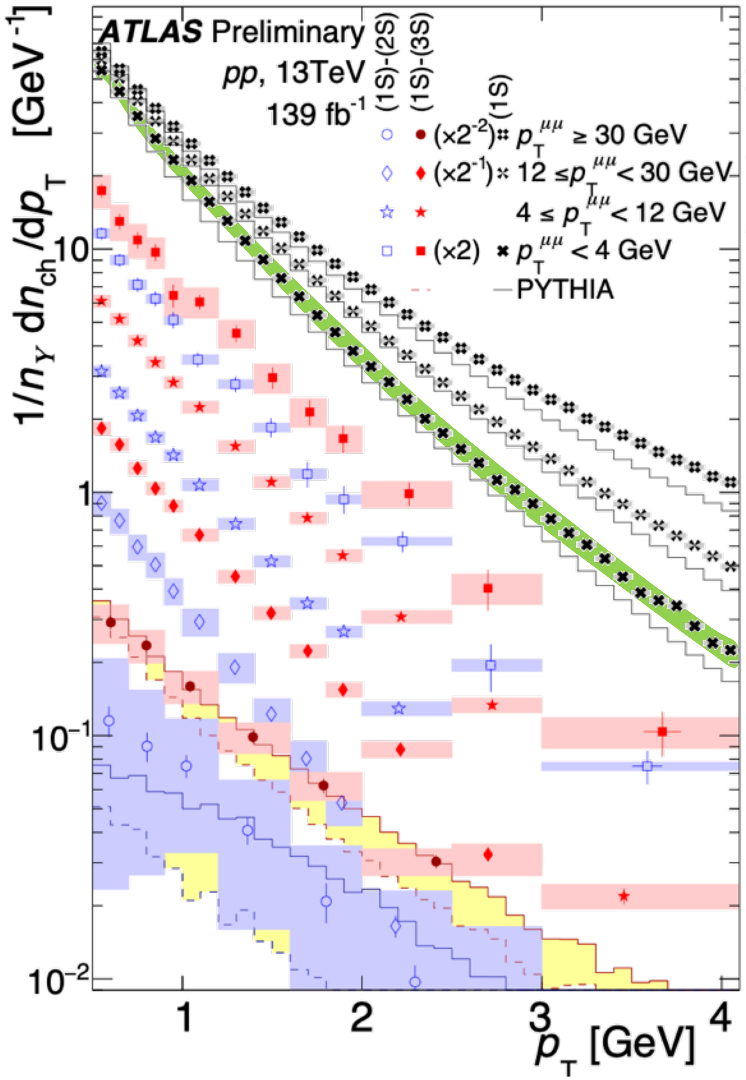
Kinematic distributions

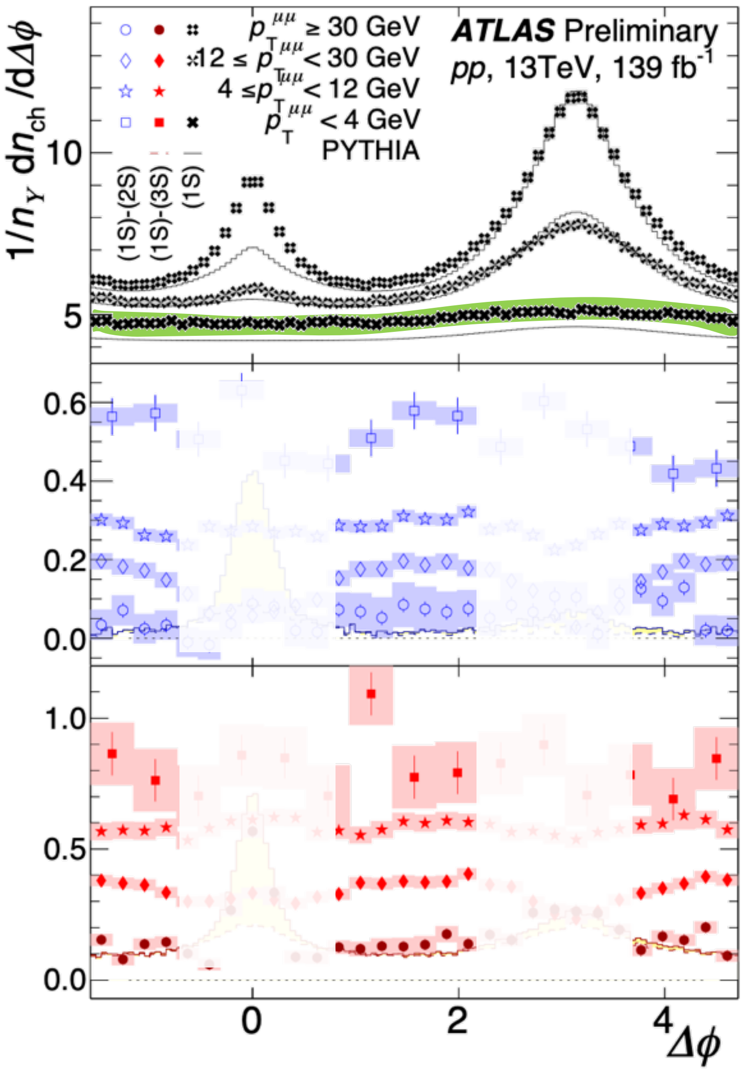




- Distributions for $\Upsilon(1S)$
- Pythia does not describe well
- One cannot measure the UE, but $p_T < 4$ GeV is the closest to it, jet part that is correlated to $\Upsilon(nS)$

Kinematic distributions





- Distributions for $\Upsilon(1S)$
- Pythia does not describe well
- One cannot measure the UE, but $p_T < 4$ GeV is the closest to it, jet part that is correlated to $\Upsilon(nS)$
- Subtracted distributions look like UE at rather high $\Upsilon(nS)$ p_T. At the highest p_T there are feed-downs
- Away from jets there are regions with charged particles

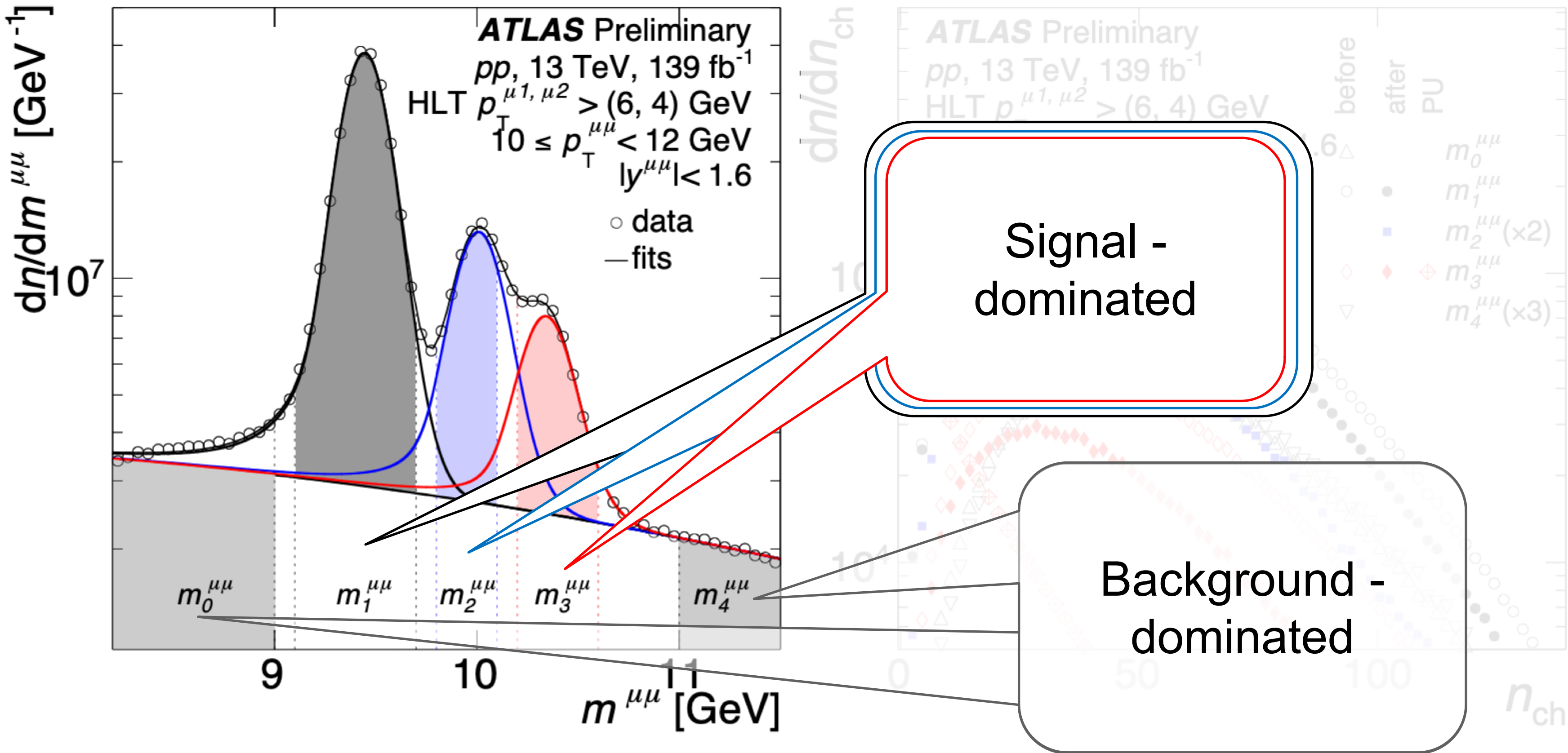
Signal extraction

$$\text{fit}(m) = \sum_{nS} N_{\Upsilon(nS)} F_n(m) + N_{\text{bkg}} F_{\text{bkg}}(m)$$

$$F_n(m) = (1-\omega_n) CB_n(m) + \omega_n G_n(m)$$

$$F_{\text{bkg}}(m) = \sum_{i=0}^{3} a_i (m - m_0)^i ; a_0 = 1$$

- Define 3+2 regions

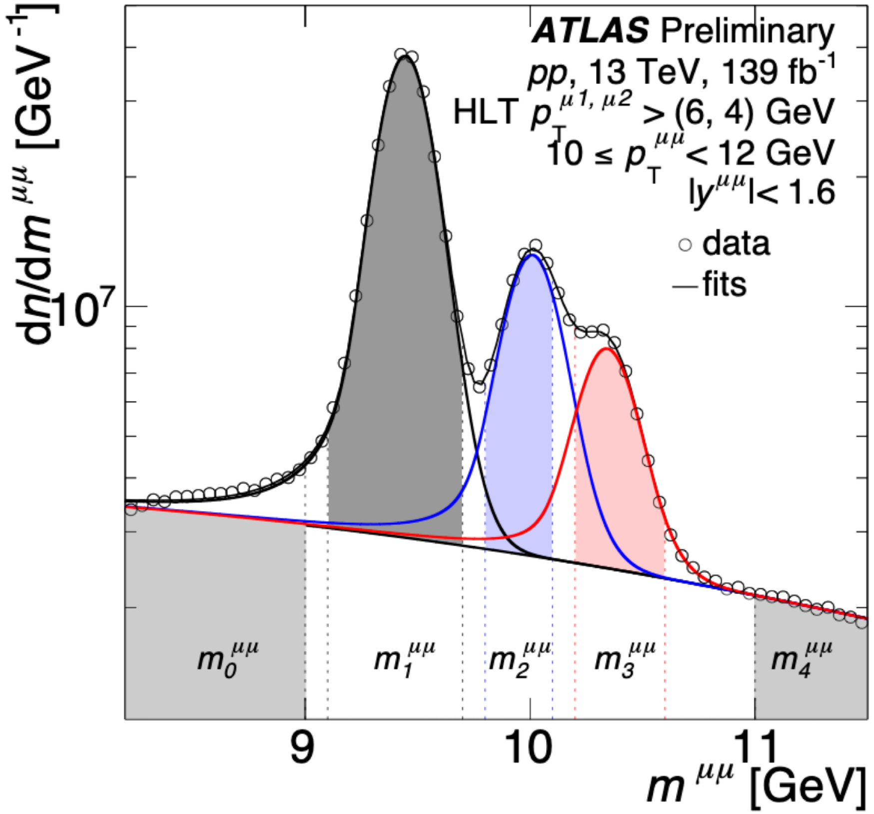


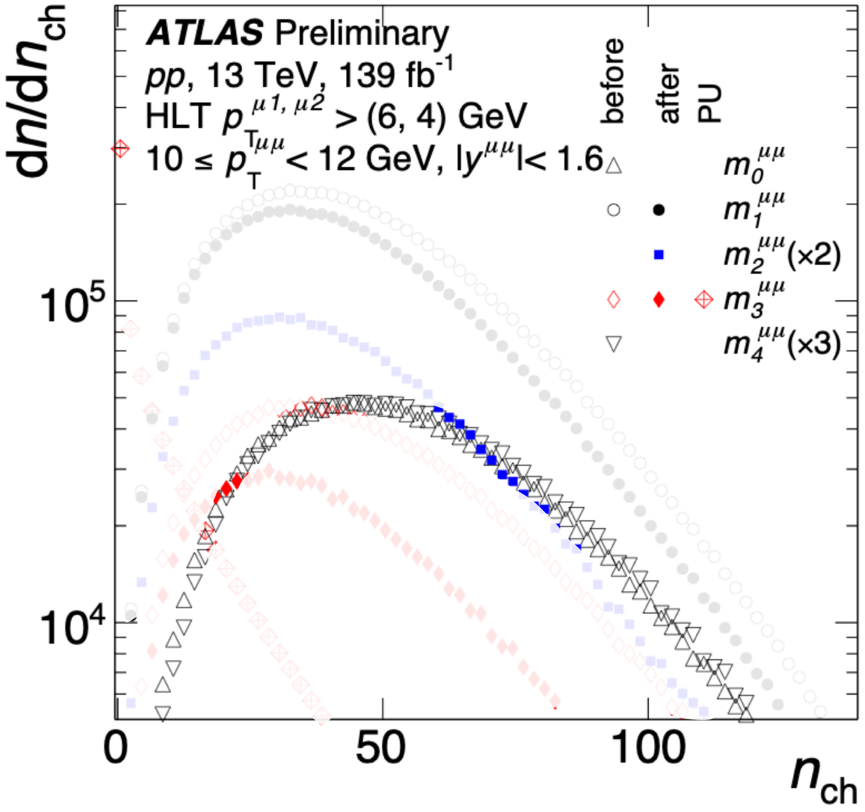
Signal extraction

$$s_n = \frac{\int_{m_n^{\mu\mu}} N_{\Upsilon(nS)} F_n(m)\, dm}{\int_{m_n^{\mu\mu}} \mathrm{fit}(m)\, dm}$$

$$f_{nk} = \frac{\int_{m_n^{\mu\mu}} N_{\Upsilon(kS)} F_k(m)\, dm}{\int_{m_n^{\mu\mu}} \mathrm{fit}(m)\, dm} \qquad k_n = \frac{\langle F_{\mathrm{bkg}}(m)\rangle|_{m_4^{\mu\mu}} - \langle F_{\mathrm{bkg}}(m)\rangle|_{m_n^{\mu\mu}}}{\langle F_{\mathrm{bkg}}(m)\rangle|_{m_4^{\mu\mu}} - \langle F_{\mathrm{bkg}}(m)\rangle|_{m_0^{\mu\mu}}}$$

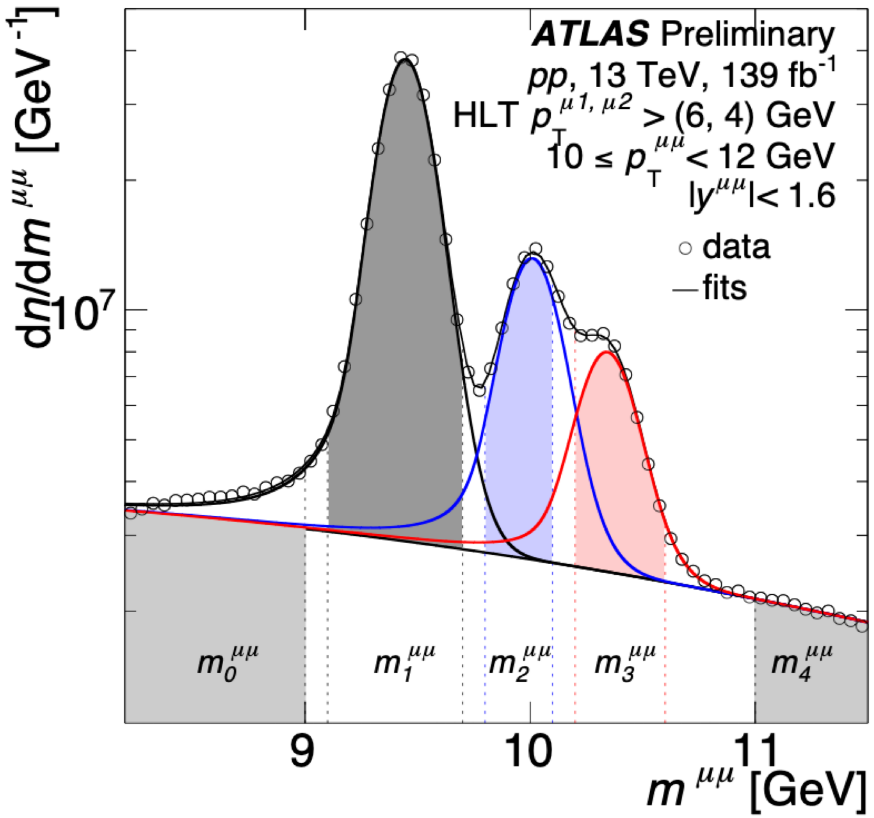
- Define 3+2 regions
- Bkg shapes are similar – interpolate

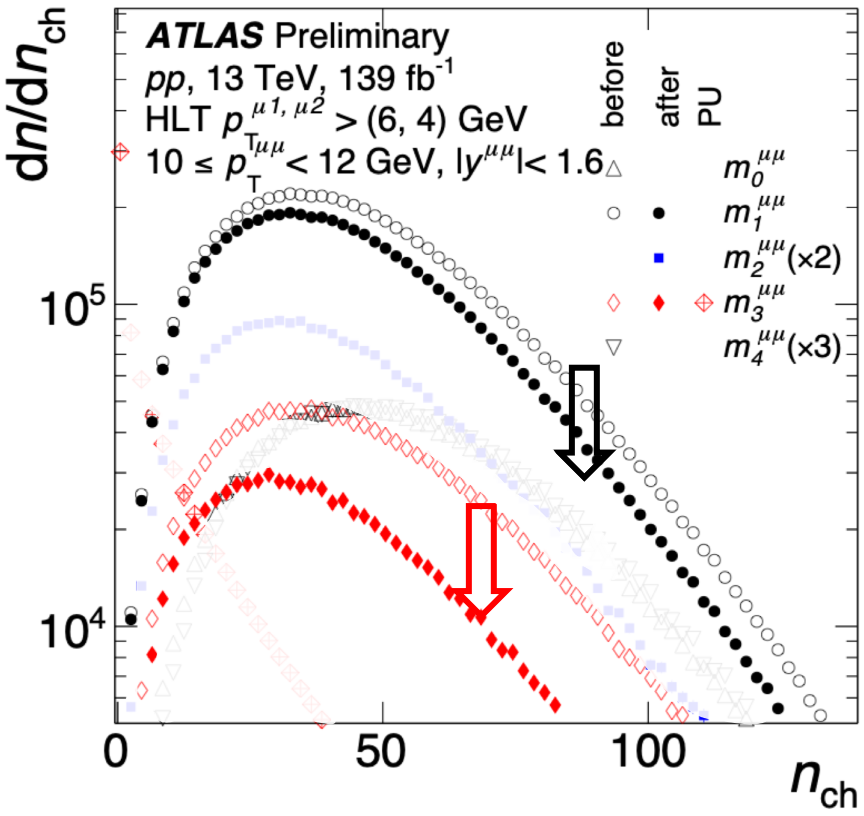




Signal extraction

$$\begin{pmatrix} P(m_0^{\mu\mu}) \\ P(m_1^{\mu\mu}) \\ P(m_2^{\mu\mu}) \\ P(m_3^{\mu\mu}) \\ P(m_4^{\mu\mu}) \end{pmatrix} = \begin{pmatrix} 1-f_{01} & f_{01} & 0 & 0 & 0 \\ k_1\,(1-s_1) & s_1 & 0 & 0 & (1-k_1)\,(1-s_1) \\ k_2\,(1-s_2-f_{21}-f_{23}) & f_{21} & s_2 & f_{23} & (1-k_2)\,(1-s_2-f_{21}-f_{23}) \\ k_3\,(1-s_3-f_{32}) & 0 & f_{32} & s_3 & (1-k_3)\,(1-s_3-f_{32}) \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} P_0 \\ P(\Upsilon(1S)) \\ P(\Upsilon(2S)) \\ P(\Upsilon(3S)) \\ P_4 \end{pmatrix}$$

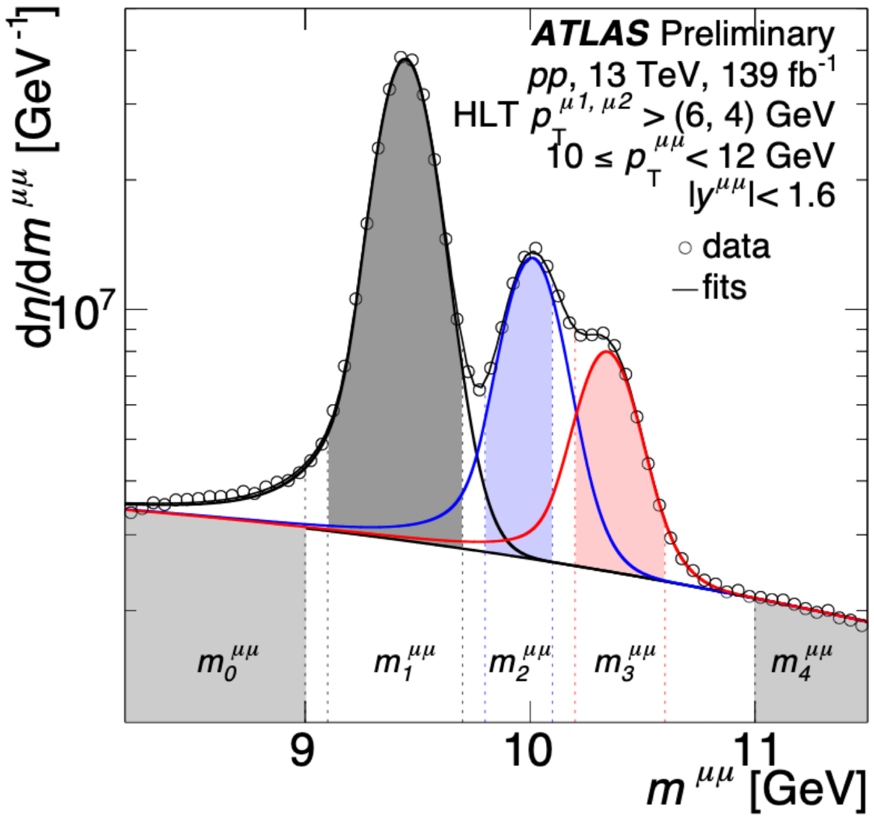


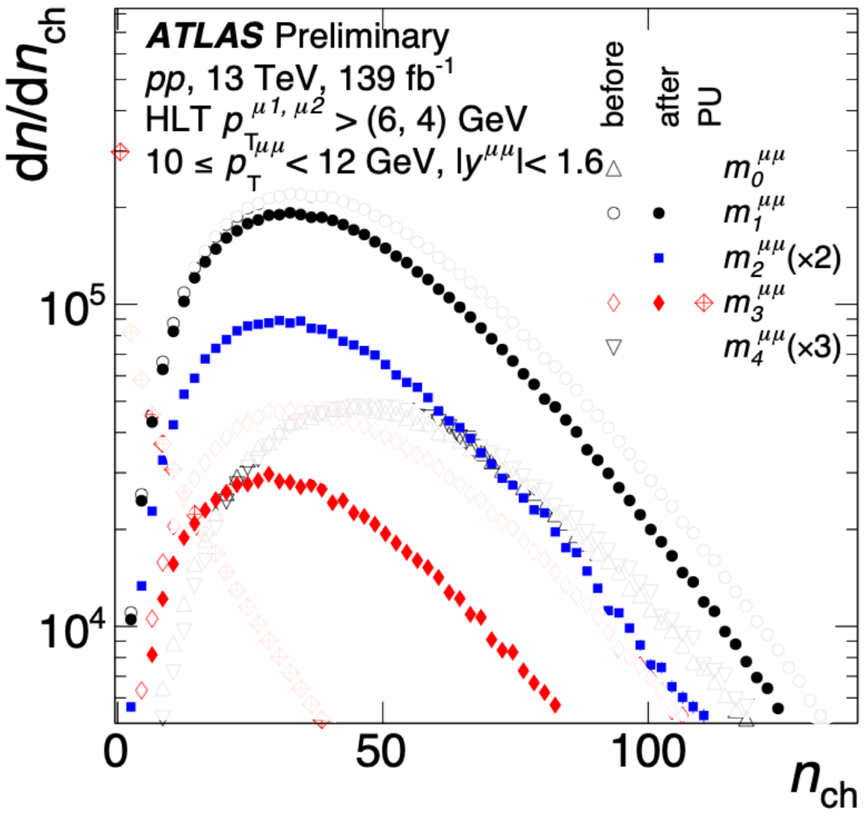


- Define 3+2 regions
- Bkg shapes are similar – interpolate
- Bkg subtraction for $\Upsilon(1S)$ and $\Upsilon(3S)$

Signal extraction

$$\begin{pmatrix} P(m_0^{\mu\mu}) \\ P(m_1^{\mu\mu}) \\ P(m_2^{\mu\mu}) \\ P(m_3^{\mu\mu}) \\ P(m_4^{\mu\mu}) \end{pmatrix} = \begin{pmatrix} 1-f_{01} & f_{01} & 0 & 0 & 0 \\ k_1(1-s_1) & s_1 & 0 & 0 & (1-k_1)(1-s_1) \\ k_2(1-s_2-f_{21}-f_{23}) & f_{21} & s_2 & f_{23} & (1-k_2)(1-s_2-f_{21}-f_{23}) \\ k_3(1-s_3-f_{32}) & 0 & f_{32} & s_3 & (1-k_3)(1-s_3-f_{32}) \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} P_0 \\ P(\Upsilon(1S)) \\ P(\Upsilon(2S)) \\ P(\Upsilon(3S)) \\ P_4 \end{pmatrix}$$

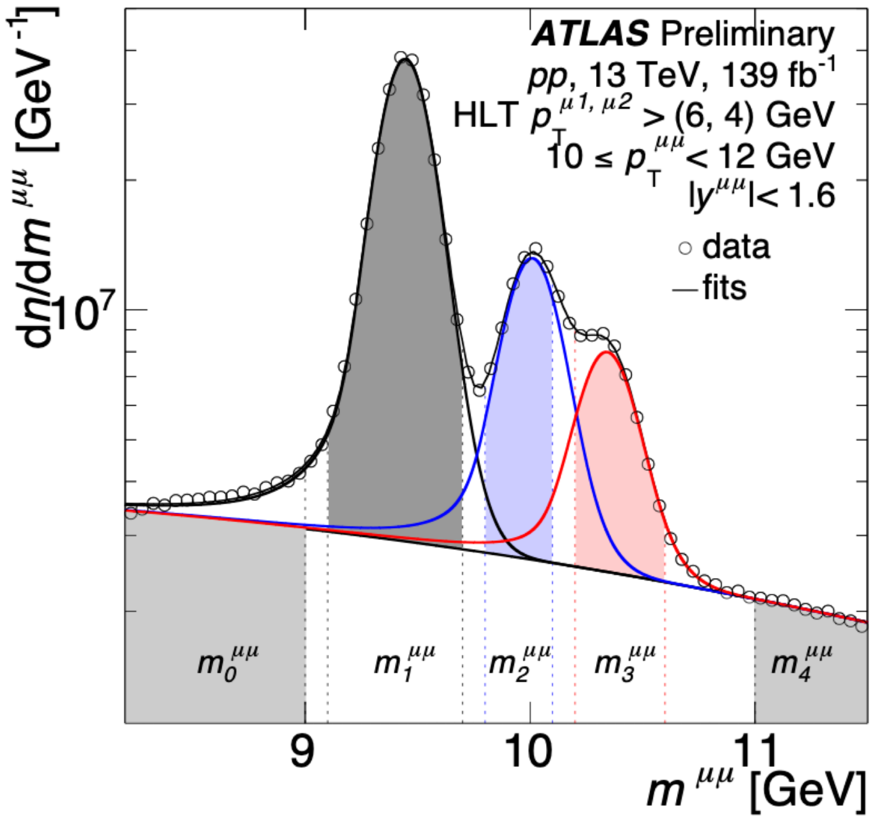


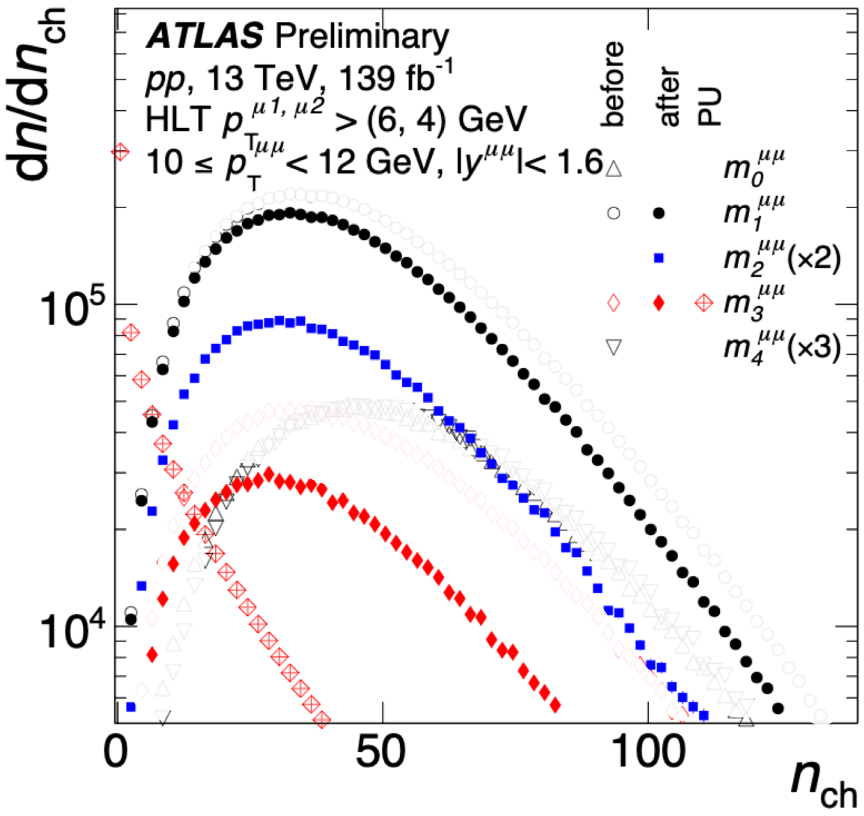


- Define 3+2 regions
- Bkg shapes are similar – interpolate
- Bkg subtraction for $\Upsilon(1S)$ and $\Upsilon(3S)$
- After subtraction n_{ch} look different

Signal extraction

Triggers are all combined together
Pileup is constructed from mixed events and is either directly subtracted or unfolded
Non-linear effects are also accounted for





- Define 3+2 regions
- Bkg shapes are similar – interpolate
- Bkg subtraction for $\Upsilon(1S)$ and $\Upsilon(3S)$
- After subtraction n_{ch} look different
- Remove pileup, same shape for all $\Upsilon(nS)$

The pileup story

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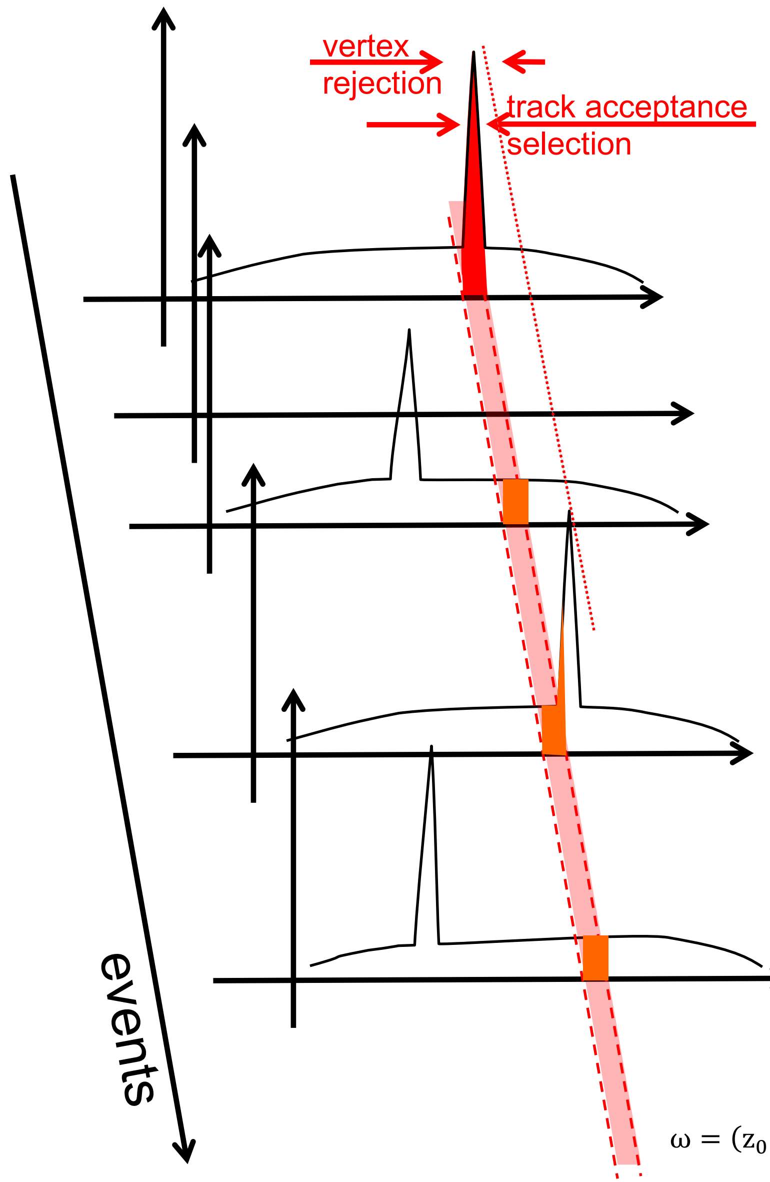


$\omega = (z_0 - z_{vtx}) \sin\theta$

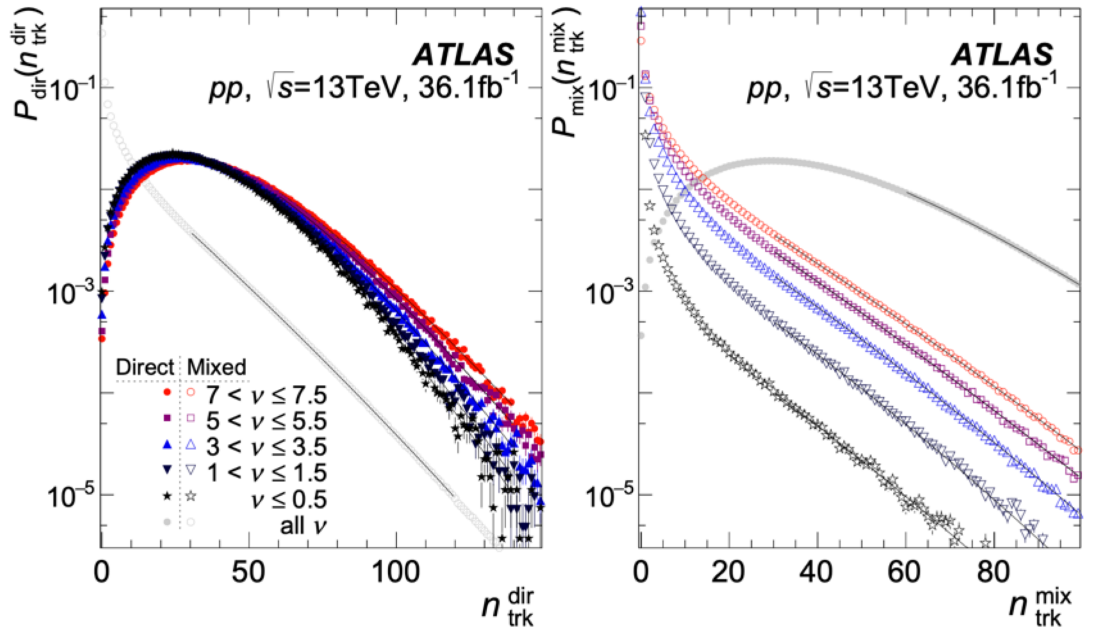
Start with the triggered event, called Direct

In the same run search for events with at the same μ

Build Mixed event from tracks with vertex pointing $|\omega| <$ 0.75 mm to the Direct event

If the other vertex is within 15mm of the Direct, discard it

Do 20 times to get statistics



$$\nu = 2\omega_0 \left. \frac{d^2 n_{trk}}{d\omega d\bar{\mu}} \right|_{\bar{z}_{vtx}=0} Gauss(\bar{z}_{vtx})\ \bar{\mu}$$

Track production (physics)

z_{vtx} distribution

Analysis selection

Instantaneous luminosity

Analysis in brief

Entire ATLAS Run-2 data: 2015 – 2018, $\sqrt{s}$ = 13 TeV, 139 fb^{-1}

Full luminosity data constrained at μ < 50 (fake production) and then at ν < 20 in 40 intervals

$\Upsilon(nS)$ are reconstructed as di-muons

6 different di-muon triggers with muon p_T from 4 to 11 GeV

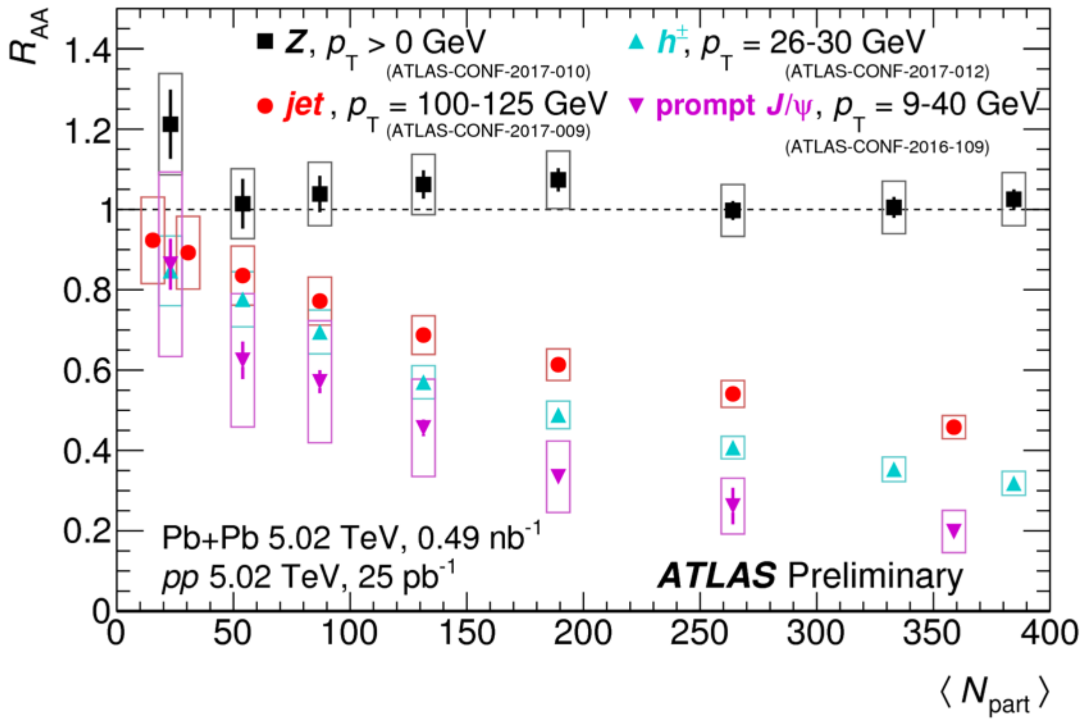
$\Upsilon(nS)$ kinematics |y| < 1.6, 0 < p_T < 70 GeV where we ran out of statistics

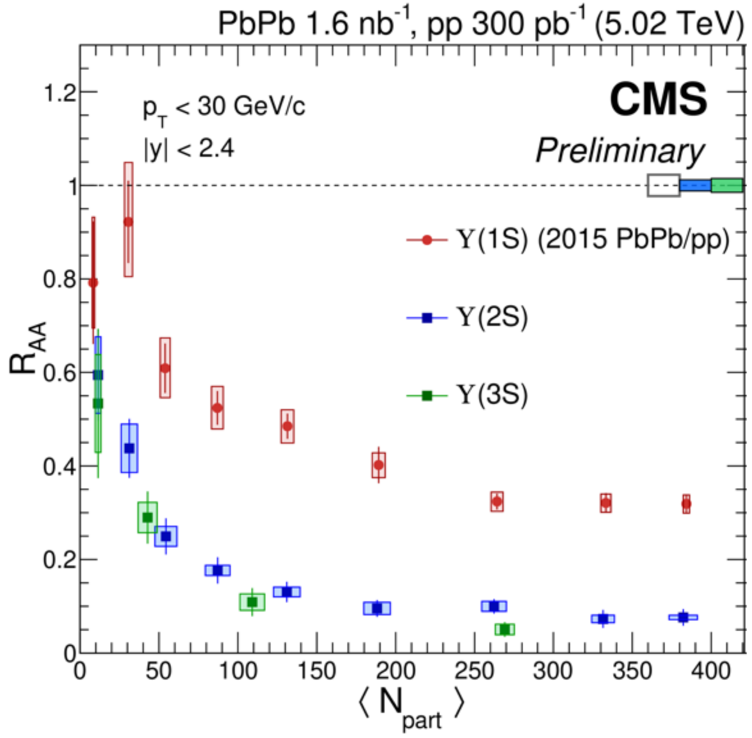
All together after cuts: $\sim 5\times 10^7 \Upsilon(1S)$, $\sim 10^7 \Upsilon(2S)$, $\sim 7\times 10^6 \Upsilon(3S)$

Charged hadrons kinematics $|\eta|$ < 2.5, 0.5 < p_T < 10 GeV, fully corrected

Dimuon invariant mass distributions are fitted to functions with 24 parameters

Back to heavy ions



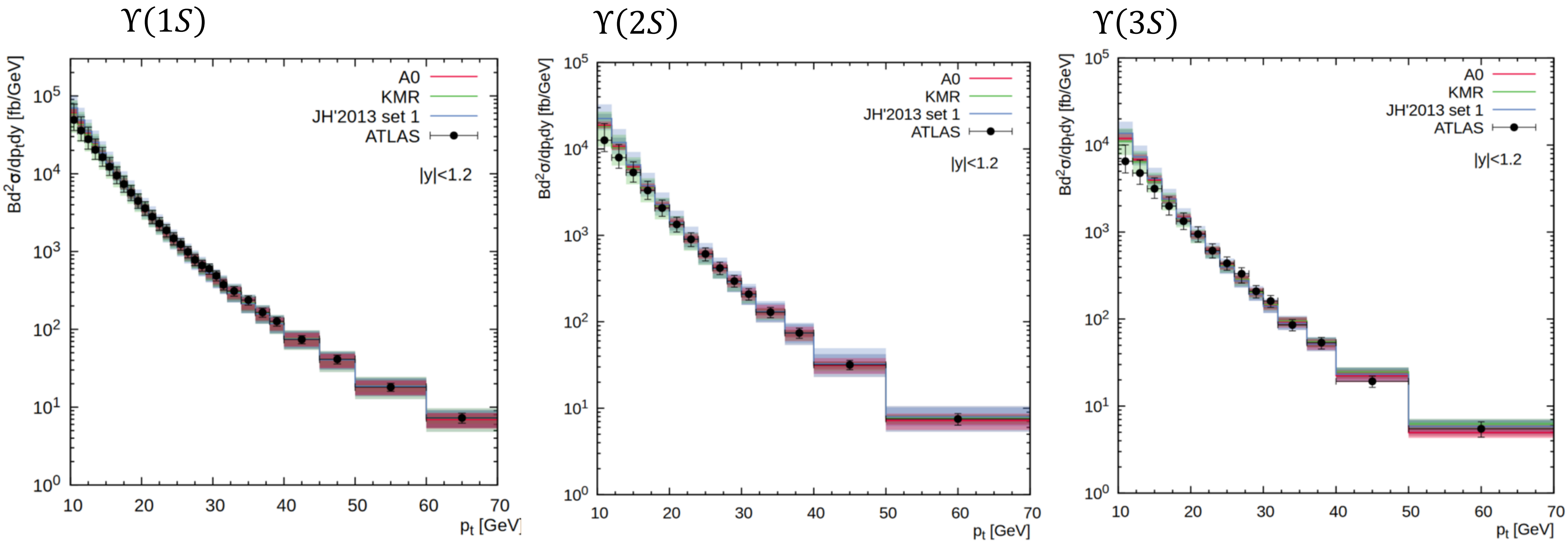


Similarity in the suppression of $\Upsilon(1S)$ and other species and the difference to higher $\Upsilon(nS)$ can be an indication of the regime change

Most particles, including $\Upsilon(1S)$ $L \geq \sqrt[3]{N_{\text{part}}} \times r_p$ volume emission

$\Upsilon(2S)$, $\Upsilon(3S)$ $L \ll \sqrt[3]{N_{\text{part}}} \times r_p$ surface emission

Theory calculation



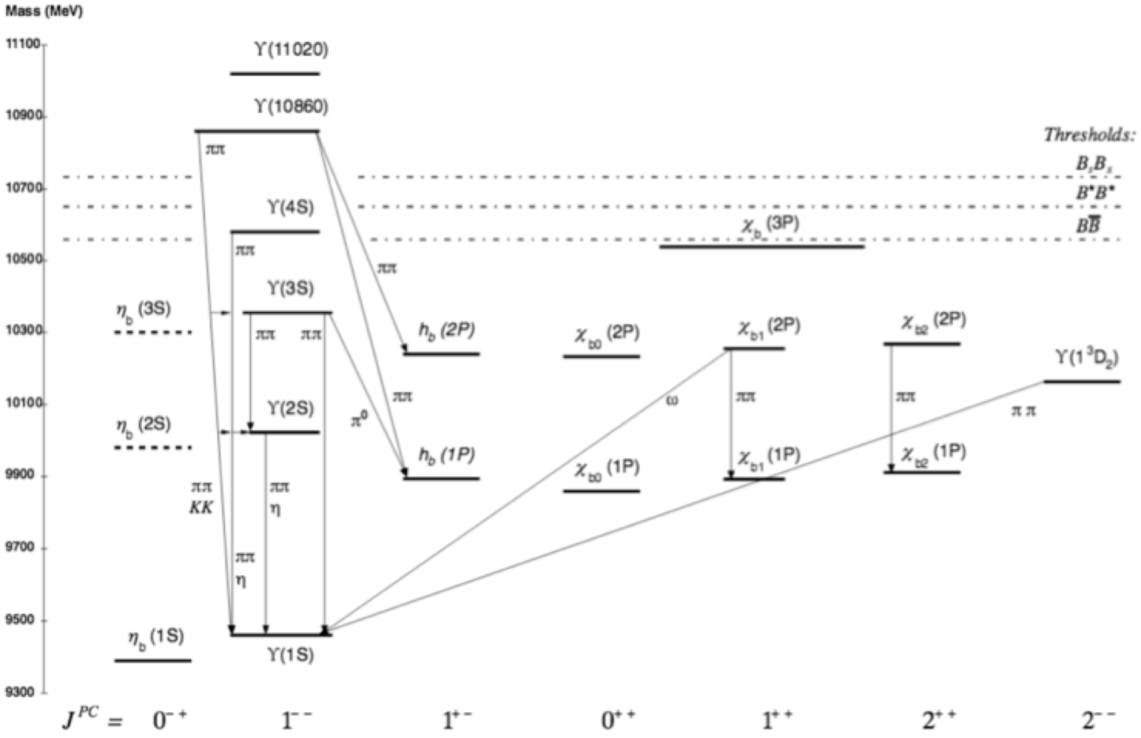
[61] N. A. Abdulov and A. V. Lipatov, Bottomonium production and polarization in the NRQCD with kT - factorization. III: Υ(1S) and χb(1P) mesons, Eur. Phys. J. C 81, 1085 (2021), arXiv:2011.13401.
[62] N. A. Abdulov and A. V. Lipatov, Bottomonia production and polarization in the NRQCD with kT - factorization. II: Υ(2S) and χb(2P) mesons, Eur. Phys. J. C 80, 486 (2020), arXiv:2003.06201.
[63] N. A. Abdulov and A. V. Lipatov, Bottomonia production and polarization in the NRQCD with kT - factorization. I: Υ(3S) and χb(3P) mesons, Eur. Phys. J. C 79, 830 (2019), arXiv:1909.05141.

Global analysis

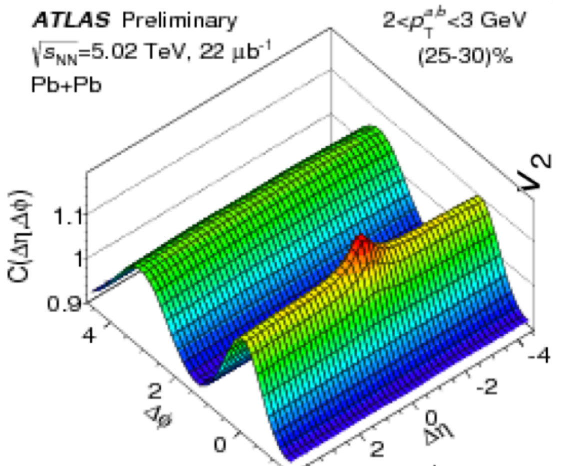
Basic principle:

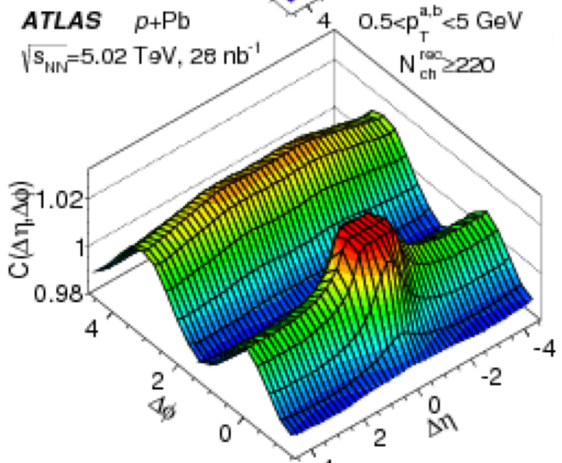
Particles with the same quark content and same masses shall have the same kinematics

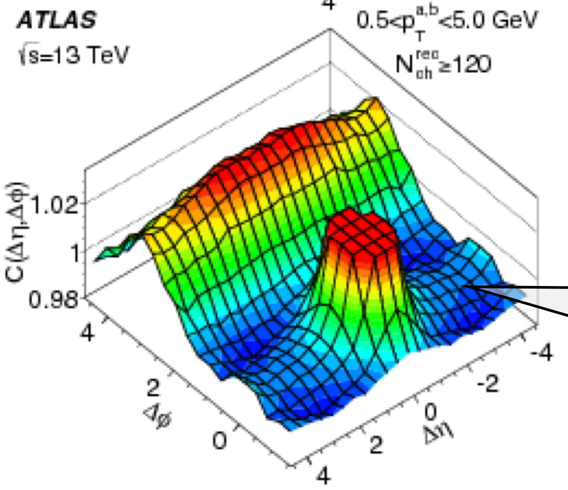
The extent of deviation due to a 10% difference in masses can be tested with the m_T – scaling

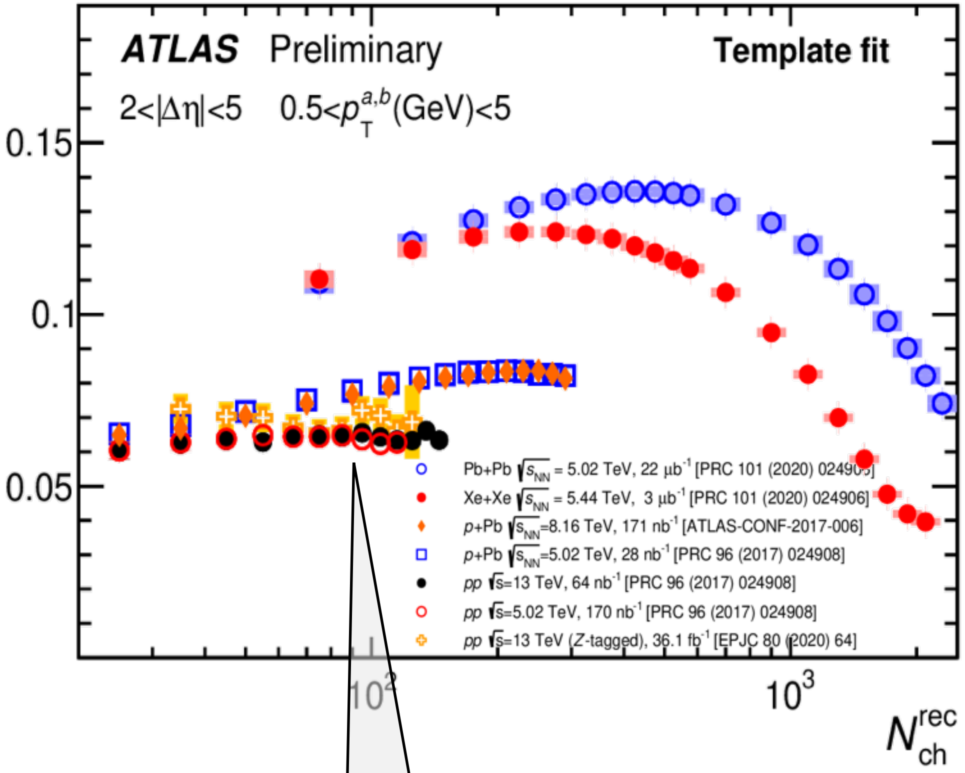


QGP signatures in small systems



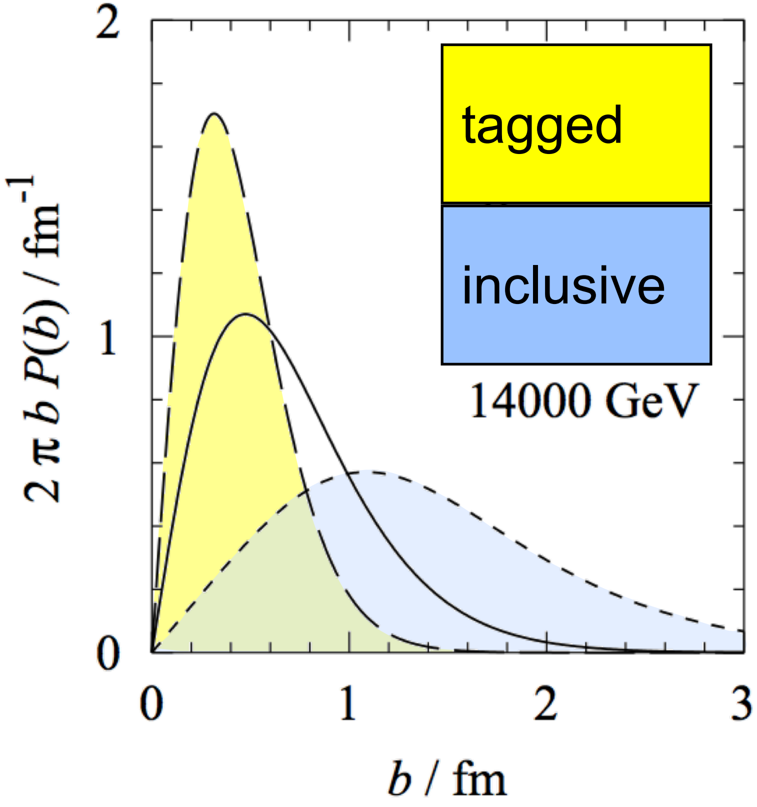






$$\frac{dN}{d\varphi} = \frac{N}{2\pi}\left[1 + 2\sum_n \mathbf{v_n} cos(n(\varphi - \Phi_n))\right]$$

Collectivity in *pp*



Two-particle correlations in *pp* are independent of n_{ch}

Do they depend on b_{imp}?

We checked it with events tagged by Z boson.