Constraining η/s through high-$p_\perp$ tomography

Bithika Karmakar
Institute of Physics Belgrade

Collaborators: Dusan Zigic, Igor Salom, Jussi Auvinen, Pasi Huovinen, Marko Djordjevic and Magdalena Djordjevic

ExploreQGP

30 May 2023

European Research Council
Established by the European Commission

Outline

- Introduction
- η/s of the medium : Soft-to-hard boundary
- High-$p_\perp$ energy loss: Generalized DREENA-A
- Phenomenological approach to constrain η/s
- Theoretical approach to evaluate η/s
- Conclusion

Introduction

- Low-$p_\perp$ observables are used to explore the bulk properties of the QGP created in heavy-ion collisions.
- High-$p_\perp$ probes also become powerful tomography tools; Sensitive to global QGP features, e.g., different temperature profiles or initial conditions.
- The near perfect fluidity of QGP has been investigated extensively in heavy-ion collision experiments.
- η/s is well constrained by Bayesian analysis in low-$p_\perp$ sector in the temperature range $T_c \lesssim T \lesssim 1.5T_c$ and weakly constrained at larger temperatures.
- We try to put constraints on η/s by analyzing high-$p_\perp$ observables using the generalized DREENA-A.

η/s of the medium : Soft-to-hard boundary

- QGP is expected to behave as weakly interacting gas: Weakly coupled
- Fluid dynamics predicts the η/s to be very low: Strongly coupled
- QGP may behave as perfect fluid near T_c (soft regime) and η/s may increase at high temperature (hard regime).
- Testing the soft-to-hard hypothesis is difficult: Anisotropy is weakly affected by the η/s at high temperature.
- High-$p_\perp$ data/theory can serve as complementary tool.

High-$p_\perp$ energy loss : Generalized DREENA-A

- **D**ynamical **R**adiative and **E**lastic **EN**ergy loss **A**pproach
  - ○ Based on finite temperature field theory and generalized HTL approach

    M. Djordjevic, PRC 74, 064907, (2006) ; PRC 80, 064909 (2009), M. Djordjevic and U. Heinz, PRL 101, 022302
  - ○ Finite size dynamical QCD medium is considered
  - ○ Takes into account both radiative and collisional energy losses
  - ○ Generalized to the case of magnetic mass and running coupling
  - ○ No fitting parameter in the theory
- Takes arbitrary temperature profile as input.

  D. Zigic, I. Salom, J. Auvinen, P. Huovinen, M. Djordjevic Front.in Phys. 10 (2022) 957019
- Optimized to incorporate any arbitrary event-by-event fluctuating temperature profile.

  D. Zigic, J. Auvinen, I. Salom, M. Djordjevic, P. Huovinen Phys.Rev.C 106 (2022) 4, 044909
- DREENA-A is available on http://github.com/DusanZigic/DREENA-A (Details in talk by Dusan Zigic, next session)

Phenomenological approach

- Three different $(\eta/s)(T)$ parametrizations have been considered.
- Parameters are adjusted to reproduce low-$p_\perp$ data.
- Temperature profile is generated for each case.
- High-$p_\perp$ predictions found using generalized DREENA-A.
- Compared with high-$p_\perp$ data.

Modeling the bulk evolution

- Initial entropy profiles are generated using TRENTo model.
- 10^4 events for Pb+Pb ($s\ = 5.02$ TeV) and Au+Au ($s\ = 200$ GeV) collisions.
- Events are sorted in centrality classes.
- Initial free streaming is not preferred by high-$p_\perp$ data.

  S. Stojku, J. Auvinen, M. Djordjevic, P. Huovinen and M. Djordjevic, Phys. Rev. C 105 (2022) 2, L021901
- Onset time for hydrodynamics: $\tau_0 = 1fm$.

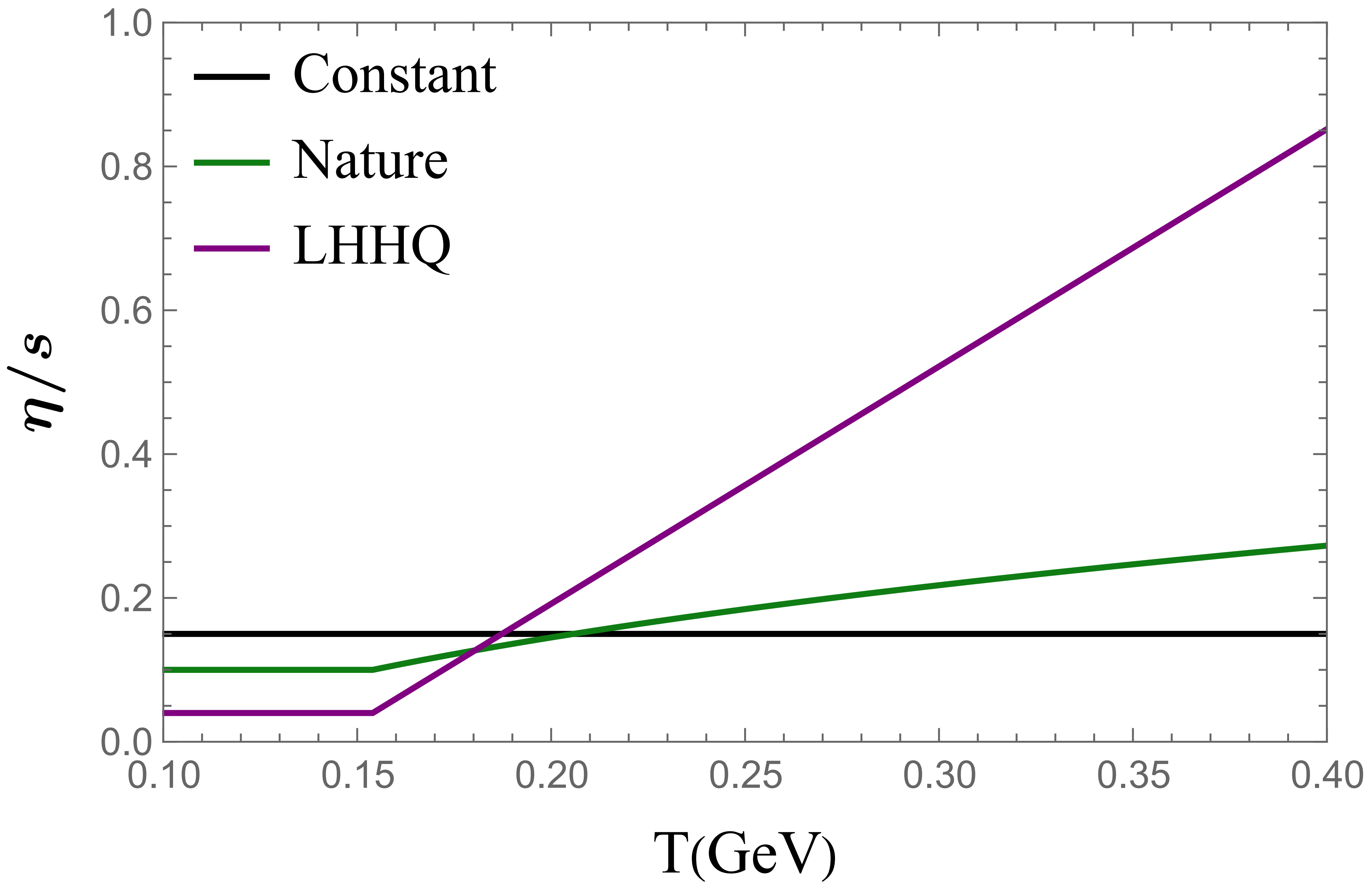
  S. Stojku, J. Auvinen, M. Djordjevic, P. Huovinen and M. Djordjevic, Phys. Rev. C 105 (2022) 2, L021901
- (2+1)-dimensional fluid dynamical model (VISHNew) used to simulate the medium evolution.

Temperature dependence of η/s

$$(\eta/s)(T) = \begin{cases} (\eta/s)_{\text{min}}, & T < T_c, \\ (\eta/s)_{\text{min}} + (\eta/s)_{\text{slope}}(T - T_c)\dfrac{T}{T_c}^{(\eta/s)_{\text{crv}}}, & T > T_c \end{cases}$$

Nature: Nature Phys. 15, no. 11, 1113-1117 (2019)

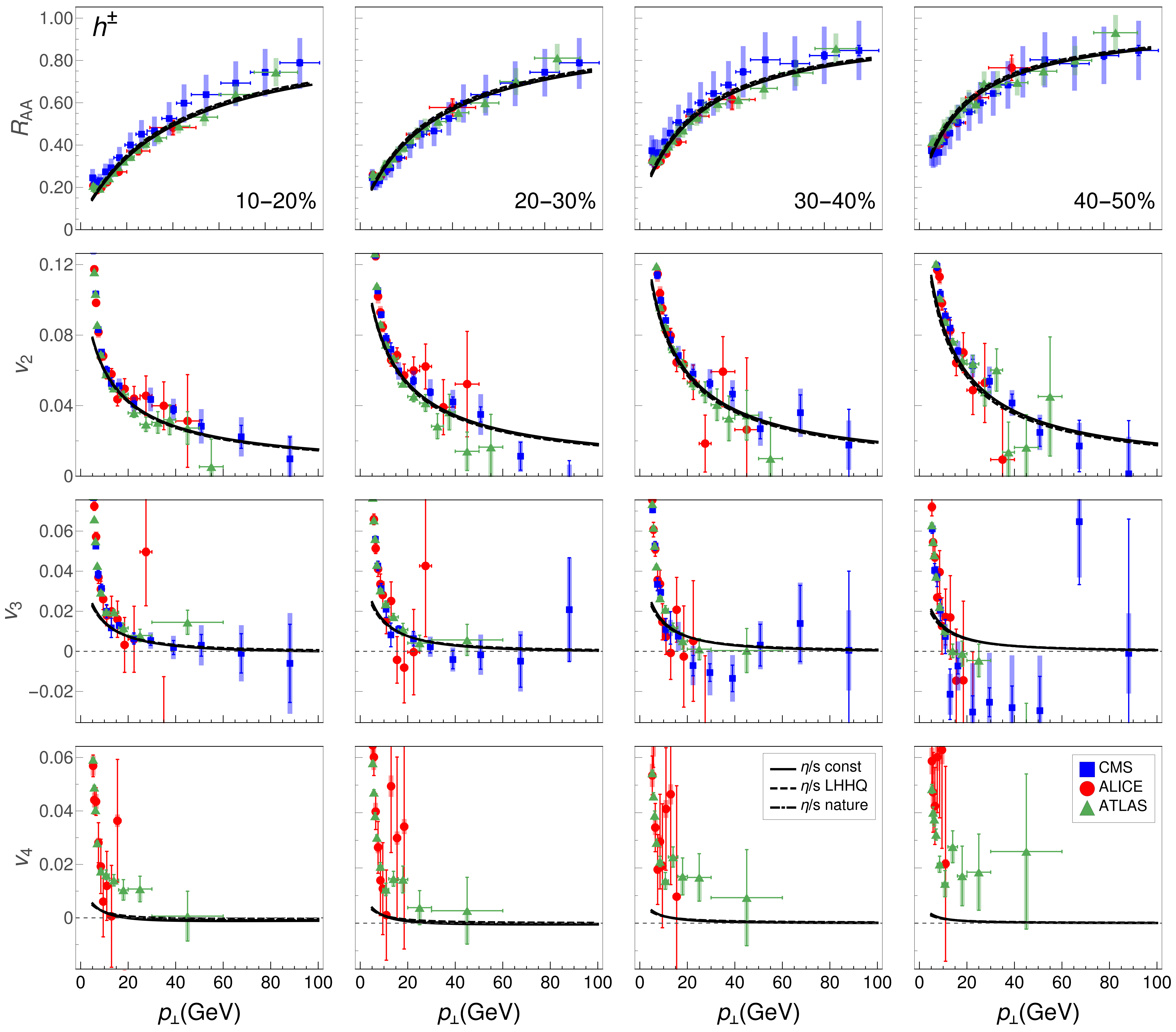
LHHQ: Phys. Rev. Lett. 106, 212302 (2011)



- Pion, kaon, proton multiplicities and $v_2\{4\}$ are reproduced by varying the TRENTo normalization factor for three η/s parametrizations.

Results

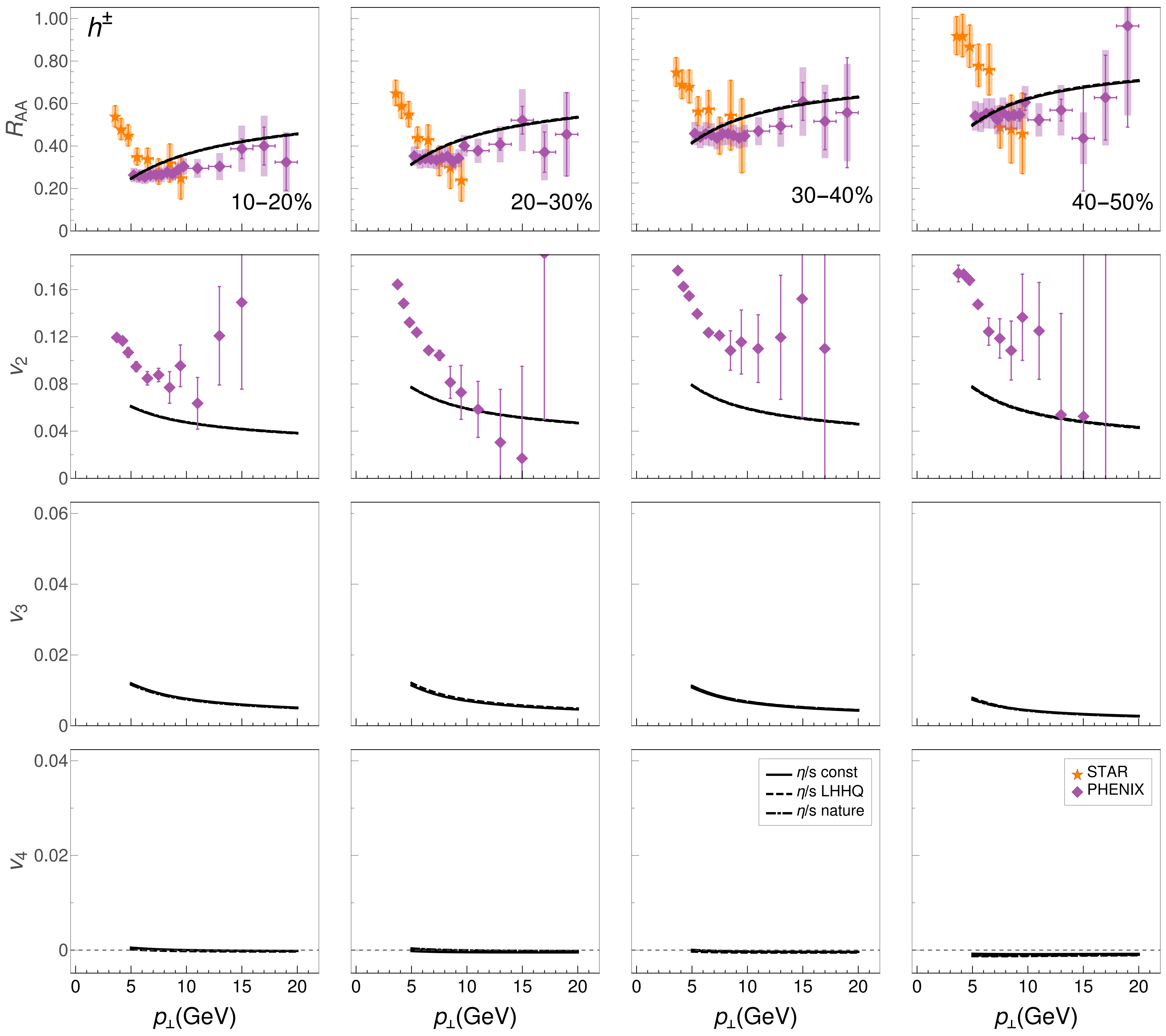
BK, D. Zigic, I. Salom, J. Auvinen, P. Huovinen, M. Djordjevic and M. Djordjevic arXiv:2305.11318



Pb + Pb (s = 5.02TeV)

Results

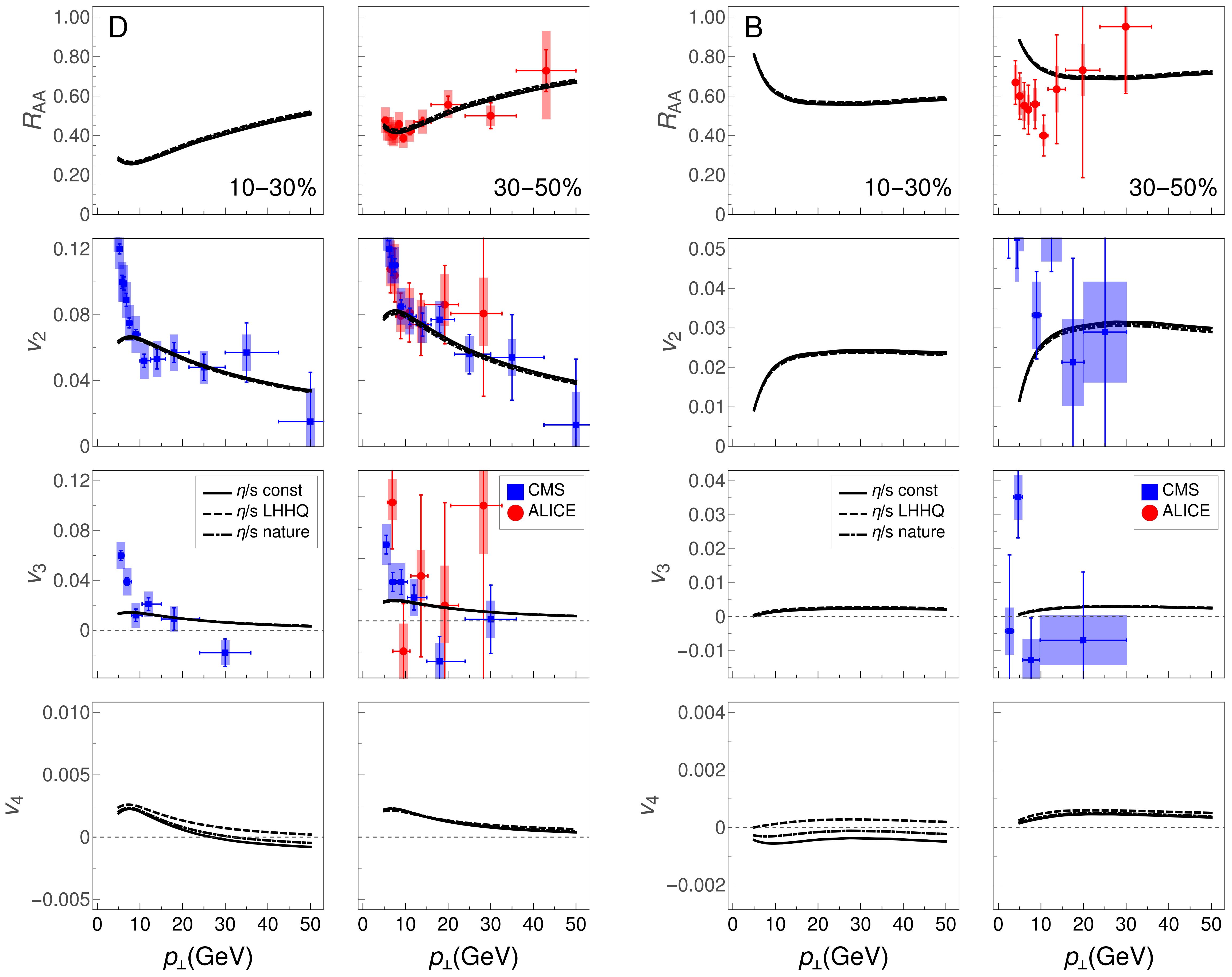
BK, D. Zigic, I. Salom, J. Auvinen, P. Huovinen, M. Djordjevic and M. Djordjevic arXiv:2305.11318



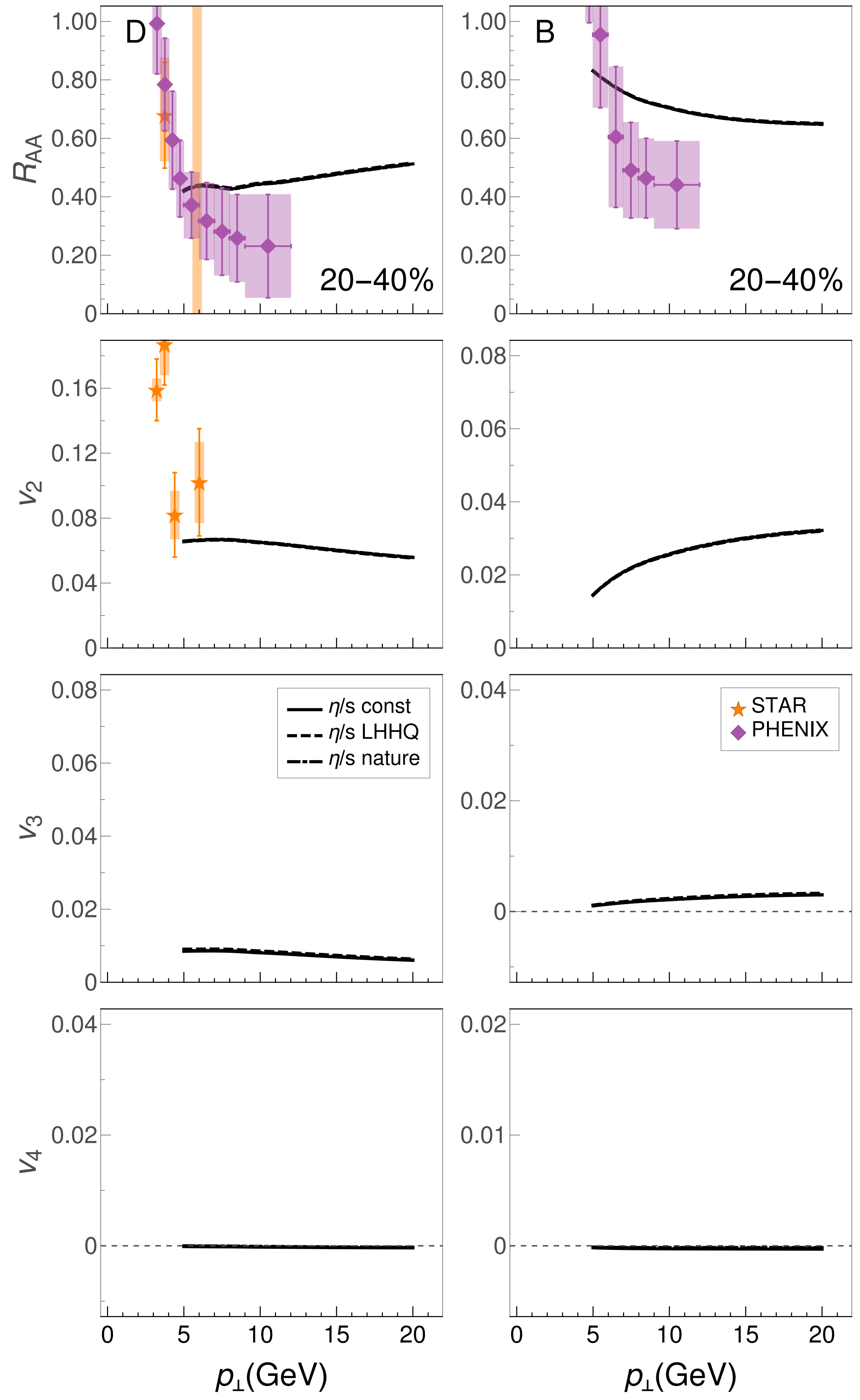
$\mathrm{Au} + \mathrm{Au}\ (s\ = 200GeV)$

Results

BK, D. Zigic, I. Salom, J. Auvinen, P. Huovinen, M. Djordjevic and M. Djordjevic arXiv:2305.11318



Pb + Pb ($s = 5.02 TeV$)

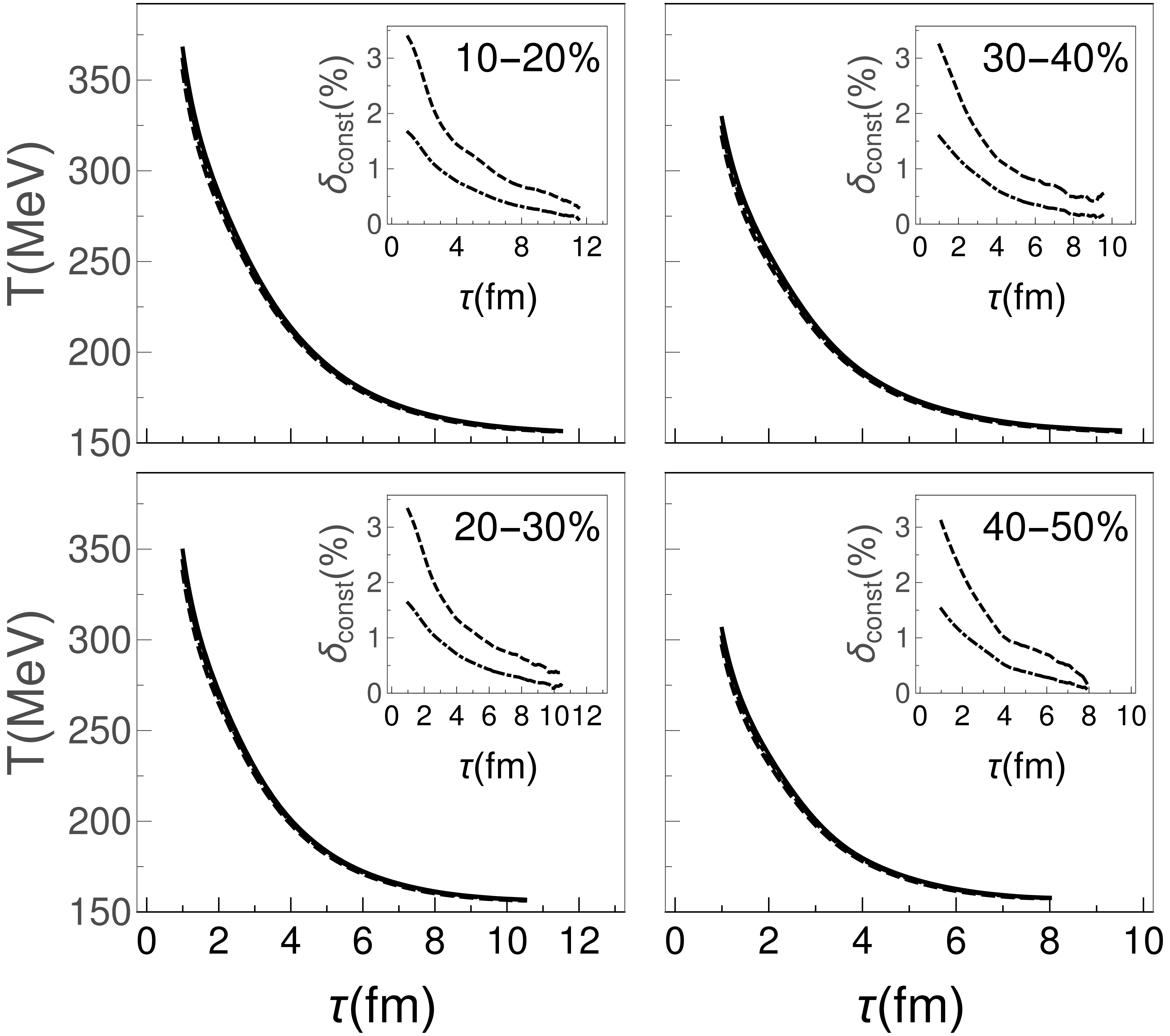


Au + Au ($s = 200 GeV$)

Average jet perceived temperature

BK, D. Zigic, I. Salom, J. Auvinen, P. Huovinen, M. Djordjevic and M. Djordjevic arXiv:2305.11318

- Pb + Pb $s = 5.02$ TeV
- Full = LHHQ; DotDashed = Nature,
  Dashed = Constant
- Inset: Dotdashed = Nature,
  Dashed = LHHQ
- Temperature difference during evolution
  is very small.
- Insufficient to lead to observable difference
  in the results.



Theoretical approach :

Transport coefficient from dynamical energy loss formalism

BK, D. Zigic, I. Salom, J. Auvinen, P. Huovinen, M. Djordjevic and M. Djordjevic arXiv:2305.11318

- Transport coefficient ($\hat{q}$) ≡ Squared average transverse momentum exchange between the medium and the fast parton per unit length

- Interaction between the parton and medium is characterized by the HTL resummed elastic collision rate:

$$\frac{d\Gamma_{el}}{d^2q} = 4C_A\, 1 + \frac{n_f}{6} T^3 \frac{\alpha_s^2}{q^2\, q^2 + \mu_E^2}$$

- After including running coupling and finite magnetic mass:

$$\frac{d\Gamma_{el}}{d^2q} = \frac{C_A}{\pi} T\alpha(ET) \frac{\mu_E^2 - \mu_M^2}{(q^2 + \mu_E^2)(q^2 + \mu_M^2)}$$

- In fluid rest frame:

$$\hat{q} = \int_0^{6ET} d^2q\, q^2 \cdot \frac{d\Gamma_{el}}{d^2q} = C_A T \frac{4\pi}{(11 - \frac{2}{3}n_f)} \frac{\mu_E^2 \ln \frac{6ET + \mu_E^2}{\mu E^2} - \mu_M^2 \ln \frac{6ET + \mu_M^2}{\mu_M^2}}{\ln(\frac{ET}{\Lambda^2})}$$

η/s from the transport coefficient

BK, D. Zigic, I. Salom, J. Auvinen, P. Huovinen, M. Djordjevic and M. Djordjevic arXiv:2305.11318

- In the limit $ET \to \infty$:

$$\hat{q} = C_A \frac{4\pi}{11 - \frac{2}{3}n_F}^{\,2} \frac{4\pi\, 1 + \frac{n_F}{6}}{W(\xi(T))}(1 - x_{ME}^2)T^3$$

- $\xi(T) = \frac{1 + \frac{n_f}{6}}{11 - \frac{2}{3}n_f} \frac{4\pi T^2}{\Lambda}$, $W \equiv$ Lambert's W function, $x_{ME} = \mu_M/\mu_E$
- $\hat{q}$ is weakly dependent on jet energy E.
- In weakly coupled limit:

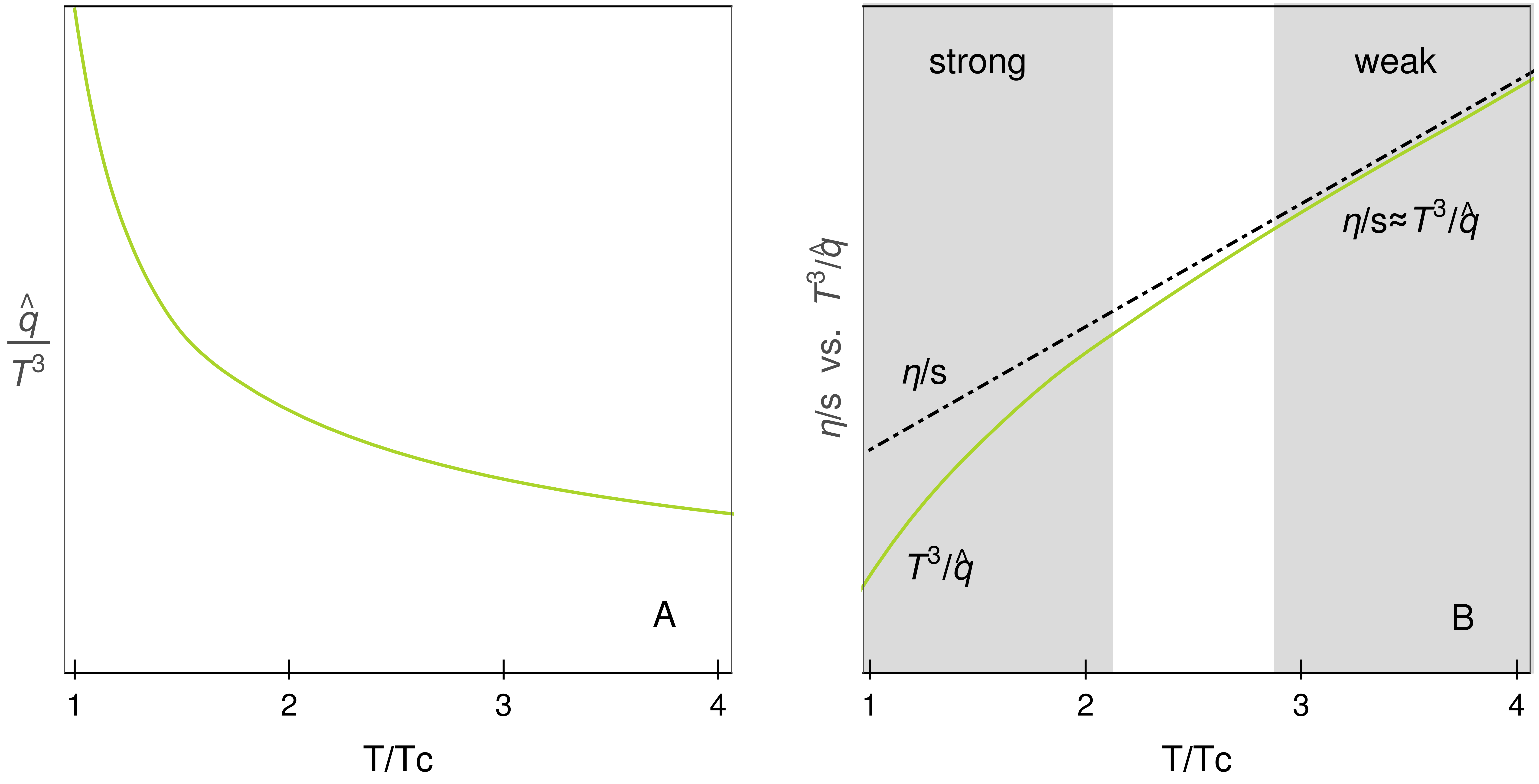
$$\eta/s \approx 1.25T^3/\hat{q}$$

Phys. Rev. Lett. 99 192301 (2007), Phys. Rev. D 104, L071501 (2021)

η/s from the transport coefficient

BK, D. Zigic, I. Salom, J. Auvinen, P. Huovinen, M. Djordjevic and M. Djordjevic arXiv:2305.11318

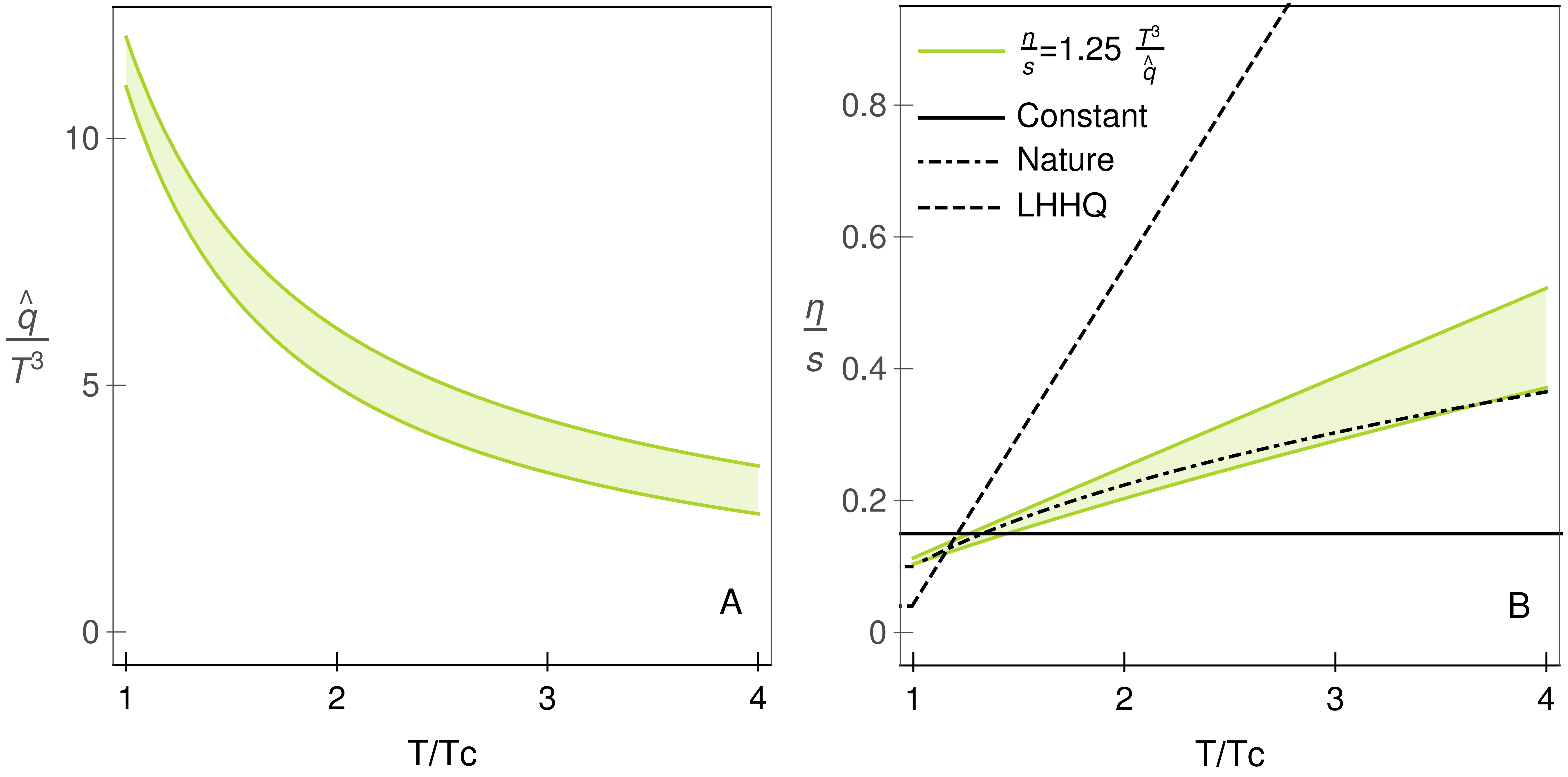
- $\hat{q}$ quantifies the parton coupling strength in the medium
- $\hat{q}/T^3$ must rise rapidly near T_c from above.
- Our formalism valid in weakly coupled regime.
- $T^3/\hat{q}$ and η/s should agree in the weak coupling regime.
- Soft-to-hard boundary



η/s from the transport coefficient

BK, D. Zigic, I. Salom, J. Auvinen, P. Huovinen, M. Djordjevic and M. Djordjevic arXiv:2305.11318

- $\hat{q}/T^3$ shows expected behavior.
- Enhanced quenching near T_c.
- η/s is surprisingly close to the constraints from Bayesian analysis.



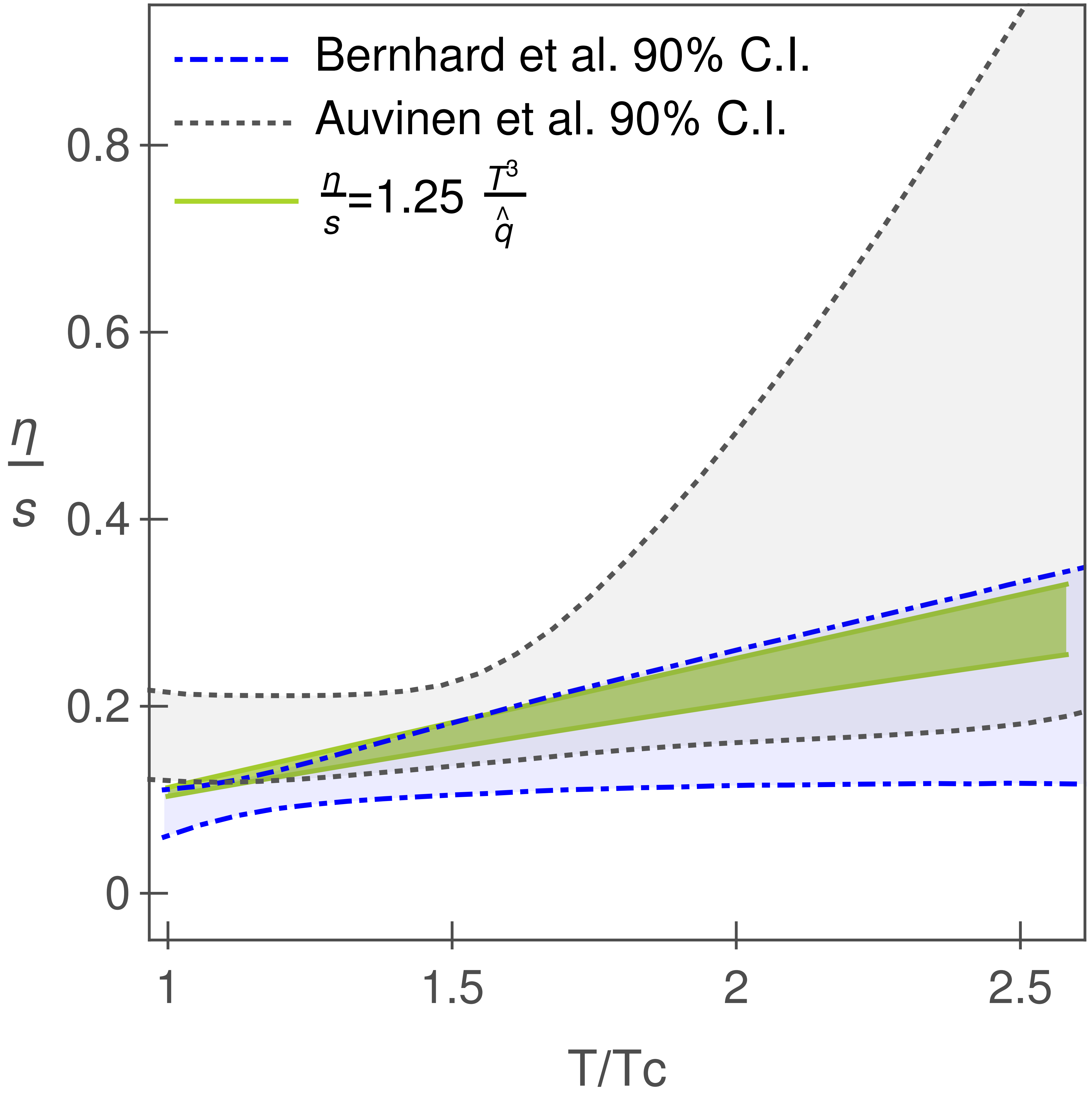
η/s from the transport coefficient

BK, D. Zigic, I. Salom, J. Auvinen, P. Huovinen, M. Djordjevic and M. Djordjevic arXiv:2305.11318

- Uncertainty due to initial jet energy is very small
- Surprisingly close to the parametrization inspired by the Bayesian analysis.
- Does not drop significantly below the inferred η/s values near T_c .
- No soft-to-hard boundary.

Blue → Nature Phys. 15, no. 11, 1113-1117 (2019)

Black → Phys. Rev. C 102, 044911 (2020)



Conclusion

- We use generalized DREENA-A to compute high-$p_\perp$ energy loss.
- In the phenomenological approach:
  - Three different $(\eta/s)(T)$ parametrizations have been considered.
  - The predictions from the generalized DREENA-A for three η/s scenarios lead to plots that are almost indistinguishable.
  - The difference in the average jet-perceived temperature for the three cases is less than 2%.
  - High-$p_\perp$ observables are not sensitive to such small temperature difference.
- In the theoretical approach:
  - Transport coefficient and jet quenching strength are calculated from the dynamical energy loss formalism.
  - η/s shows surprisingly good agreement all the way to T_c with constraints extracted from existing Bayesian analyses. Provides much smaller uncertainties at high temperature.
  - Intriguing hypothesis: quasiparticle picture is consistent at the entire temperature range.
  - No guidance on locating soft-to-hard boundary.

Thank you for your attention

- Mass of light quark $M = \mu_E/6$
- Mass of charm and bottom quark 1.2 GeV and 4.75 GeV
- Gluon mass $m_g = \mu_E/2$
- $\mu_M/\mu_E = 0.6$
- $\Lambda_{QCD} = 0.2GeV$

○ Constant η/s (0.15 for Pb+Pb and 0.12 for Au+Au collision)

○ Nature: $(\eta/s)_{\text{min}} = 0.1,(\eta/s)_{\text{slope}} = 1.11,(\eta/s)_{\text{crv}} = -0.48$

○ LHHQ: $(\eta/s)_{\text{min}} = 0.04,(\eta/s)_{\text{slope}} = 3.30,(\eta/s)_{\text{crv}} = 0$

Modeling the bulk evolution

- Particlization is performed using Cooper-Frye prescription at isothermal space time hypersurface at 151 MeV.
- UrQMD is used to simulate microscopic dynamics of hadronic system.
- Bulk viscosity parametrized as

$$(\zeta/s)(T) = \frac{(\zeta/s)_{\text{max}}}{1 + \frac{T - T_0}{(\zeta/s)_{\text{width}}}^{2}}$$

with $(\zeta/s)_{max} = 0.03$, $(\zeta/s)_{width} = 0.022$ and $T_0 = 0.183 GeV$.

- We use lattice QCD based EoS from HotQCD (high temperature) + HRG (low temperature) EoS.