

# Fluctuations and correlations of the anisotropic flow

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# Motivation

- azimuthal anisotropies of hadron transverse momentum distribution  $\Leftarrow$  anisotropies in the initial conditions
- measurable access to the initial state—profile of the energy-momentum deposition
- initial states change from event to event  $\Rightarrow$  e-by-e fluctuations of anisotropic flow
- Here: initial transverse energy-momentum profile depends on (space-time) rapidity  $\Rightarrow$  (pseudo-)rapidity dependence of anisotropic flow
- breaking of Bjorken symmetry (even at high collision energies)
  - $\Rightarrow$  information about energy deposition into the initial state
  - $\Rightarrow$  discriminatory power over initial state models
- Test of initial conditions also at lower collision energies
- This report: Au+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV and  $\sqrt{s_{NN}} = 27$  GeV

# The hybrid model setup

- Initial conditions
  - Glauber Monte Carlo: GLISSANDO 2  
[M. Rybczyński, G. Stefanek, W. Broniowski, P. Božek, Comp.Phys.Commun. 185 (2014) 1579]
  - uRQMD  
[S.A.Bass, et al., Prog. Part. Nucl. Phys. 41 (1998) 225. M. Bleicher, et al., J. Phys. G: Nucl. Part. Phys. 25 (1999) 1859]
- hydrodynamic evolution: 3D viscous model vHLLÉ  
[Iu. Karpenko, P. Huovinen, M. Bleicher, Comput.Phys.Commun. 185 (2014) 3016]
- hadronic phase simulated by transport model: uRQMD

## GLISSANDO – extension to 3rd dimension

- original GLISSANDO: entropy density in transverse plane according to the mixed model

$$s(x, y, \quad) = \kappa \sum_i \left[ (1 - \alpha) + N_i^{\text{coll}} \alpha \right] \exp \left( - \frac{(x - x_i)^2 + (y - y_i)^2}{2\sigma^2} \right)$$

## GLISSANDO – extension to 3rd dimension

- original GLISSANDO: entropy density in transverse plane according to the mixed model
- space-time rapidity profile  $f_{\pm}(\eta_s)$  added

P. Božek, W. Broniowski, Phys. Rev. C 85 (2012) 044910

$$s(x, y, \eta_s) = \kappa \sum_i f_{\pm}(\eta_s) \left[ (1 - \alpha) + N_i^{\text{coll}} \alpha \right] \exp \left( -\frac{(x - x_i)^2 + (y - y_i)^2}{2\sigma^2} \right)$$

$$f_{\pm}(\eta_s) = \frac{\eta_M \pm \eta_s}{2\eta_M} H(\eta_s) \quad \text{for } |\eta_s| < \eta_M$$

$$H(\eta_s) = \exp \left( -\frac{(|\eta_s| - \eta_0)^2 \Theta(|\eta_s| - \eta_0)}{2\sigma_{\eta}^2} \right)$$

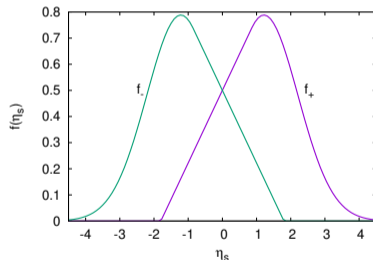
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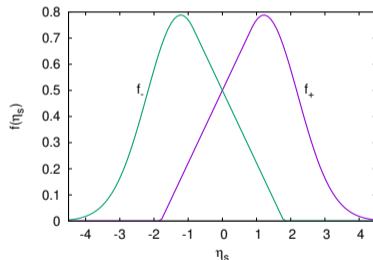
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Baryon number deposition:

$$n_B(x, y, \eta_s) = \kappa_B \sum_i \exp \left( -\frac{(\eta_B \pm \eta_s)^2}{2\sigma_B^2} \right) \exp \left( -\frac{(x - x_i)^2 + (y - y_i)^2}{2\sigma^2} \right)$$

- energy, momentum, baryon number and charge extracted from the participating hadrons
- smeared among the hydrodynamic cells

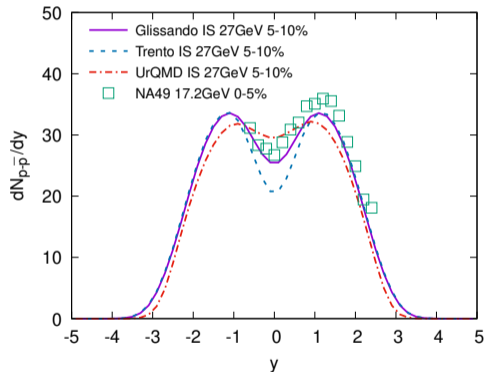
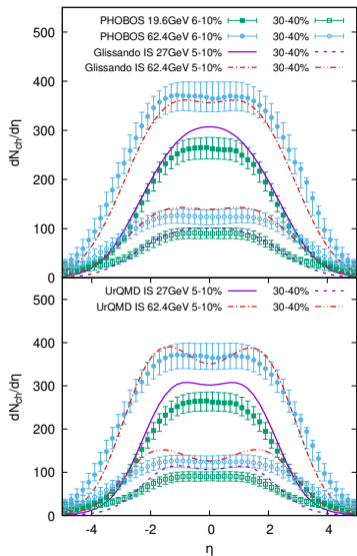
$$w \propto \exp \left( -\frac{(x_h - x_c)^2}{R_T^2} - \frac{(y_h - y_c)^2}{R_T^2} - \gamma^2 \tau_0^2 \frac{(\eta_h - \eta_c)^2}{R_\eta^2} \right)$$



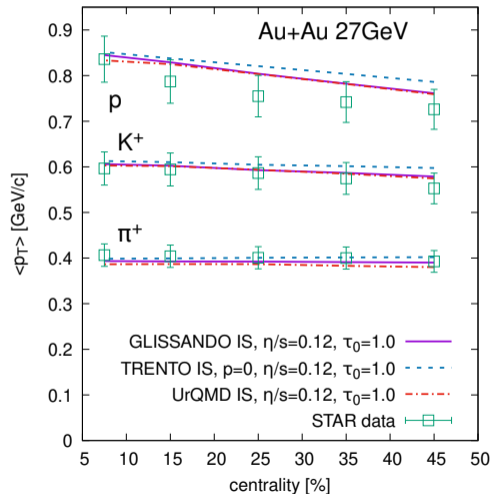
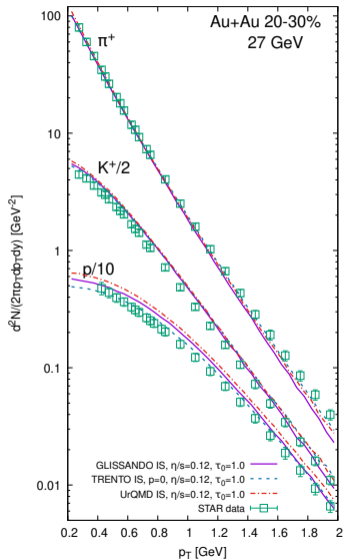
# Hydro and transport part

- Hydrodynamics: vHLLE
  - Muller-Israel-Stewart equations
  - chiral EoS (crossover)
  
- particlisation
  - Cooper-Frye sampling of hadron production
  - viscous corrections included
  - oversampling
  
- transport model: uRQMD

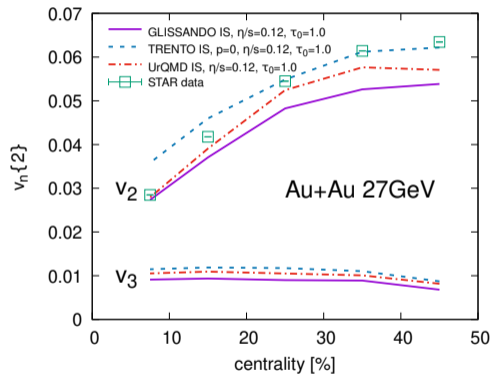
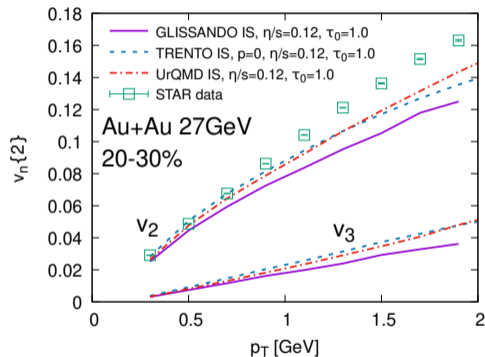
# Tuning of the model—fitting basic observables—rapidity spectra



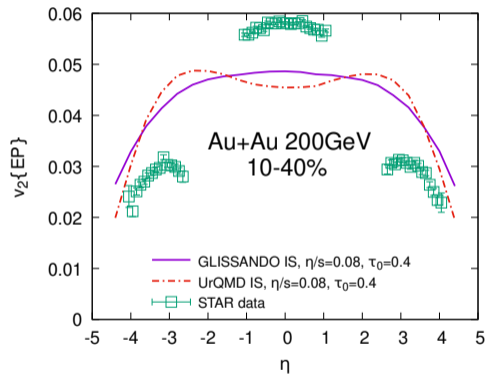
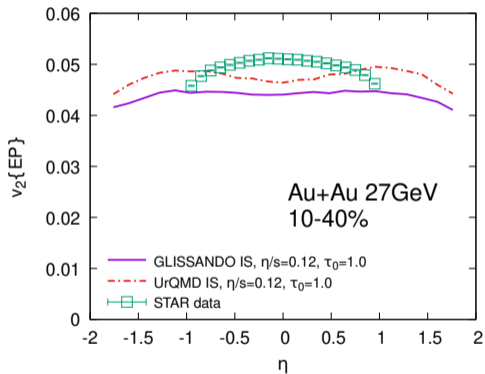
# Tuning of the model—fitting basic observables— $p_T$ spectra



# Tuning of the model—fitting basic observables—elliptic flow

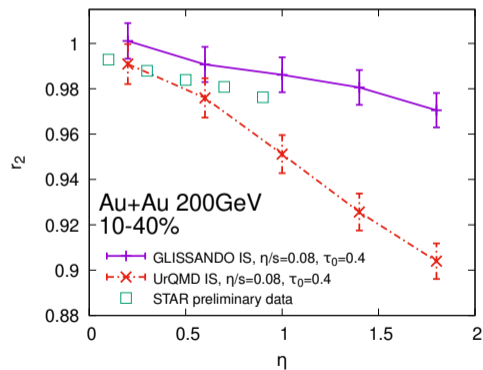
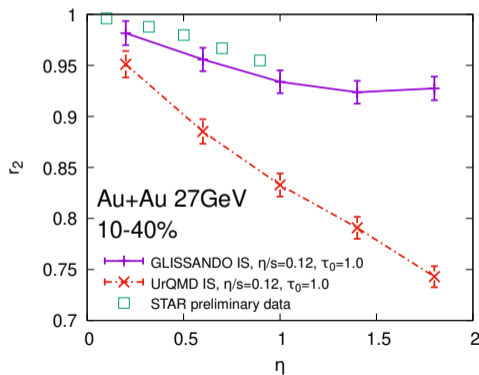


# Results: rapidity dependence of the elliptic flow



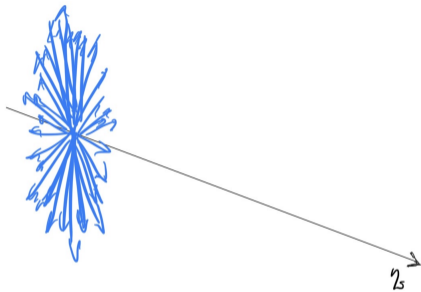
# Longitudinal decorrelation of anisotropic flow

$$r_n(\eta) = \frac{\langle q_n(-\eta)q_n^*(\eta_{\text{ref}}) \rangle}{\langle q_n(\eta)q_n^*(\eta_{\text{ref}}) \rangle}.$$



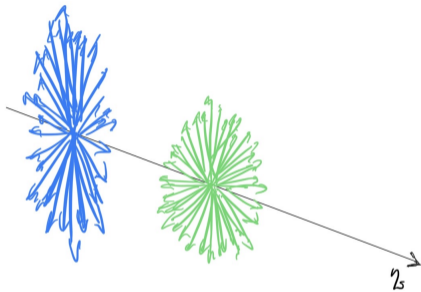
# The mechanism of decorrelation

varying amplitude of the flow anisotropy



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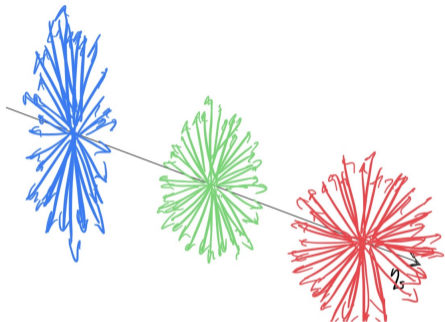
varying amplitude of the flow anisotropy





# The mechanism of decorrelation

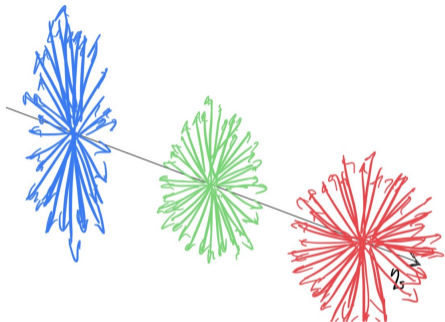
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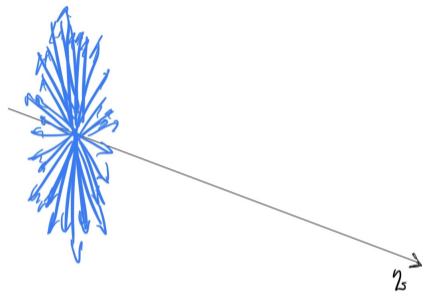
$$r_n^v(\eta) = \frac{\langle v_n(-\eta)v_n(\eta_{\text{ref}}) \rangle}{\langle v_n(\eta)v_n(\eta_{\text{ref}}) \rangle}$$

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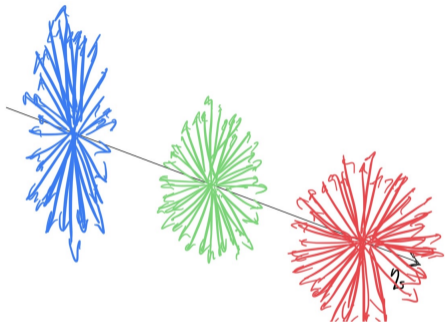
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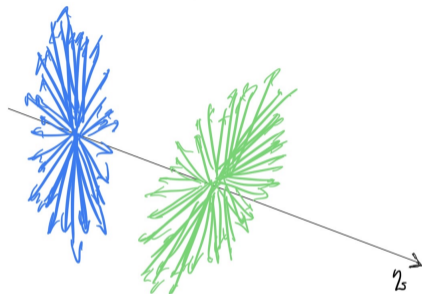
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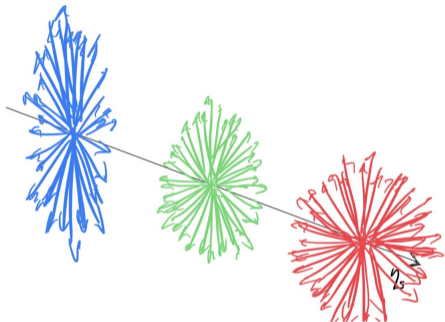
varying angle of the flow anisotropy



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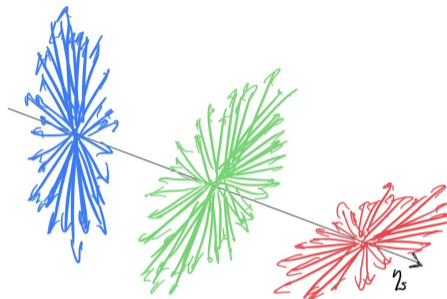
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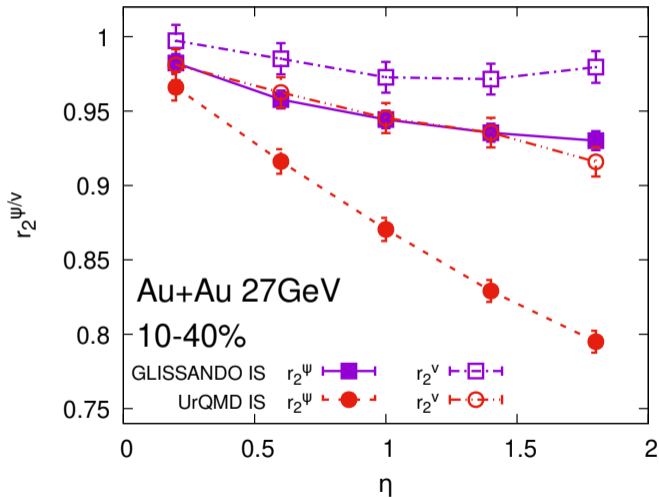
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varying angle of the flow anisotropy

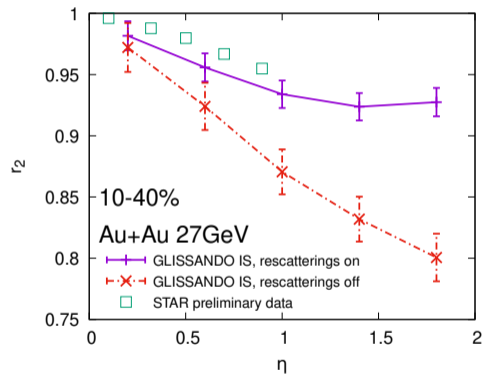
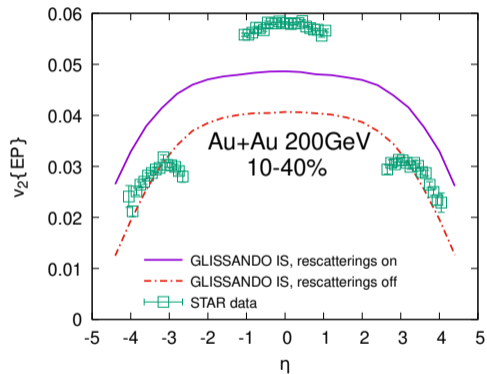


$$r_n^\psi(\eta) = \frac{\langle \cos[n(\Psi_n(-\eta) - \Psi_n(\eta_{\text{ref}}))] \rangle}{\langle \cos[n(\Psi_n(\eta) - \Psi_n(\eta_{\text{ref}}))] \rangle}$$

# The mechanism of anisotropy decorrelation

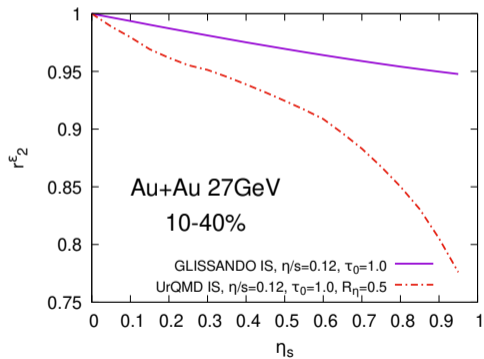
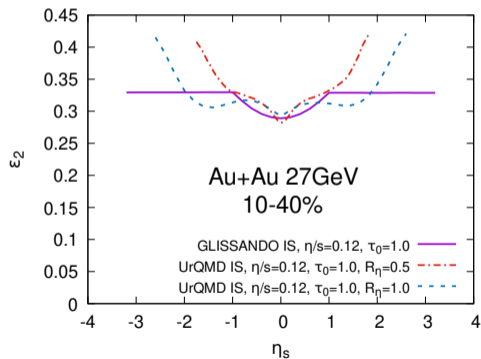


# The effect of rescattering



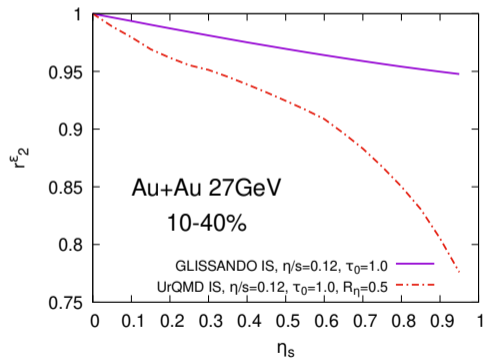
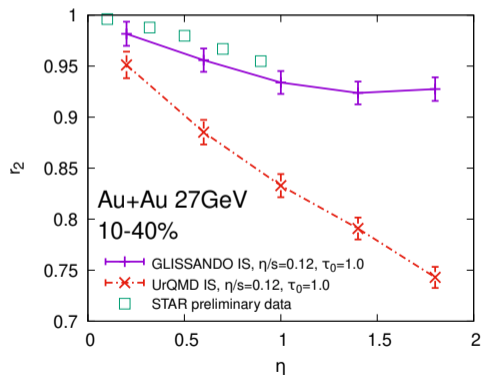
# Decorrelation in the initial conditions

$$r_n^\epsilon(\eta_s) = \frac{\langle \epsilon_n(-\eta_s) \epsilon_n(\eta_{s,\text{ref}}) \cos[n(\Psi_n(-\eta_s) - \Psi_n(\eta_{s,\text{ref}}))] \rangle}{\langle \epsilon_n(\eta_s) \epsilon_n(\eta_{s,\text{ref}}) \cos[n(\Psi_n(\eta_s) - \Psi_n(\eta_{s,\text{ref}}))] \rangle}.$$



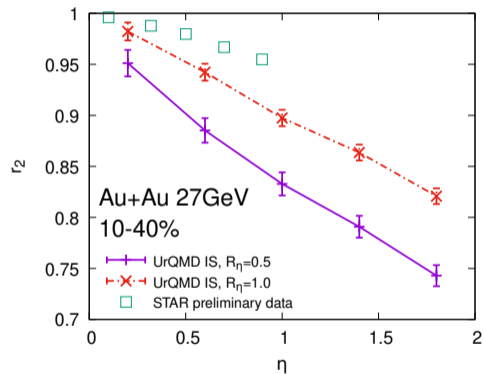
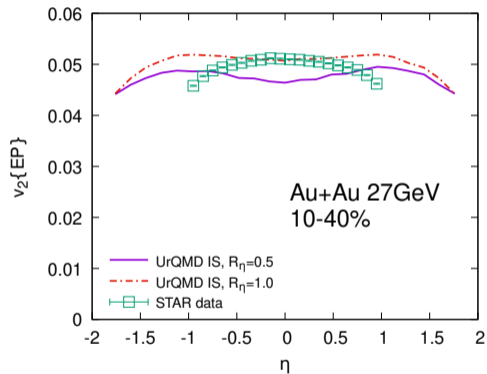
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# The effect of smearing in uRQMD initial conditions



# Conclusions

- longitudinal flow anisotropy decorrelation seems to prefer Glauber MC over uRQMD initial conditions for Au+Au collisions at 27 GeV
  - mainly caused by flow angle decorrelations (like at 200 GeV)
  - corresponds to decorrelation of the initial conditions
  - too small rapidity smearing in uRQMD initial conditions
  
- Outlook: look at decorrelation in transverse momentum  
[ALICE collab: arXiv:2206.04574]

## References:

- J. Cimerman, *et al.*, Phys.Rev.C 104 (2021) 014904  
J. Cimerman, *et al.*, Phys.Rev.C 103 (2021) 034902

# BACKUP

## Values of parameters

$\sqrt{s_{NN}}$ [GeV]	27	72	200
$\tau_0$ [fm/c]	1.0	0.7	0.4
$\eta/s$	0.12	0.08	0.08
$\eta_0$	$0.89 - 0.2\chi$	1.8	1.5
$\sigma_\eta$	$1.09 - 0.2\chi$	0.7	1.4
$\eta_M$	1.0	1.8	3.36
$\eta_B$	$1.33 - 0.32\chi$	2.2	
$\sigma_B$	$0.79 - 0.21\chi$	1.0	