

Numerical Simulations of Faraday Waves in Binary Bose-Einstein Condensates*

Antun Balaž¹ and Alexandru Nicolin²

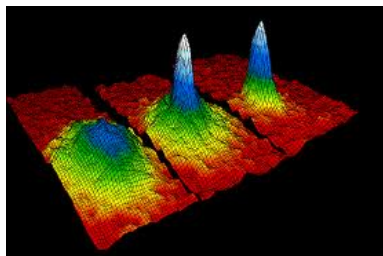
¹Scientific Computing Laboratory, Institute of Physics Belgrade,
University of Belgrade, Pregrevica 118, 11080 Belgrade, Serbia

²“Horia Hulubei” National Institute for Physics and Nuclear Engineering,
Department of Theoretical Physics, P.O.B. MG-6, 077125 Bucharest, Romania

*Supported by: Serbian Ministry of Education and Science under Grant ON171017, CNCS-UEFISCDI under Grant PD122, Contract No. 35/28.07.2010, and European Commission through projects EGI-InSPIRE, PRACE-1IP, PRACE-2IP, and HP-SEE.

Why ultracold atoms are interesting?

- Intensive progress in the field of ultracold atoms has been recognized by Nobel prize for physics in 2001 for experimental realization of Bose-Einstein condensation
- Cold alkali atoms:
Rb, Na, Li, K...
 $T \sim 1 \text{ nK}$, $\rho \sim 10^{14} \text{ cm}^{-3}$
- Cold bosons, cold fermions
- Harmonic trap, optical lattice
- Short-range interactions,
long-range dipolar interactions
- Tunable quantum systems concerning dimensionality, type and strength of interactions



What are Faraday waves?

- M. Faraday, Philos. Trans. R. Soc. London **121**, 299 (1831):

When the upper surface of a plate vibrating so as to produce sound is covered with a layer of water, the water usually presents a beautifully crisped appearance, the crispations being produced more readily and beautifully when there is a certain quantity than when there is less. For small crispations, the water should flow upon the surface freely. Large crispations require more water than small ones. Too much water sometimes interferes with the beauty of the appearance, but the crispation is not incompatible with much fluid, for the depth may amount to eight, ten, or twelve inches, and is probably unlimited.

[crispation: curled condition; curliness; an undulation. (rare)]

- Faraday patterns became a standard topic in nonlinear physics due to experiments with liquids in the 1980s

Faraday waves in BEC

- Faraday patterns not relevant to BEC community until the seminal paper of Staliunas et al., PRL **89**, 210406 (2002)
- Their main point was that by periodically modulating the scattering length of a magnetically trapped 3D one excites a series of patterns similar to those in fluid mechanics
- The paper uses full 3D simulations to show the patterns in the density profile of the condensate but no systematic computations are performed. They use the Mathieu equations only to show that there is an instability
- Their PRA **70**, 011601 (2004) paper addresses cigar-shaped and pancake-like condensates, i.e. quasi-1D and 2D setups
- They show the formation of the waves through direct integration of the GP and give analytical arguments based on multiple-scale analysis

Experimental observation of Faraday waves

P. Engels, C. Atherton, M. A. Hoefer, PRL **98**, 095301 (2007)

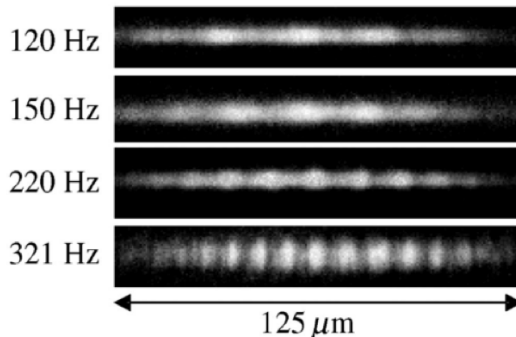


FIG. 1. In-trap absorption images of Faraday waves in a BEC. Frequency labels for each image represent the driving frequency at which the transverse trap confinement is modulated.

Two-component BEC systems

- Experimentally realized with a broad variety of types of atoms and parameters of a system
 - heterogeneous systems: different types of atoms
 - homogeneous systems: same type of atoms, different internal (usually spin) states
- Rich dynamics and interplay of the parameters
- Several possible ground states
- A variety of possible dynamical evolutions
- We focus on the study of Faraday waves and patterns in cigar-shaped two-component ^{87}Rb BECs, with strong radial confinement, which is harmonically modulated
- We also study resonant waves, which appear for specific values of the frequency of radial modulation

Mean-field description of 2C BEC

- The system is described by a coupled system of GP equations:

$$i\hbar \frac{\partial \Psi_1(\vec{r}, t)}{\partial t} = \left[-\frac{\hbar^2}{2m_1} \Delta + V(\vec{r}, t) + g_{11} |\Psi_1(\vec{r}, t)|^2 + g_{12} |\Psi_2(\vec{r}, t)|^2 \right] \Psi_1(\vec{r}, t)$$

$$i\hbar \frac{\partial \Psi_2(\vec{r}, t)}{\partial t} = \left[-\frac{\hbar^2}{2m_2} \Delta + V(\vec{r}, t) + g_{21} |\Psi_1(\vec{r}, t)|^2 + g_{22} |\Psi_2(\vec{r}, t)|^2 \right] \Psi_2(\vec{r}, t)$$

where the couplings are given by:

$$g_{11} = \frac{4\pi\hbar^2 a_1}{m_1}, \quad g_{22} = \frac{4\pi\hbar^2 a_2}{m_2}, \quad g_{12} = g_{21} = \frac{2\pi\hbar^2 a_{\text{int}}}{m_{\text{eff}}}$$

- Typical experimental values we consider for two hyperfine states of ^{87}Rb :

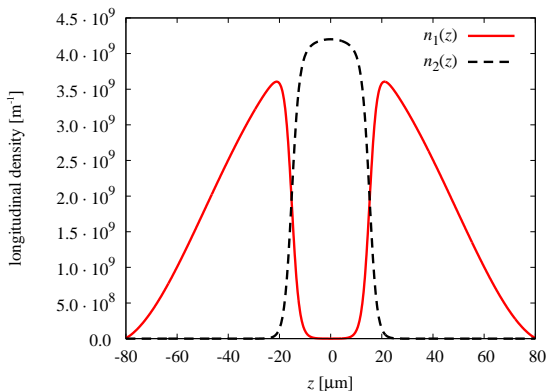
$$N_1 = 2.5 \cdot 10^5, \quad N_2 = 1.25 \cdot 10^5$$

$$a_1 = 100.4 a_0, \quad a_2 = 98.98 a_0, \quad a_{\text{int}} = 100.4 a_0$$

$$\omega_\rho(t) = \omega_{\rho,0}(1 + \epsilon \sin \omega_m t), \quad \omega_{\rho,0} = 160 \cdot 2\pi \text{ Hz}$$

$$\omega_m = 250 \cdot 2\pi \text{ Hz}, \quad \epsilon = 0.1, \quad \omega_z = 7 \cdot 2\pi \text{ Hz}$$

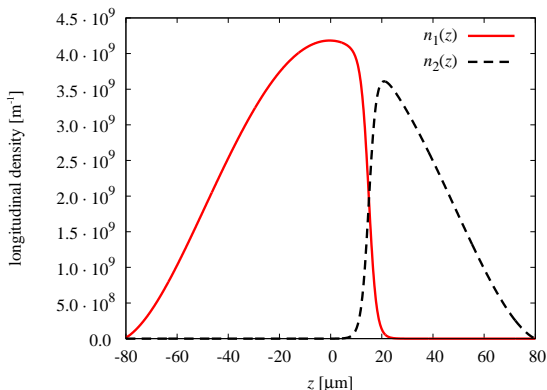
Ground state - imaginary-time propagation



Density profile of the converged eigenstate obtained by propagation in the imaginary time. Discretization parameters: $N_\rho = N_z = 2000$, $\varepsilon = 10^{-4}/\omega_z$.



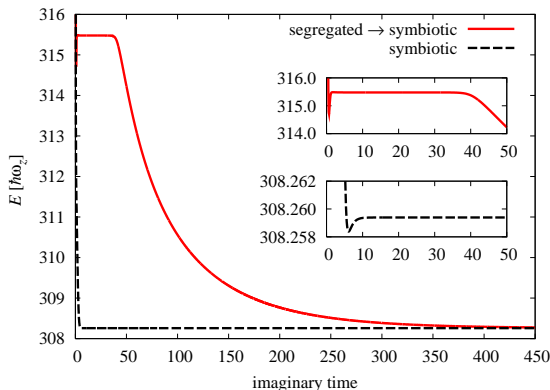
Segregated state - imaginary-time propagation



First excited eigenstate obtained by imaginary-time propagation.

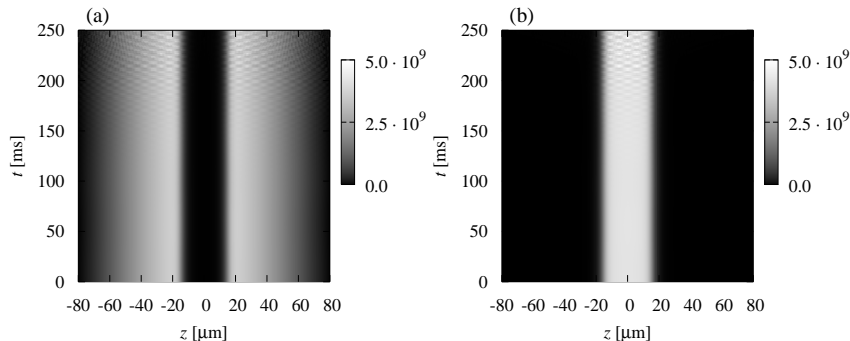
Discretization parameters: $N_\rho = N_z = 2000$, $\varepsilon = 10^{-4}/\omega_z$.

Energy: time dependence



Convergence of the total energy of the system during the imaginary-time propagation for the symbiotic pair and for the segregated state.

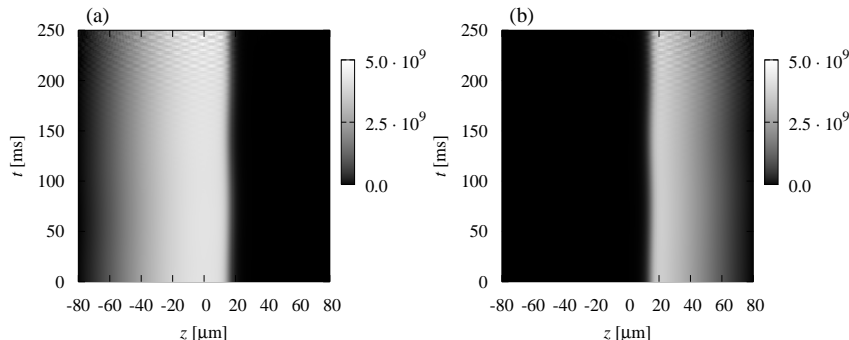
Symbiotic pair ground state - Faraday waves



Emergence of Faraday waves as a result of real-time propagation. The radial frequency of the trap is modulated at the non-resonant frequency

$$\omega_m = 250 \cdot 2\pi \text{ Hz}, \epsilon = 0.1.$$

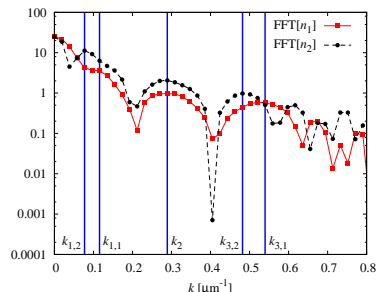
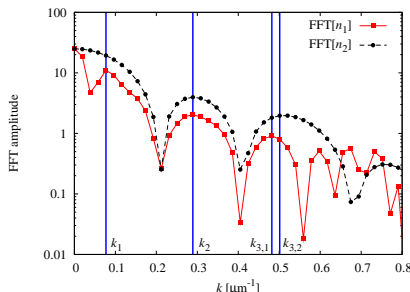
Segregated state - Faraday waves



Emergence of Faraday waves as a result of real-time propagation. The radial frequency of the trap is modulated at the non-resonant frequency

$$\omega_m = 250 \cdot 2\pi \text{ Hz}, \epsilon = 0.1.$$

Periods of Faraday waves



FFT of density profiles for the two condensates at $t = 200$ ms. For the symbiotic pair, the periods of waves are found to be $13.0 \mu\text{m}$ and $12.5 \mu\text{m}$, while for the segregated state the periods are $11.6 \mu\text{m}$ and $13.0 \mu\text{m}$.

Variational approach - symbiotic pair (1)

- Variational ansätze for the symbiotic pair wave functions:

$$\psi_1(\rho, z, t) = \mathcal{N}_1 \exp \left(-\frac{\rho^2}{2w_\rho^2(t)} + i\rho^2\alpha^2(t) \right) \left[1 - \exp \left(-\frac{z^2}{2w_z^2} \right) \right]$$

$$\psi_2(\rho, z, t) = \mathcal{N}_2 \exp \left(-\frac{\rho^2}{2w_\rho^2(t)} - \frac{z^2}{2w_z^2} + i\rho^2\alpha^2(t) \right) [1 + (u(t) + iv(t)) \cos kz]$$

- ψ_1 is considered to be unperturbed, acting as an additional confinement for ψ_2
- Variational analysis leads to a Mathieu-type equation:

$$\ddot{u}(\tau) + u(\tau) [a(k, \omega) + \epsilon b(k, \omega) \sin 2\tau] = 0$$

$$a(k, \omega) = \frac{k^4}{\omega^2} + \frac{k^2}{\omega^2} \Lambda_{\text{sym}}, \quad b(k, \omega) = \frac{k^2}{\omega^2} \Lambda_{\text{sym}}, \quad \omega t = 2\tau$$

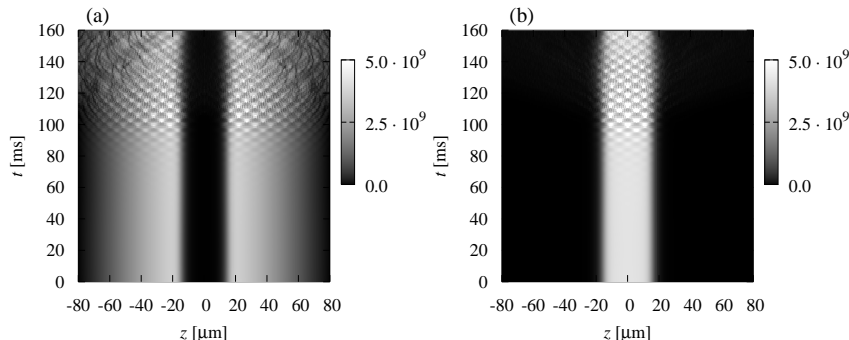
Variational approach - symbiotic pair (2)

- For small positive ϵ and positive $b(k, \omega)$, the Mathieu equation has solutions of the form $\exp(\pm i\mu\tau) \sin \sqrt{a}\tau$ and $\exp(\pm i\mu\tau) \cos \sqrt{a}\tau$, where $\text{Im}[\mu]$ consists of a series of lobes positioned around the solution of the equation $a(k, \omega) = n^2$
- The lobe centered around $a(k, \omega) = 1$ is the largest, and yields the most unstable solutions, determined by:

$$k_{F, \text{sym}} = \sqrt{-\frac{\Lambda_{\text{sym}}}{2}} + \sqrt{\frac{\Lambda_{\text{sym}}^2}{4} + \omega^2}$$

- This dispersion relation yields a period of $12.0 \mu\text{m}$ for the Faraday waves, which is in excellent agreement with the numerical results

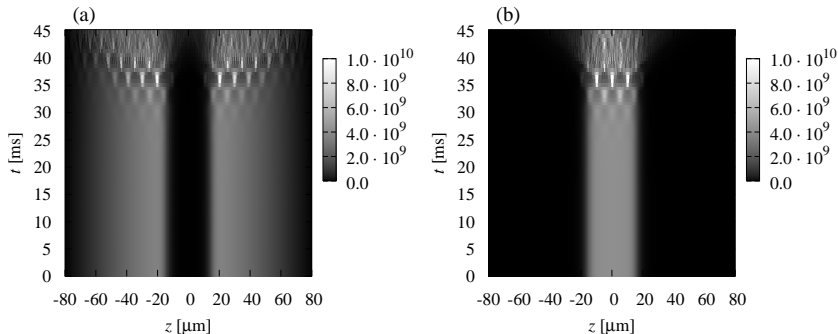
Self-resonance - symbiotic pair



Emergence of Faraday waves as a result of real-time propagation. The radial frequency of the trap is modulated at the resonant frequency

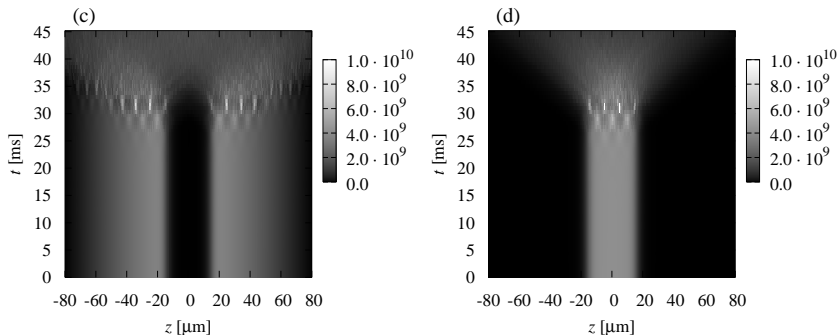
$$\omega_m = 160 \cdot 2\pi \text{ Hz}, \epsilon = 0.1.$$

Second-harmonic resonance - symbiotic pair (1)



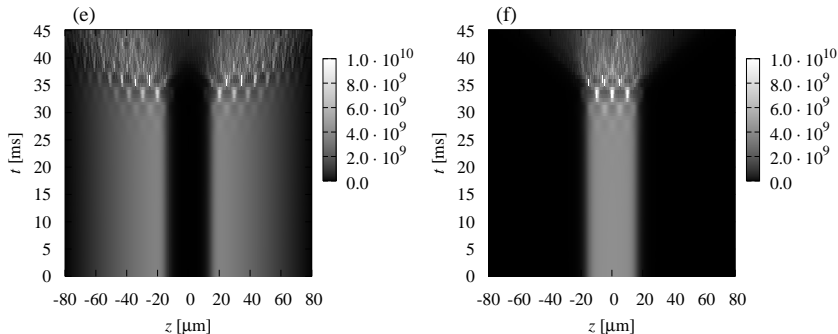
Emergence of resonant waves in a two-component BEC system radially modulated at $\omega_m = 300 \cdot 2\pi$ Hz.

Second-harmonic resonance - symbiotic pair (2)



Emergence of resonant waves in a two-component BEC system radially modulated at $\omega_m = 320 \cdot 2\pi$ Hz.

Second-harmonic resonance - symbiotic pair (3)



Emergence of resonant waves in a two-component BEC system radially modulated at $\omega_m = 340 \cdot 2\pi$ Hz.

Conclusions

- We have studied the emergence of surface waves in two-component BEC systems
- First, we have calculated two initial states of interest: ground state or symbiotic pair, and the first excited state or segregated state
- For non-resonant modulation of the radial confinement, the usual Faraday waves are observed, with the similar period in both components
- For the self-resonant modulation of the radial confinement, the expected resonant waves are observed
- For the second-harmonic resonance, much stronger and faster-emerging resonant waves are observed, turning the system from the non-miscible to the miscible state

Outlook

- Study of resonances using the Mathieu-type analysis
- Study of miscible two-component systems
- Pancake-shaped two-component systems
- Faraday waves for two-component BEC loaded into an optical lattice